

Diagrammatic Monte Carlo for the Hubbard model

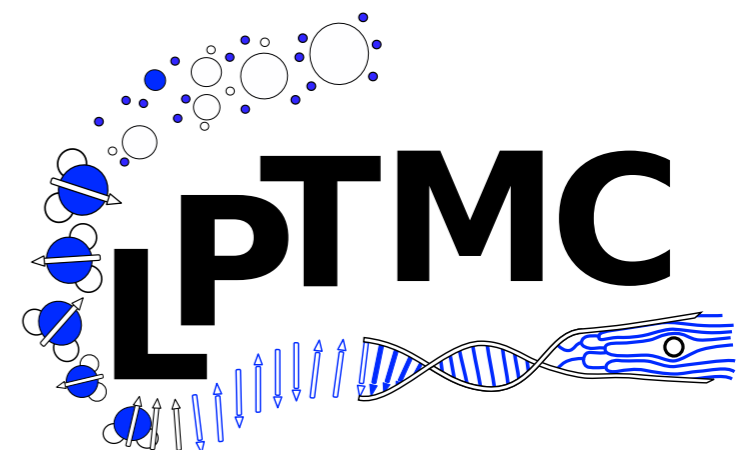
Riccardo Rossi

École Polytechnique Fédérale de Lausanne, Switzerland
CNRS, LPTMC, Sorbonne Université, Paris, France

RPMBT 22, Tsukuba, Japan

26 September 2024

EPFL



Plan

- How to compute lots of Feynman Diagrams (and why)
- Electrons interacting with AFM fluctuations
- Stripes signatures at intermediate temperature?
- Origin and fate of the pseudogap

Collaborators



Fedor Šimkovic*

Michel Ferrero

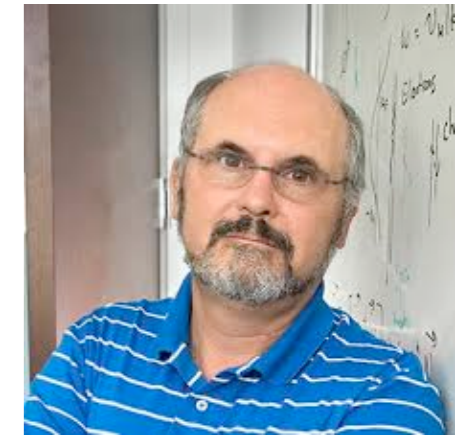
Renaud Garioud

Ecole Polytechnique, Paris



Antoine Georges

Flatiron Institute, New York



Alexei Tsvelik

Brookhaven NL



Gabriele Spada**

Félix Werner

Kris Van Houcke

Ecole Normale Supérieure, Paris



Igor Tupitsyn

Nikolay Prokof'ev

Boris Svistunov

UMass, Amherst


*now at IQM Munich

**now at University of Camerino

Why computing (a lot of) Feynman diagrams?

Fermionic sign problem

- Traditional Quantum Monte Carlo

$$\langle \hat{Q} \rangle := \frac{Z \hat{Q}}{Z} \quad \text{with } e^{\# \beta N_{\text{electrons}}}$$


unphysical non-locality for fermions

Signal-to-noise-ratio $\sim e^{-\# \beta N_{\text{electrons}}}$

Fermionic sign problem

- Traditional Quantum Monte Carlo

$$\langle \hat{Q} \rangle := \frac{Z \hat{Q}}{Z} \quad e^{\# \beta N_{\text{electrons}}}$$

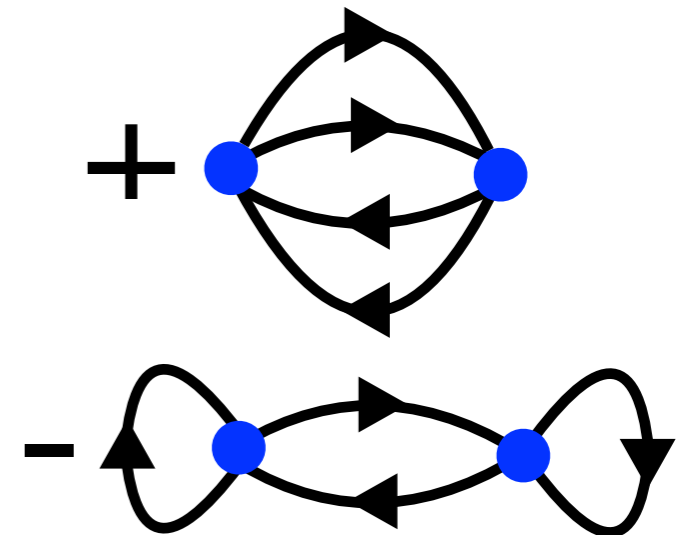
unphysical non-locality for fermions

$$\text{Signal-to-noise-ratio} \sim e^{-\# \beta N_{\text{electrons}}}$$

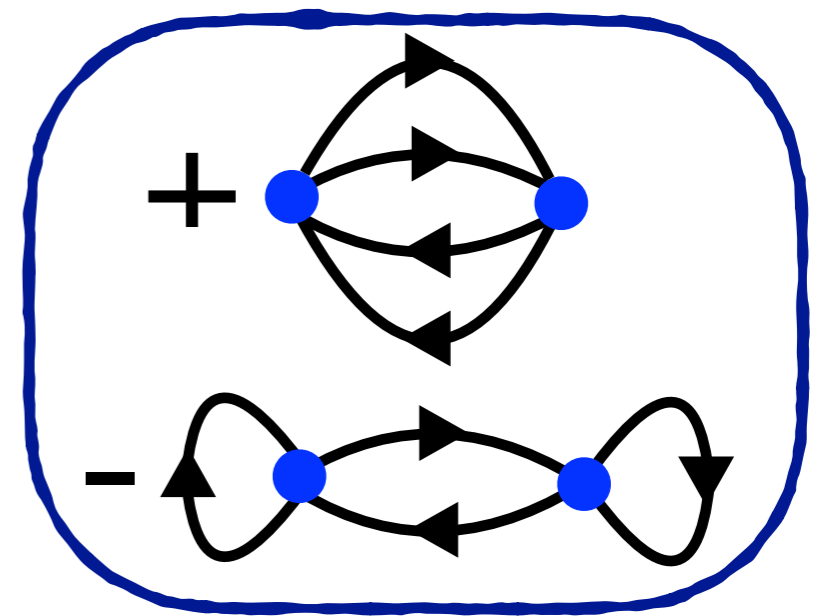
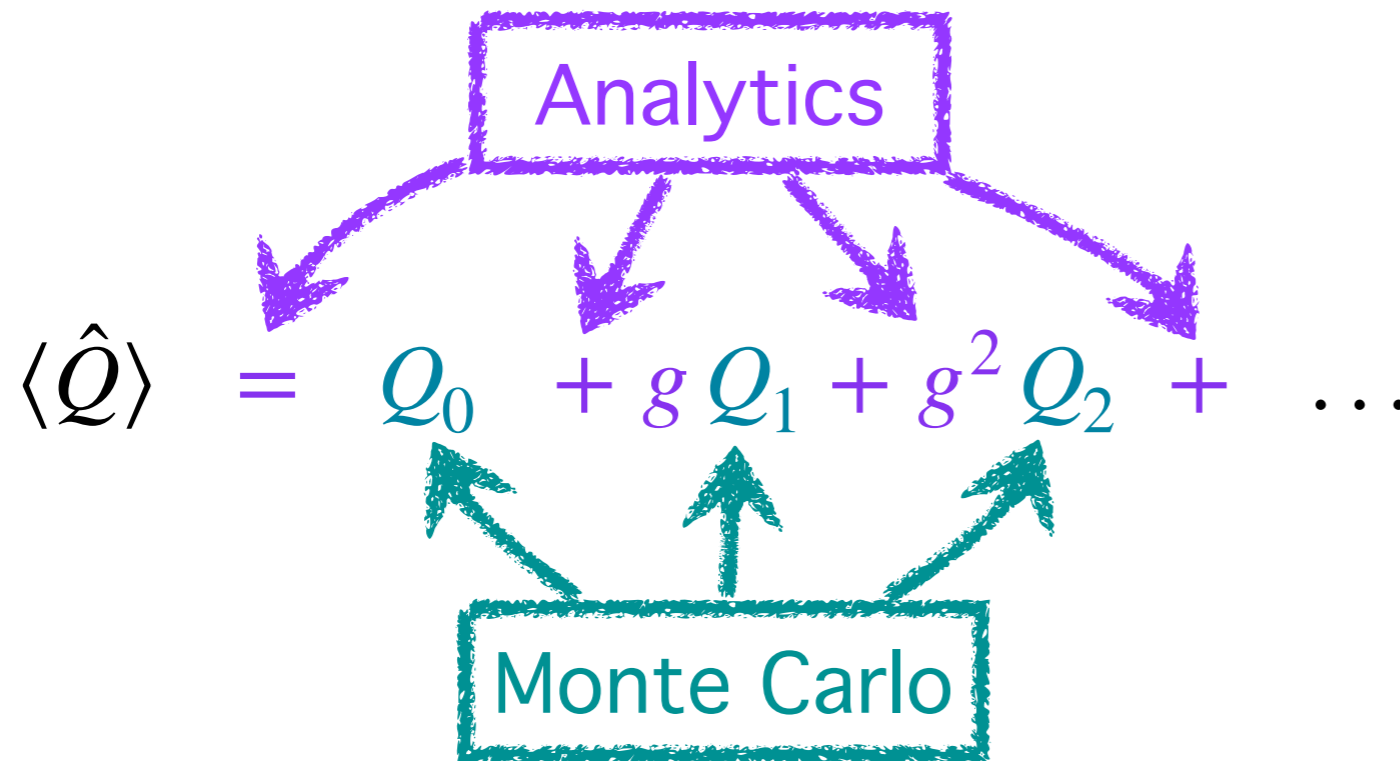
- Diagrammatic Monte Carlo

$$\langle \hat{Q} \rangle = Q_0 + g Q_1 + g^2 Q_2 + \dots$$

Small signal \iff Small correction



Diagrammatic Monte Carlo: cancellations are welcome



Bosons:

$$Q_n \sim n!$$

Fermions on a lattice:

=> analytic continuation

$$Q_n \sim (-R)^{-n}$$

Unitary Fermi Gas:

$$Q_n \sim (n!)^{1/5}$$

=> conformal-Borel resummation

Factorial cancellations

Prokof'ev, Svistunov, PRL 2008

[RR, Ohgoe, Van Houcke, Werner, PRL 2018]

Bare series convergence: lattice fermions $T > 0$

$$R_\beta := \lim_{L \rightarrow \infty} \min_n |z_{L,\beta;n}|$$

Theorem:

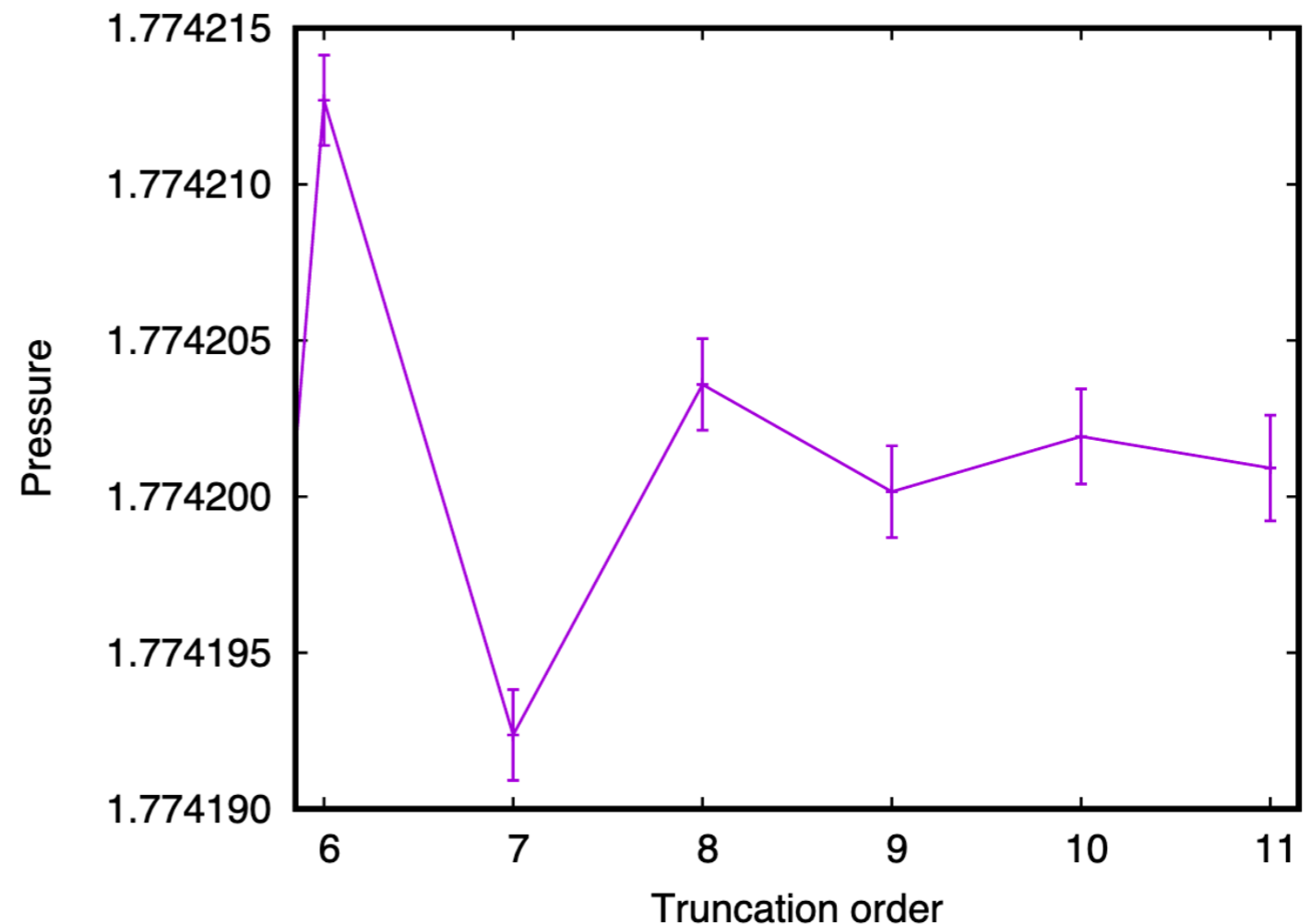
[Benfatto et al., Annales H. Poincaré 2006]:

$$R_\beta \underset{\beta \rightarrow \infty}{\gtrsim} \frac{1}{\log \beta}$$

Dyson's argument: $T > 0$ finite radius of convergence

Numerics:

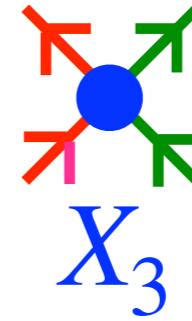
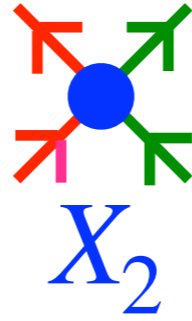
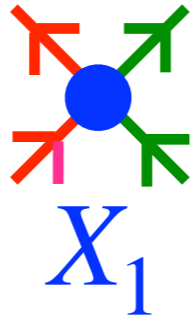
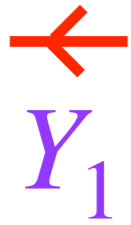
$T/t = 1/8$, $U/t = 2$, $n = 0.87500(2)$



How to compute (a lot of) Feynman diagrams?

CDet main idea

$O((n!)^2)$

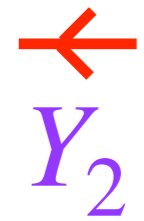
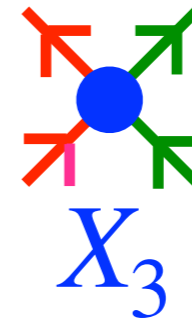
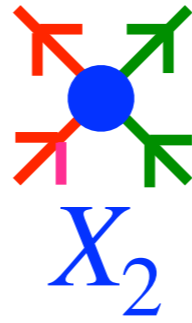
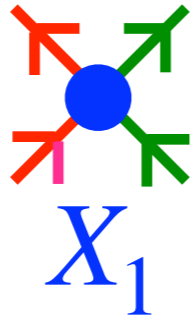
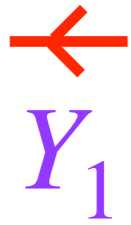


Original DiagMC

[Prokof'ev, Svistunov, PRL 1998]

CDet main idea

$O((n!)^2)$



Original DiagMC

[Prokof'ev, Svistunov, PRL 1998]

$$A_{Y_1, Y_2}(\{X_1, X_2, X_3\}) := \det \begin{pmatrix} G(X_1, X_1) & G(X_1, X_2) & G(X_1, X_3) \\ G(X_2, X_1) & G(X_2, X_2) & G(X_2, X_3) \\ G(X_3, X_1) & G(X_3, X_2) & G(X_3, X_3) \end{pmatrix} \times$$

CT-INT

[Rubtsov, arXiv2003]

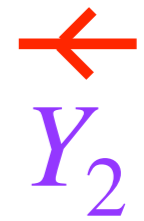
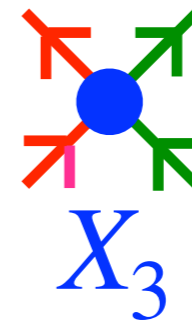
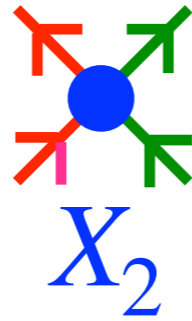
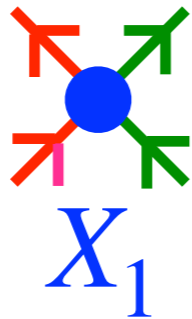
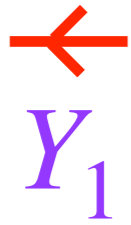
[Rubtsov, Savkin, Lichtenstein, PRB 2005]

[Bourovski, Prokof'ev, Svistunov, PRB 2004]

$$\det \begin{pmatrix} G(X_1, X_1) & G(X_1, X_2) & G(X_1, X_3) & G(X_1, Y_2) \\ G(X_2, X_1) & G(X_2, X_2) & G(X_2, X_3) & G(X_1, Y_2) \\ G(X_3, X_1) & G(X_3, X_2) & G(X_3, X_3) & G(X_1, Y_2) \\ G(Y_1, X_1) & G(Y_1, X_2) & G(Y_1, X_3) & G(Y_1, Y_2) \end{pmatrix}$$

CDet main idea

$O((n!)^2)$



Original DiagMC

[Prokof'ev, Svistunov, PRL 1998]

$$A_{Y_1, Y_2}(\{X_1, X_2, X_3\}) := \det \begin{pmatrix} G(X_1, X_1) & G(X_1, X_2) & G(X_1, X_3) \\ G(X_2, X_1) & G(X_2, X_2) & G(X_2, X_3) \\ G(X_3, X_1) & G(X_3, X_2) & G(X_3, X_3) \end{pmatrix} \times$$

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$$\det \begin{pmatrix} G(X_1, X_1) & G(X_1, X_2) & G(X_1, X_3) & G(X_1, Y_2) \\ G(X_2, X_1) & G(X_2, X_2) & G(X_2, X_3) & G(X_1, Y_2) \\ G(X_3, X_1) & G(X_3, X_2) & G(X_3, X_3) & G(X_1, Y_2) \\ G(Y_1, X_1) & G(Y_1, X_2) & G(Y_1, X_3) & G(Y_1, Y_2) \end{pmatrix}$$

CDet

$$C_{Y_1, Y_2}(\{X_1, X_2, X_3\}) := A_{Y_1, Y_2}(\{X_1, X_2, X_3\}) - \sum_{S \subsetneq \{X_1, X_2, X_3\}} C_{Y_1, Y_2}(S) A_{Y_1, Y_2}(\{X_1, X_2, X_3\} \setminus S)$$

[RR, PRL'17]

$O(3^n)$

Polynomial complexity of lattice fermions at non-zero T

CDet CPU time: $t_N \sim 3^N \times \left(\frac{|U|}{R_{A,\beta}} \right)^{2N}$

recursive formula Monte Carlo variance

Exponential cost vs
Exponential convergence:

$$t(\epsilon) \sim \epsilon^{-\alpha}$$

[RR et al, EPL'16]
[Troyer, Wiese, PRL'05]

Polynomial complexity of lattice fermions at non-zero T

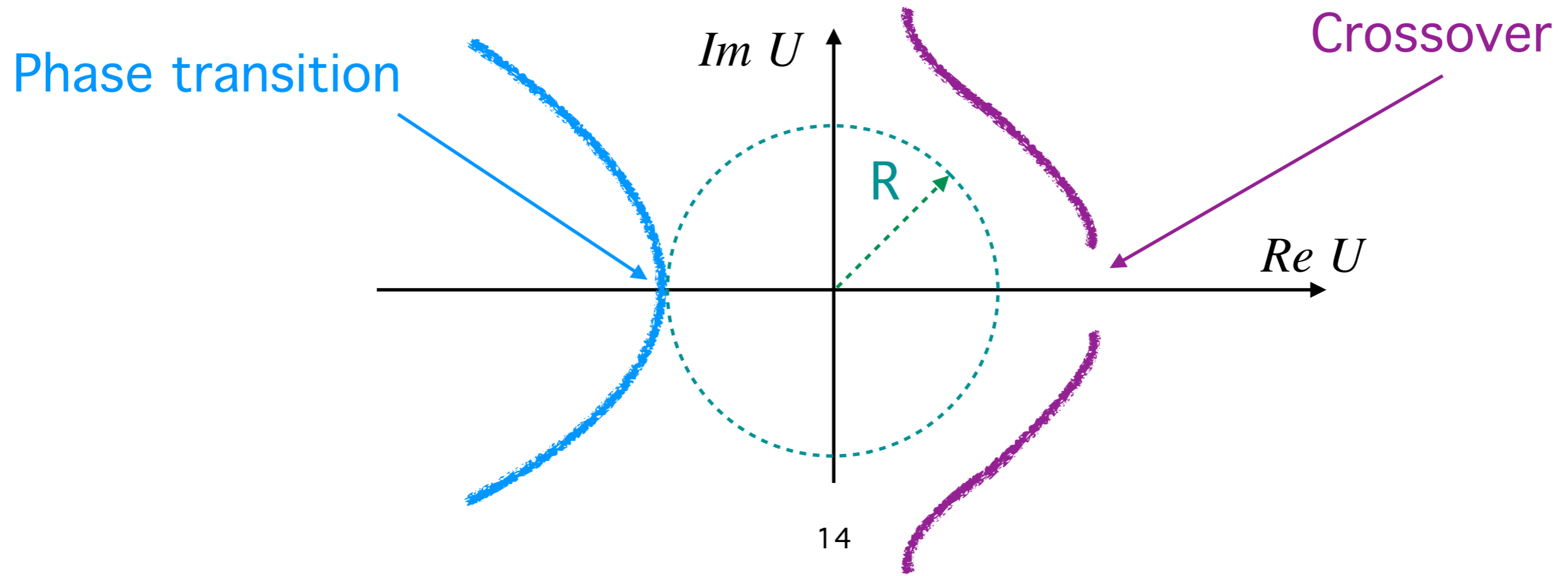
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recursive formula Monte Carlo variance

Exponential cost vs
Exponential convergence:

$$t(\epsilon) \sim \epsilon^{-\alpha}$$

[RR et al, EPL'16]
[Troyer, Wiese, PRL'05]



Algebraic renormalization

$$C[G_0, U] = \frac{A[G_0, U]}{Z[G_0, U]}, \quad C(\{X_1, \dots, X_n\}) := \frac{\delta^n}{\delta U(X_1) \dots U(X_n)} C[G_0, U] \Big|_{U=0}$$

Nilpotent polynomials (aka hyperdual numbers)

$$U(z)(X) = \sum_{j=1}^n z_j \delta(X - X_j), \quad [z_j, z_k] = 0, \quad z_j^2 = 0$$

$$C(z) := C[G_0, U(z)] = \frac{A[G_0, U(z)]}{Z[G_0, U(z)]} = \frac{A(z)}{Z(z)}, \quad \frac{\partial^n}{\partial z_1 \dots \partial z_n} C(z) = C(\{X_1, \dots, X_n\})$$

\Rightarrow CDet is just a hyperdual-number polynomial division

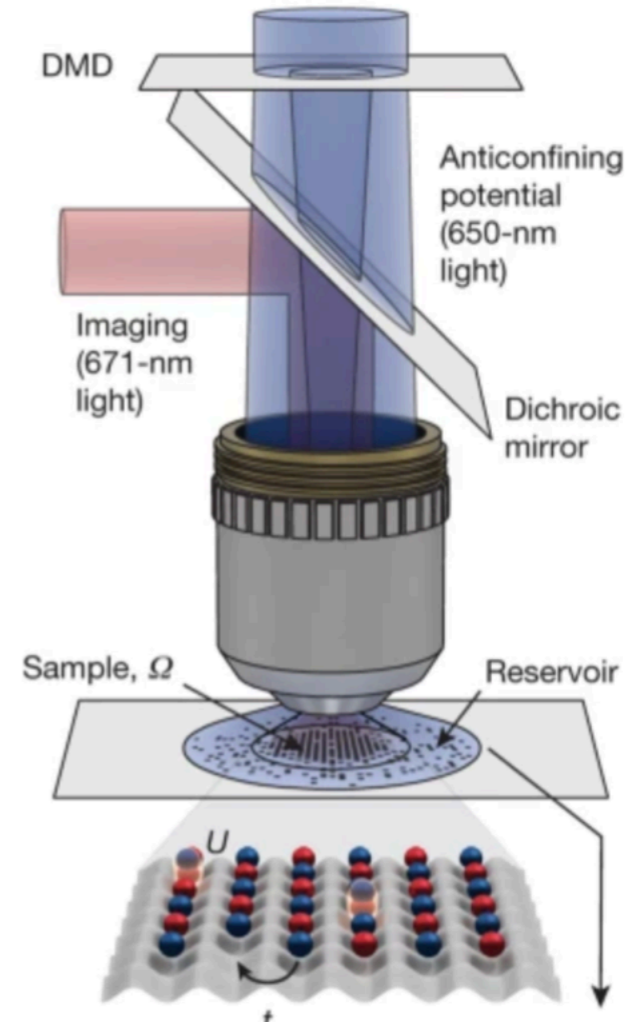
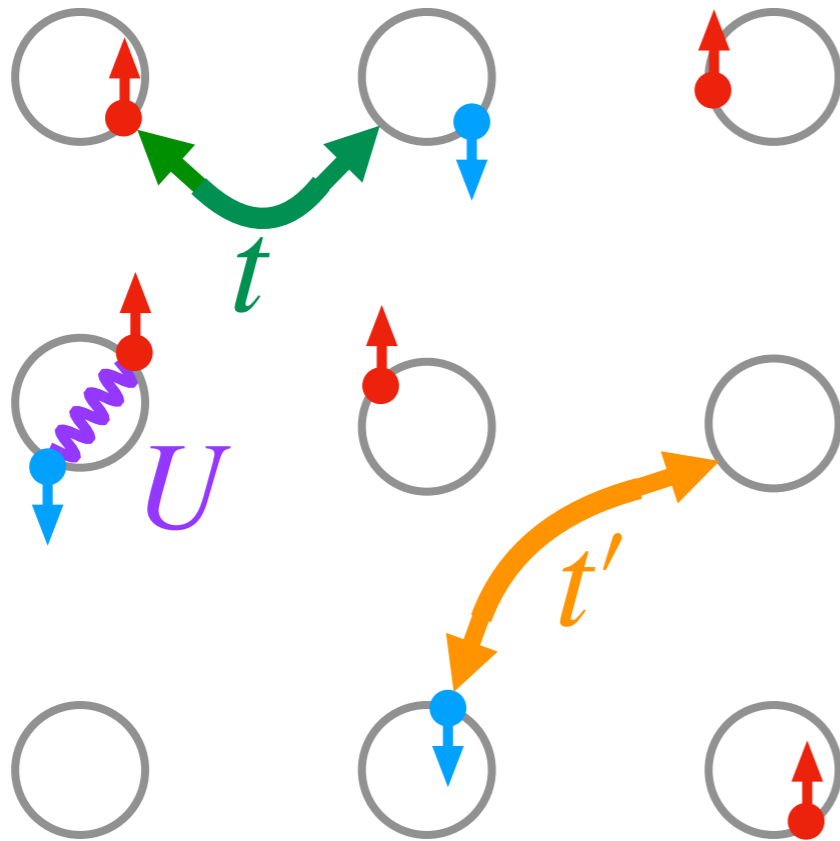
Renormalization $\equiv [G_0(z), \Gamma_0(z)]$

\Rightarrow „Bare“ CDet is all one needs

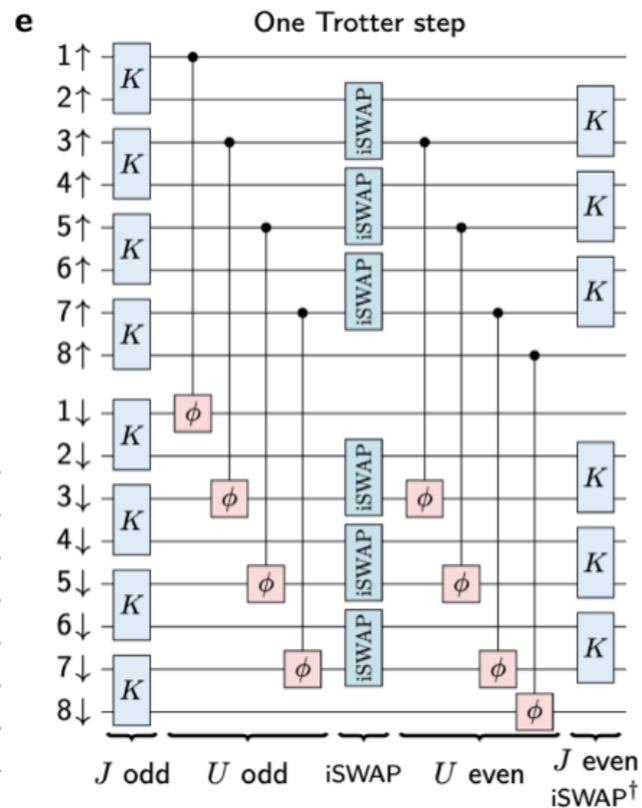
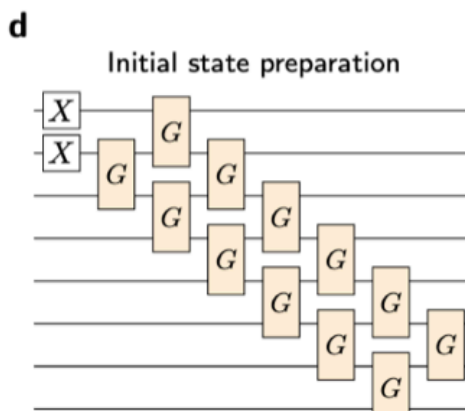
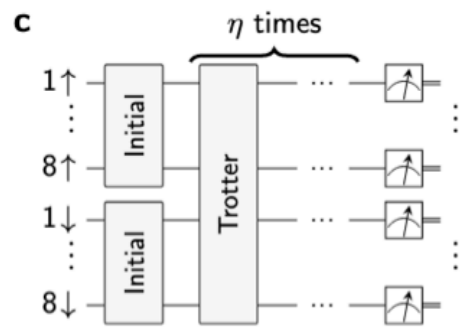
Application: fix density in g.c. ensemble by $\mu(w)$, where $w^{n+1} = 0$

2D Hubbard model

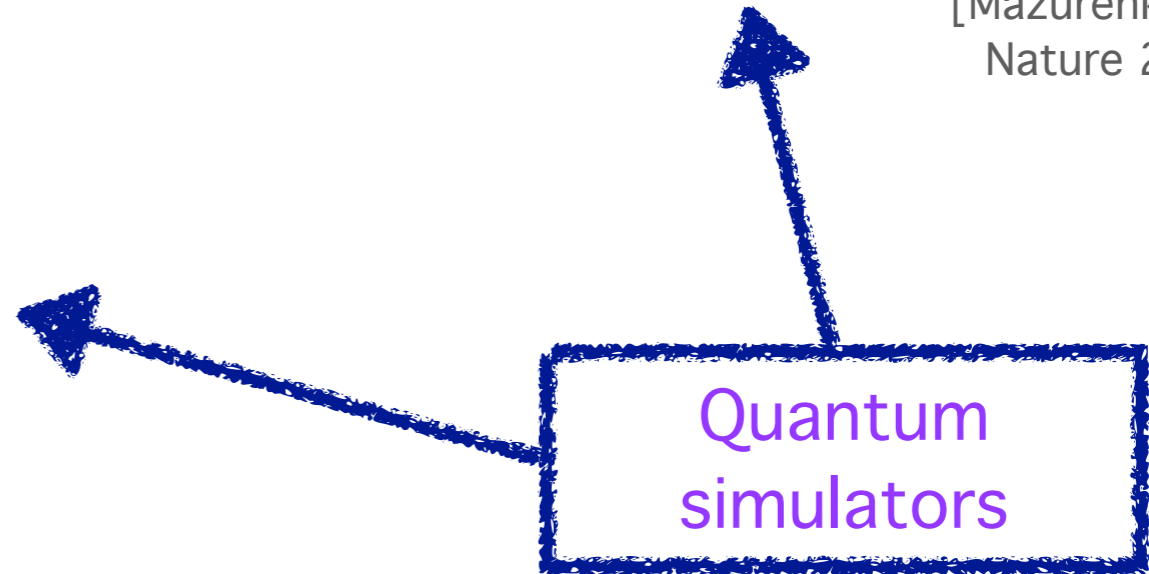
Fermi-Hubbard model



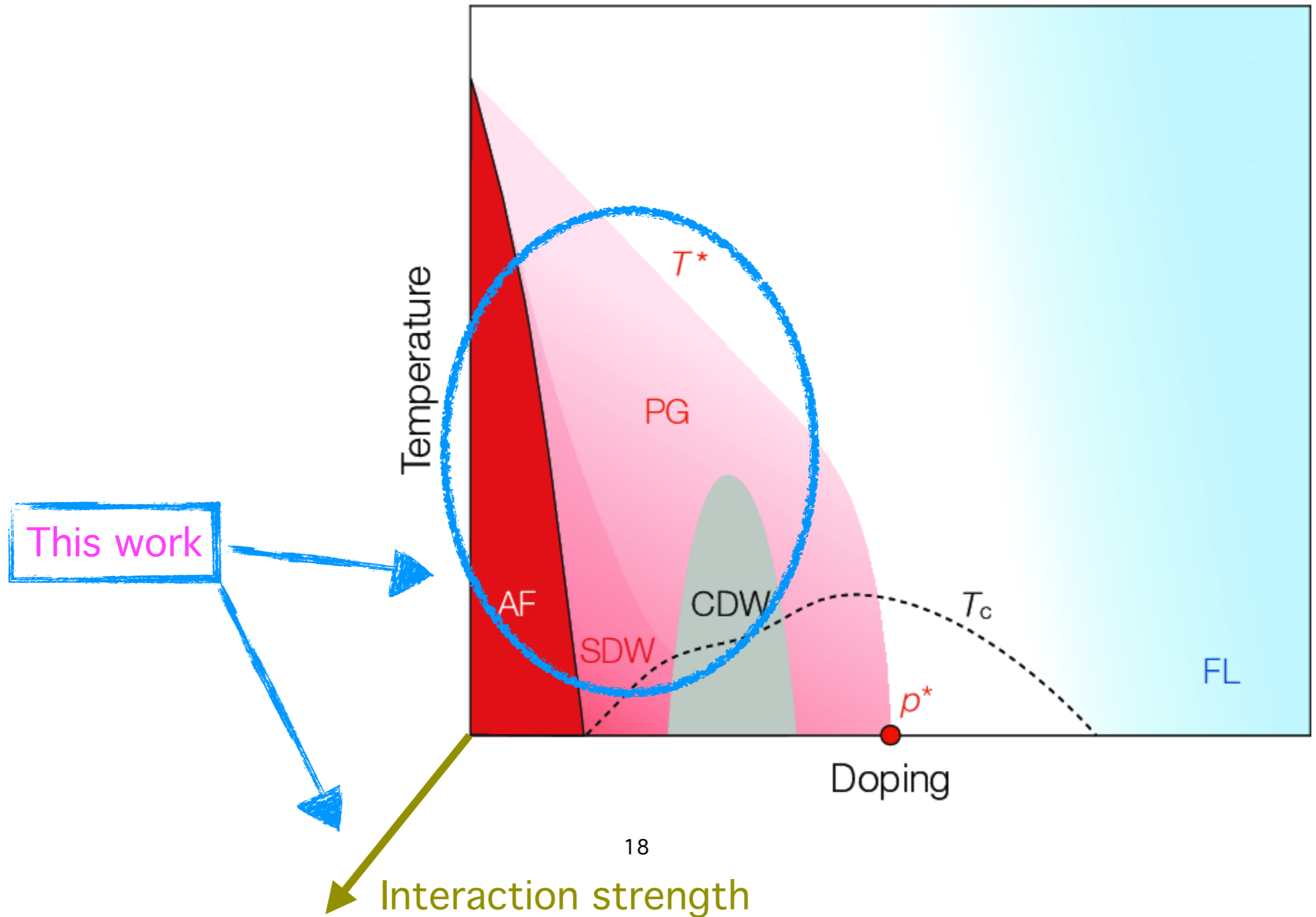
[Mazurenko et al, Nature 2017]



[Google AI Quantum, 2020]



Expected finite T situation in 2D (rough sketch)



Setting for the numerical experiment and probes

Square lattice

$$\hat{H} = \sum_{\mathbf{k}, \sigma} \epsilon_{\mathbf{k}} \hat{c}_{\mathbf{k}\sigma}^\dagger \hat{c}_{\mathbf{k}\sigma} + U \sum_{\mathbf{r}} \hat{n}_{\uparrow}(\mathbf{r}) \hat{n}_{\downarrow}(\mathbf{r}) - \mu \sum_{\mathbf{r}} \hat{n}(\mathbf{r})$$

$$\epsilon_{\mathbf{k}} = -2t(\cos k_x + \cos k_y) - 4t' \cos k_x \cos k_y$$

$$\langle \hat{O} \rangle = \frac{\text{Tr} e^{-\hat{H}/T} \hat{O}}{\text{Tr} e^{-\hat{H}/T}}$$

$$\hat{O}(\tau) = e^{\tau \hat{H}} \hat{O} e^{-\tau \hat{H}}$$

Self-energy

$$G(\mathbf{r}, \tau) = - \langle \hat{c}_{\uparrow}(\mathbf{r}, \tau) \hat{c}_{\uparrow}^\dagger(\mathbf{0}) \rangle$$

$$G^{-1} = G_0^{-1} - \Sigma$$

Spectral function proxy

$$A(\mathbf{k}) = \frac{-1}{\pi} \text{Im} G_{k, i\omega_0}$$

Spin susceptibility

$$\chi_{sp}(\mathbf{r}) = \int_0^{1/T} d\tau \langle \hat{S}_z(\mathbf{r}, \tau) \hat{S}_z(\mathbf{0}, \tau) \rangle$$

Charge susceptibility

$$\chi_{ch}(\mathbf{r}) = \int_0^{1/T} d\tau [\langle \hat{n}(\mathbf{r}, \tau) \hat{n}(\mathbf{0}, 0) \rangle - \langle \hat{n}(\mathbf{r}) \rangle \langle \hat{n}(\mathbf{0}) \rangle]$$

Strong-coupling probes

$$D = \langle \hat{n}_{\uparrow}(\mathbf{r}) \hat{n}_{\downarrow}(\mathbf{r}) \rangle \quad \langle \hat{n}(\mathbf{r}) \rangle \text{ versus } \mu$$

Alternatives

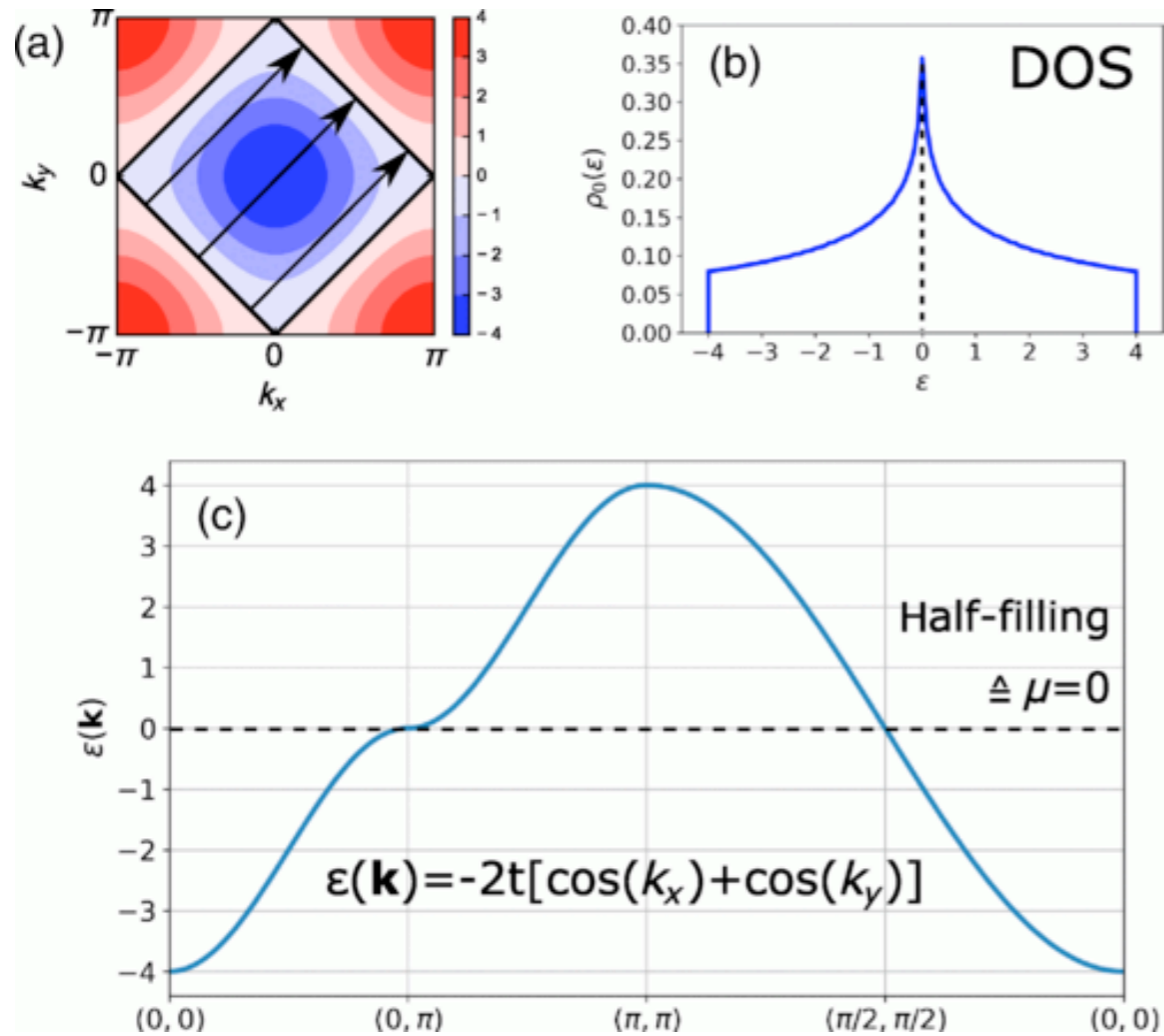
$$S_{sp}(\mathbf{r}) = \langle \hat{S}_z(\mathbf{r}) \hat{S}_z(\mathbf{0}) \rangle$$

$$S_{ch}(\mathbf{r}) = \langle \hat{n}(\mathbf{r}) \hat{n}(\mathbf{0}) \rangle$$

Fermions interacting with AFM fluctuations

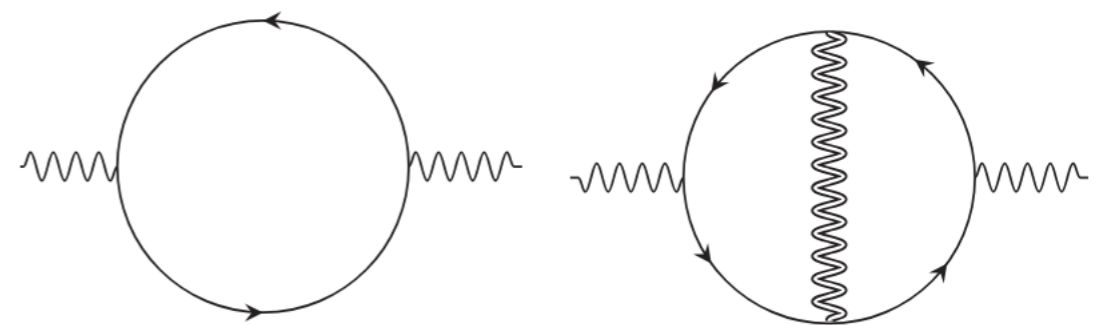
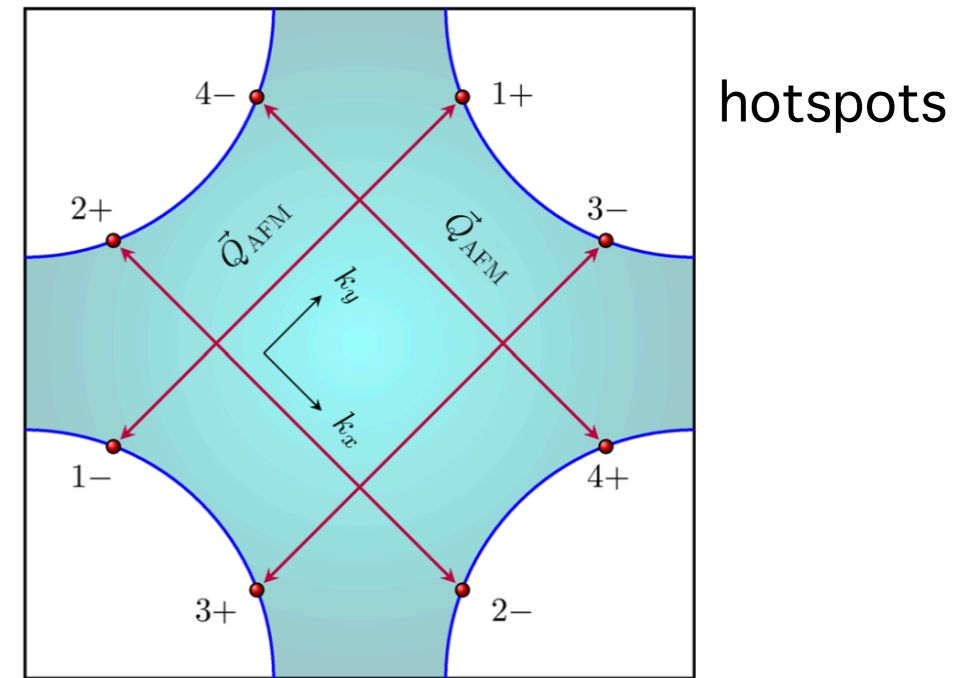
Previous works

Half-filled Hubbard with ph sym: Perfect nesting



[Simkovic et al, PRL 2020]

[Schafer et al, PRX 2021]



$$D(\mathbf{q})^{-1} = |\mathbf{q}_0| + c(v)[|\mathbf{q}_x| + |\mathbf{q}_y|]$$

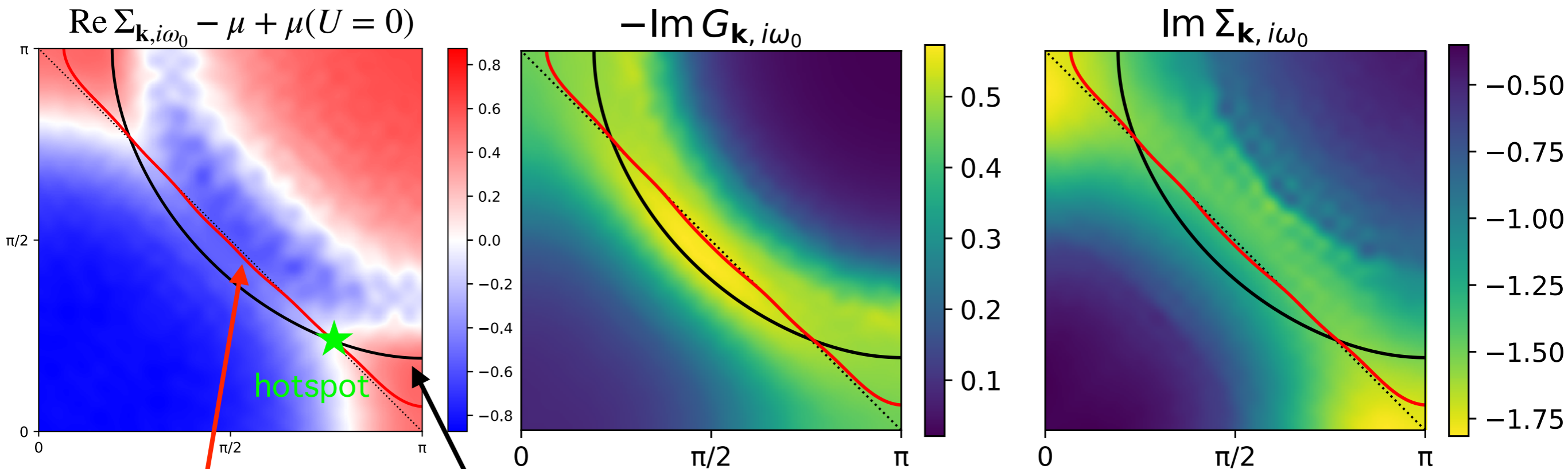
[Schlief et al, PRX 2017]

Our numerical experiment: Half-filled Hubbard model with t'

[RR, Šimkovic, Ferrero, Georges, Tsvelik, Prokofiev, Tupitsyn, PRR'24]

Flowing toward nesting? Hotspots?

Half filling, $U/t = 5.75$, $t'/t = -0.3$, $T/t = 1/7$



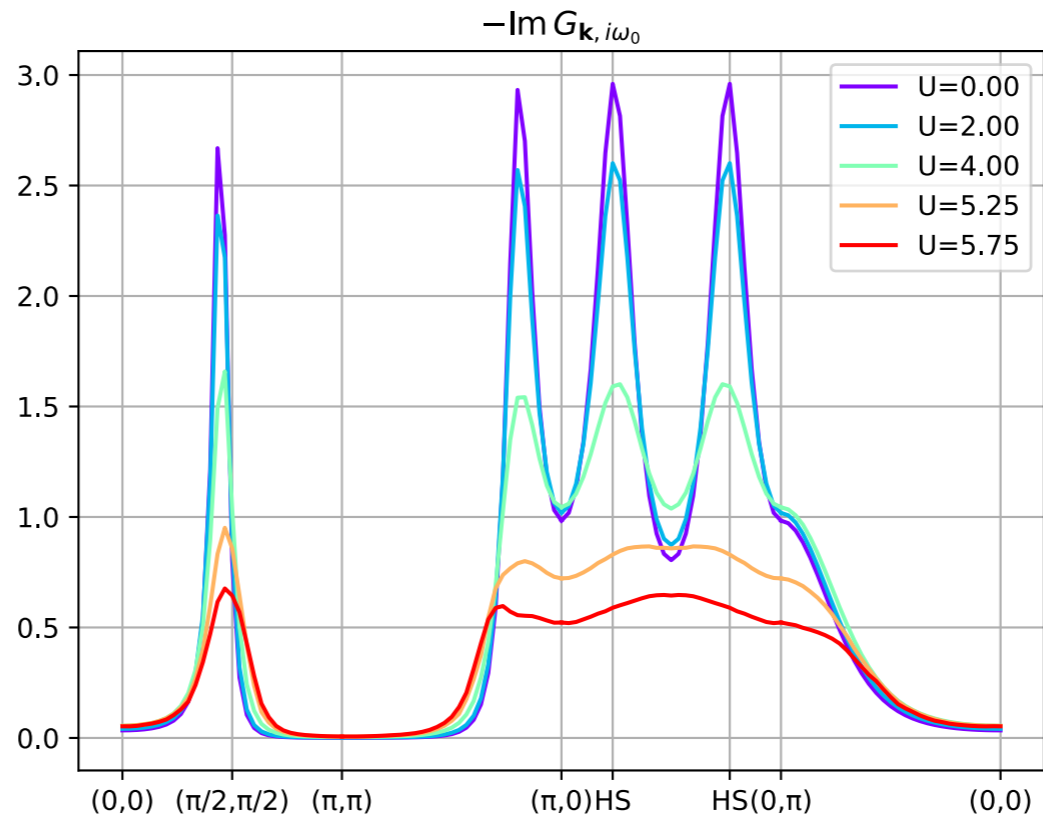
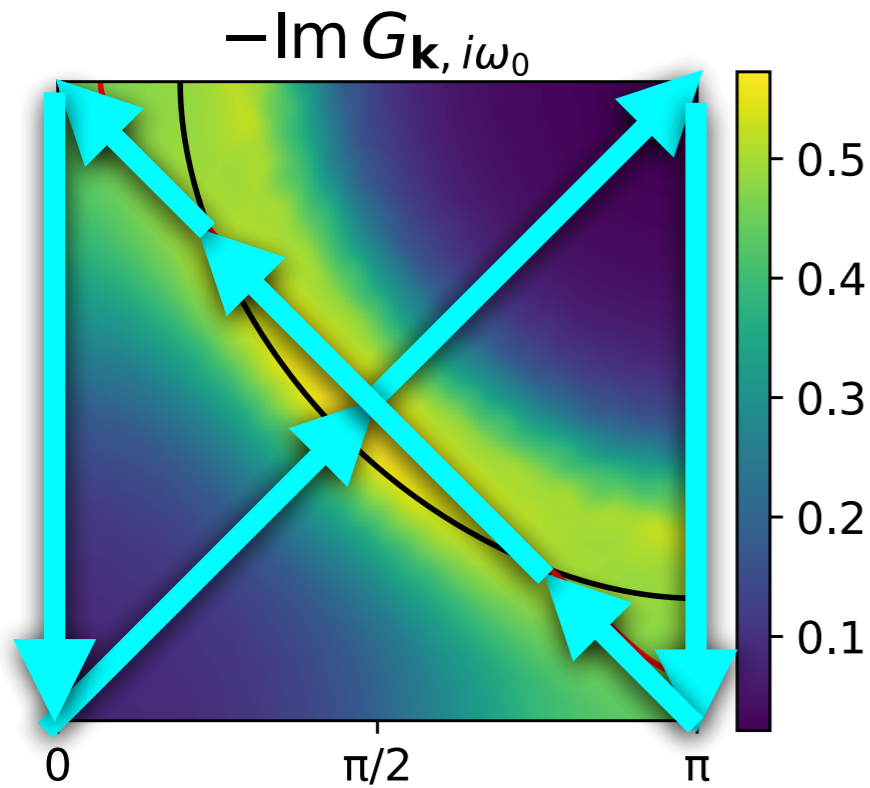
Line of zeros of renormalized dispersion relation

Non interacting Fermi surface

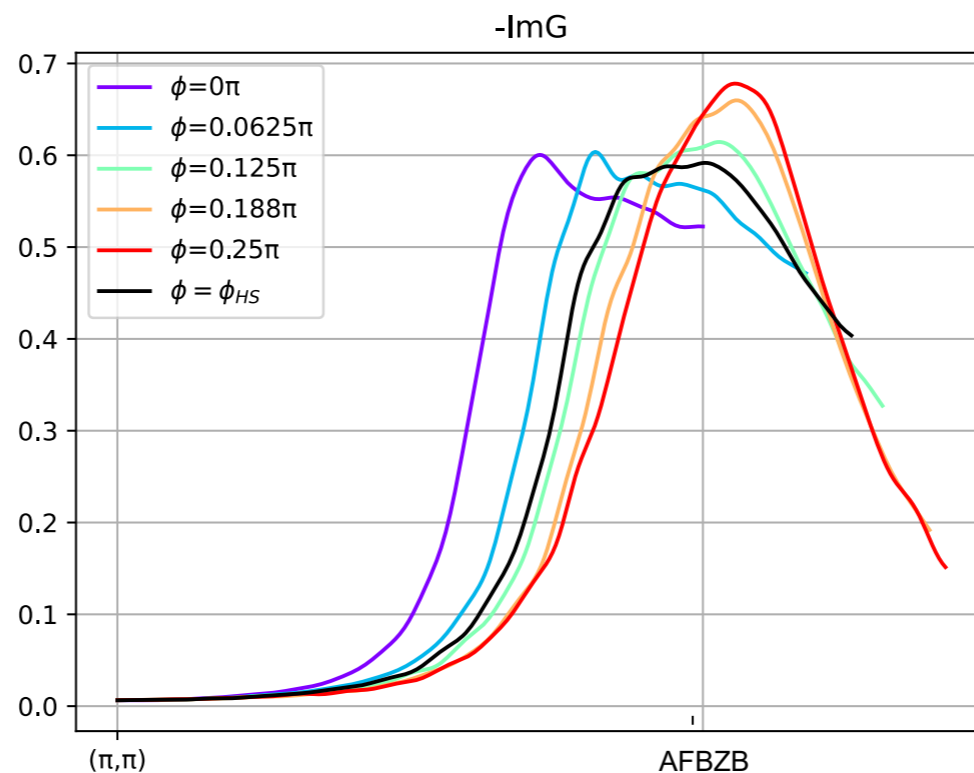
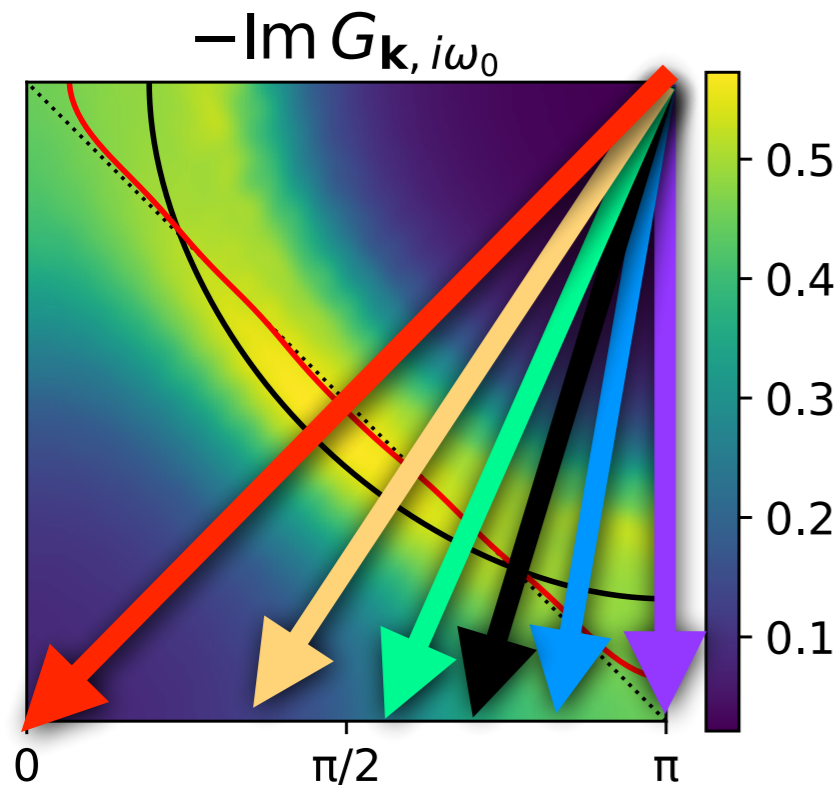
AFM spin fluctuations $\Sigma_{\mathbf{k}, i\omega_n} \sim \frac{1}{i\omega_n - \xi_{\mathbf{k}+\mathbf{Q}}}$

$\epsilon_{\mathbf{k}} + \text{Re } \Sigma_{\mathbf{k}, i\omega_0} - \mu = 0 \neq \text{FS} = \text{argmax}[-\text{Im}G_{\mathbf{k}, i\omega_0}]$ Σ pole developing at antinodes, not at the hotspots!

Fermi surface reconstruction tendencies



Spectral intensity is minimal at the hotspots



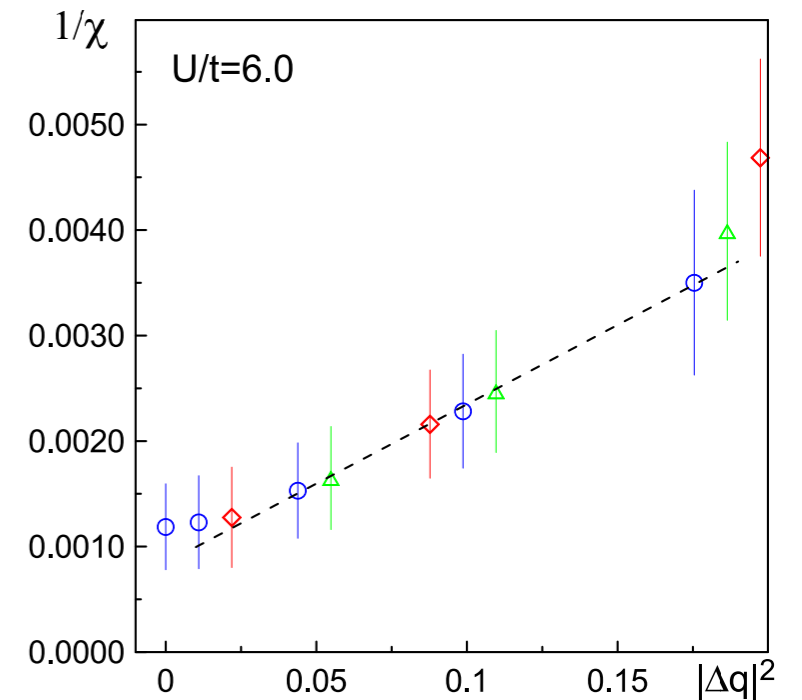
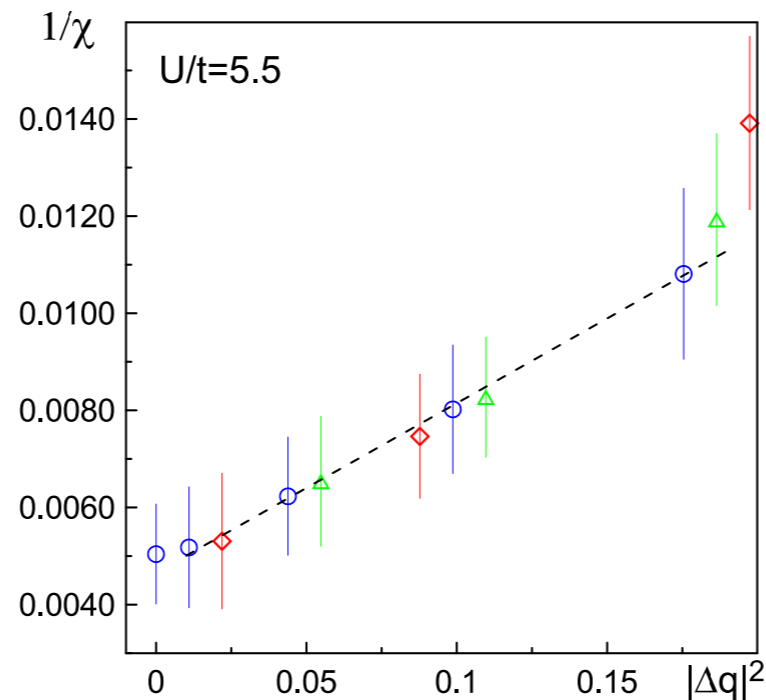
-> Signature of FS reconstruction without magnetic order

Exact RG fixed point of Schlieff-Lunts-Lee?

$$\chi_{sp}^{-1}(\mathbf{k}, \omega = \mathbf{0}) = \mathbf{c}(|\mathbf{k}_x| + |\mathbf{k}_y|)$$

[Schlieff et al, PRX 2017]

Hubbard model:
no sign of linear scaling
 for as low as $T/t = 1/10$

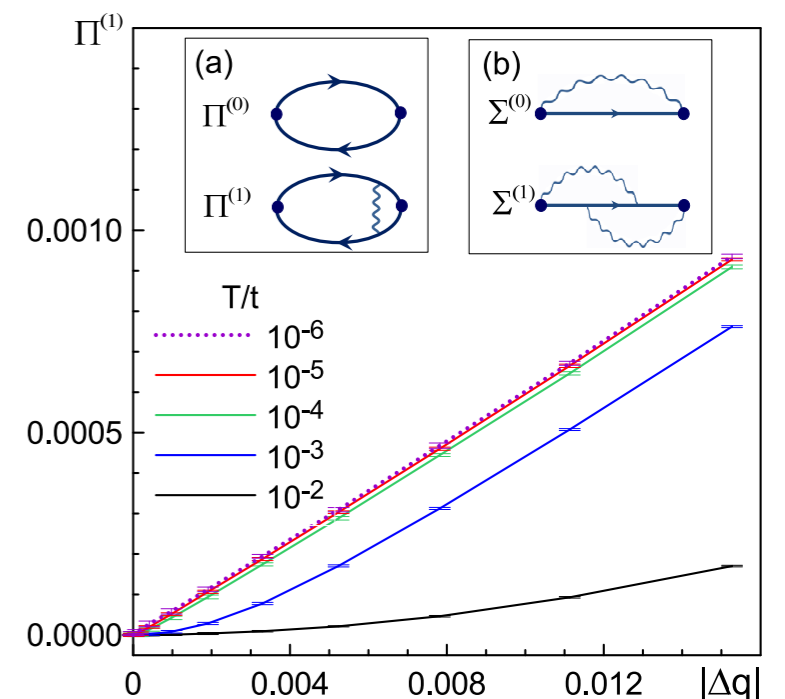


Spin-Fermion model:

$$\hat{H} = \sum_{\mathbf{k}, \alpha} \epsilon_{\mathbf{k}} c_{\mathbf{k}, \alpha}^{\dagger} c_{\mathbf{k}, \alpha} + \sum_{\mathbf{k}} \chi_0^{-1}(\mathbf{k}) \mathbf{S}(\mathbf{k}) \cdot \mathbf{S}(-\mathbf{k}) + g \sum_{\mathbf{k}, \mathbf{q}, \alpha, \beta, i} c_{\mathbf{k}+\mathbf{q}, \alpha}^{\dagger} \sigma_{\alpha, \beta}^i c_{\mathbf{k}, \beta} S_{-q}^i$$

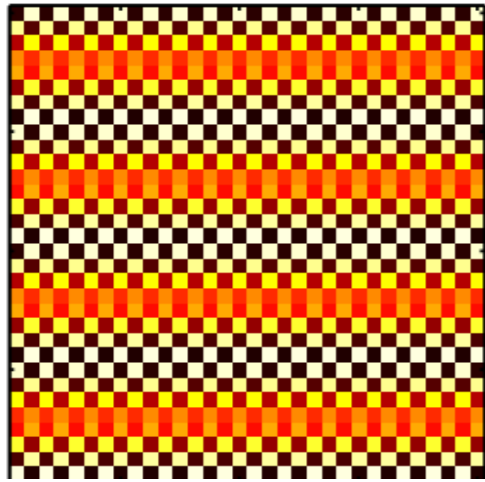
No trace of linear scaling for
 as low $T/t = 10^{-4}$

See also
 [Bauer et al, PRR 2020]



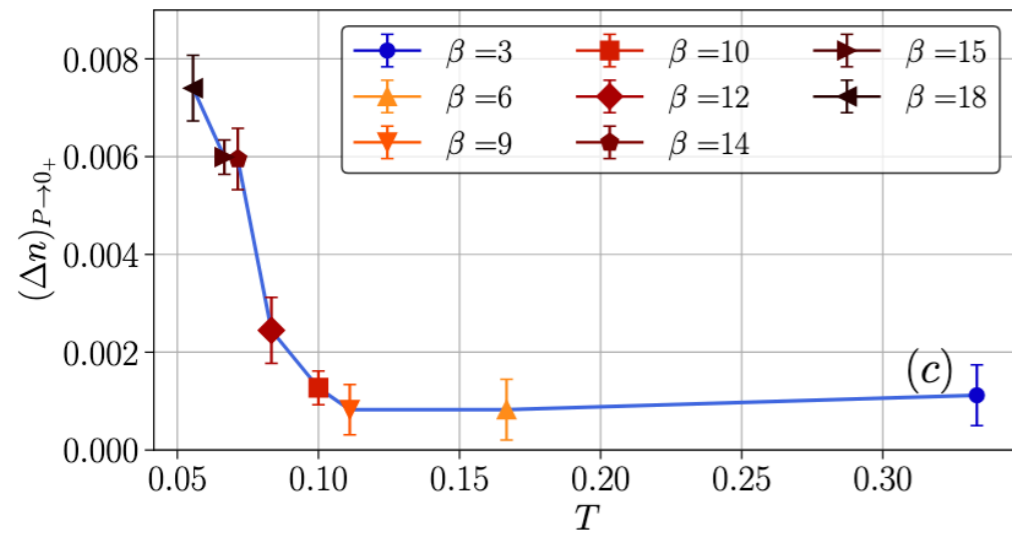
Stripe precursors at
intermediate temperature?

Stripes in the Ground state of the HM ($t'=0$)



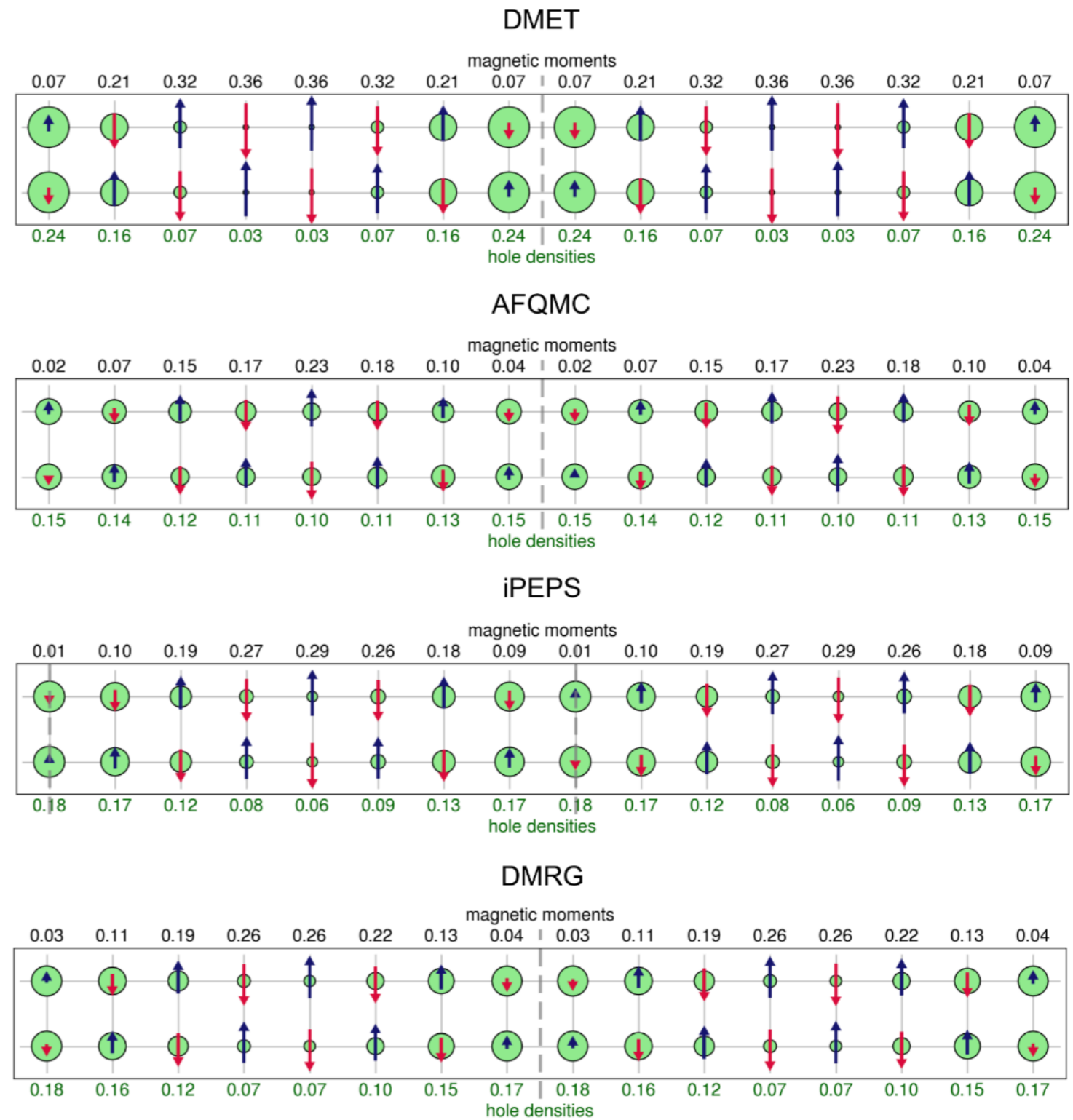
[Xu et al, J. Phys.: Condens. Matter 2011]

Also $T > 0$



[Xiao et al, PRX 2023]

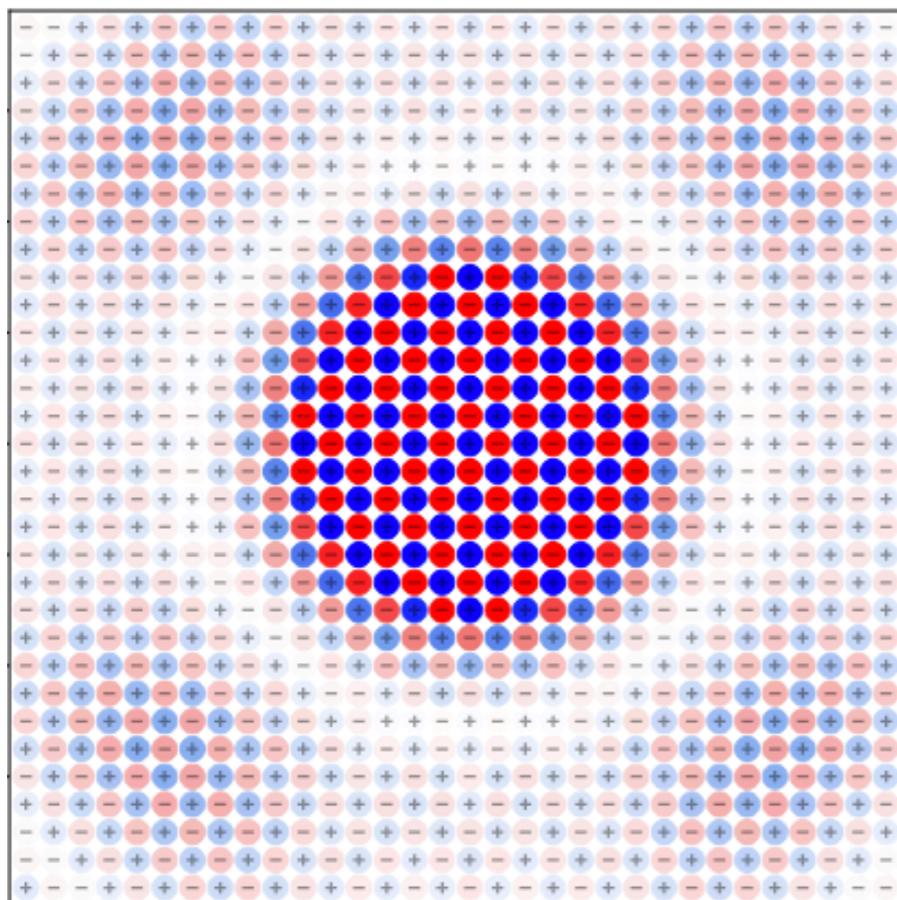
with pinning+external fields



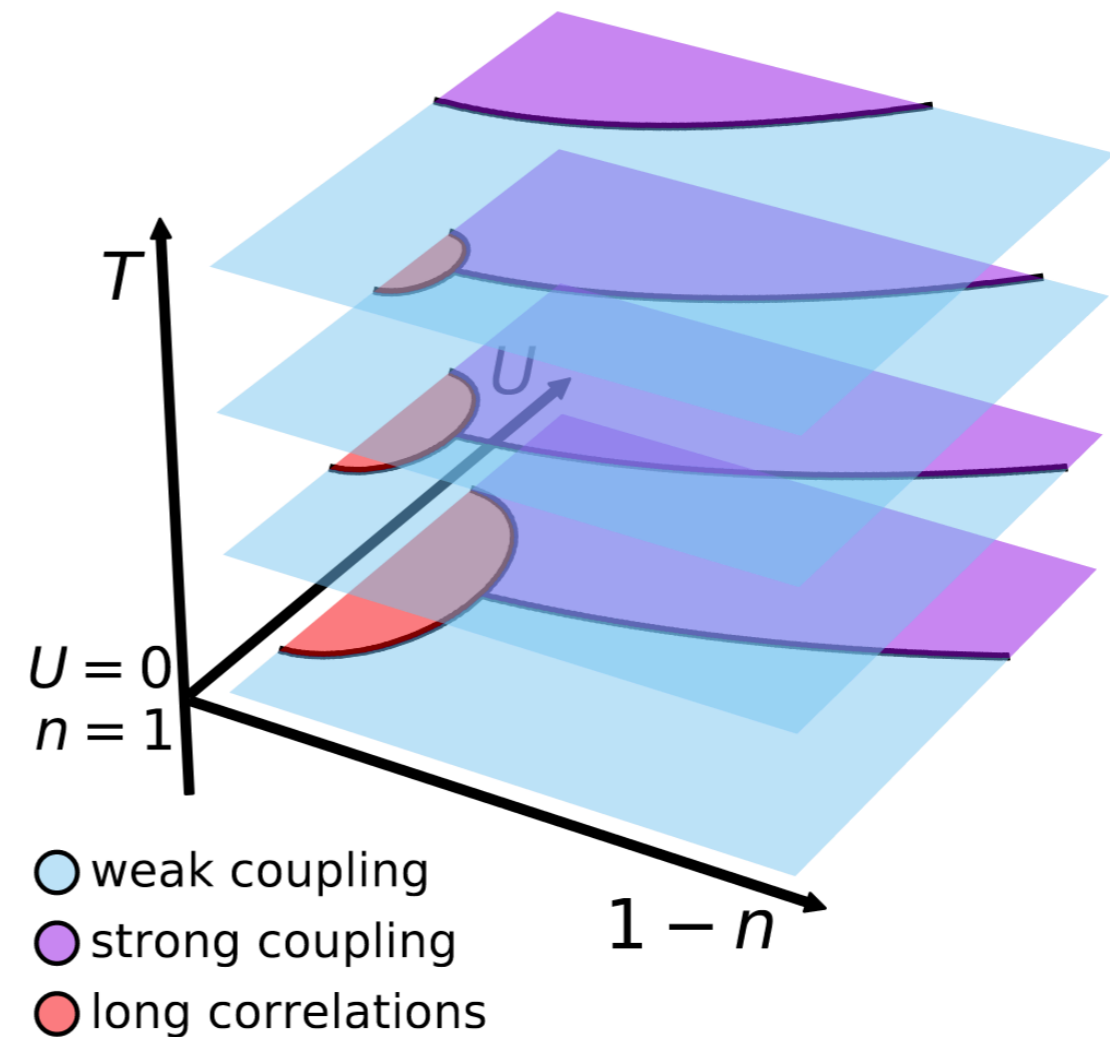
[Zheng et al, Science 2017]

The Weak, the Long, and the Strong

$U/t=4$, $T/t=0.1$, $n=0.925$

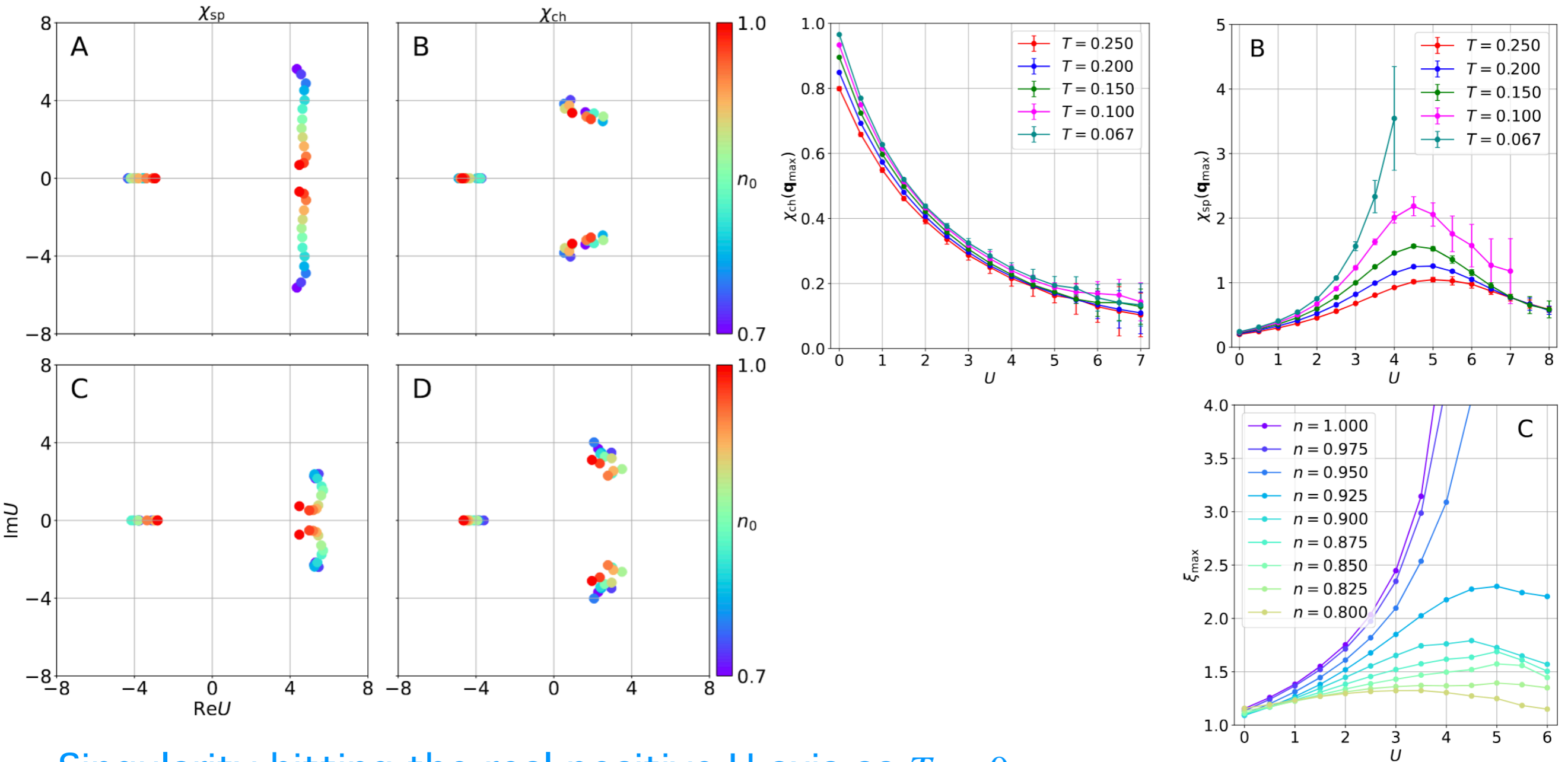


$\chi_{sp}(\mathbf{r})$



But no signature of stripe precursors!

Spin and charge decoupling from analytic structure



Singularity hitting the real positive U axis as $T \rightarrow 0$
for spin susceptibility \Rightarrow SDW crossover

No such thing happening for charge!

Origin and fate of the pseudogap phase

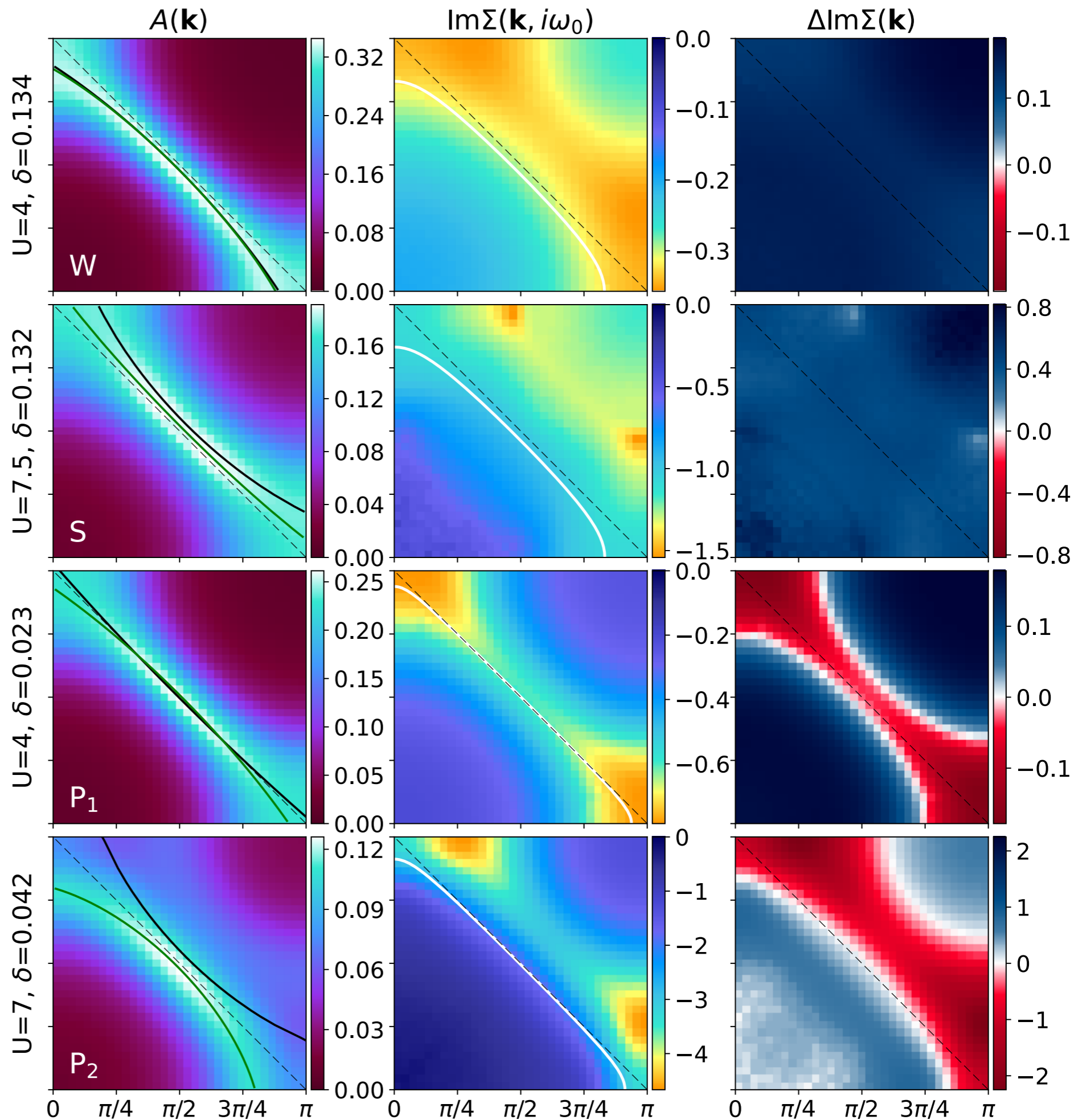
Pseudogap regimes

Pseudogap: momentum-dependent destruction of quasiparticles

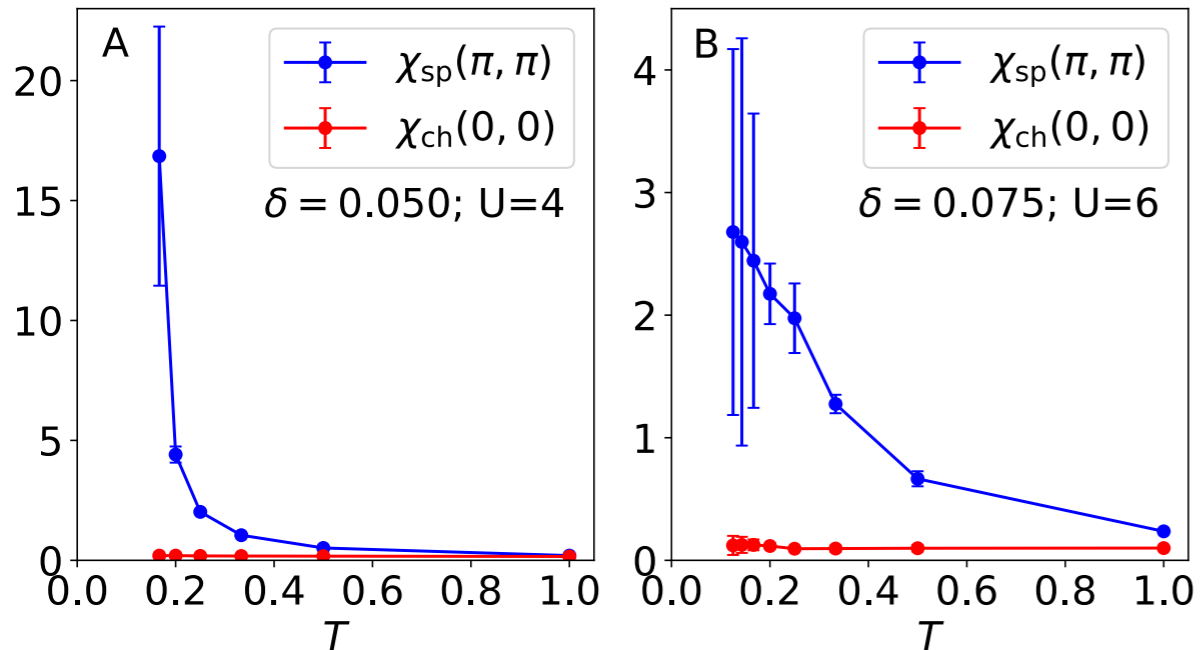
Weak-coupling pseudogap: long spin correlation length

Strong-coupling pseudogap: short-range spin correlations

[Šimkovic, RR, Georges, Ferrero, Science 2024]



Pseudogap: origin



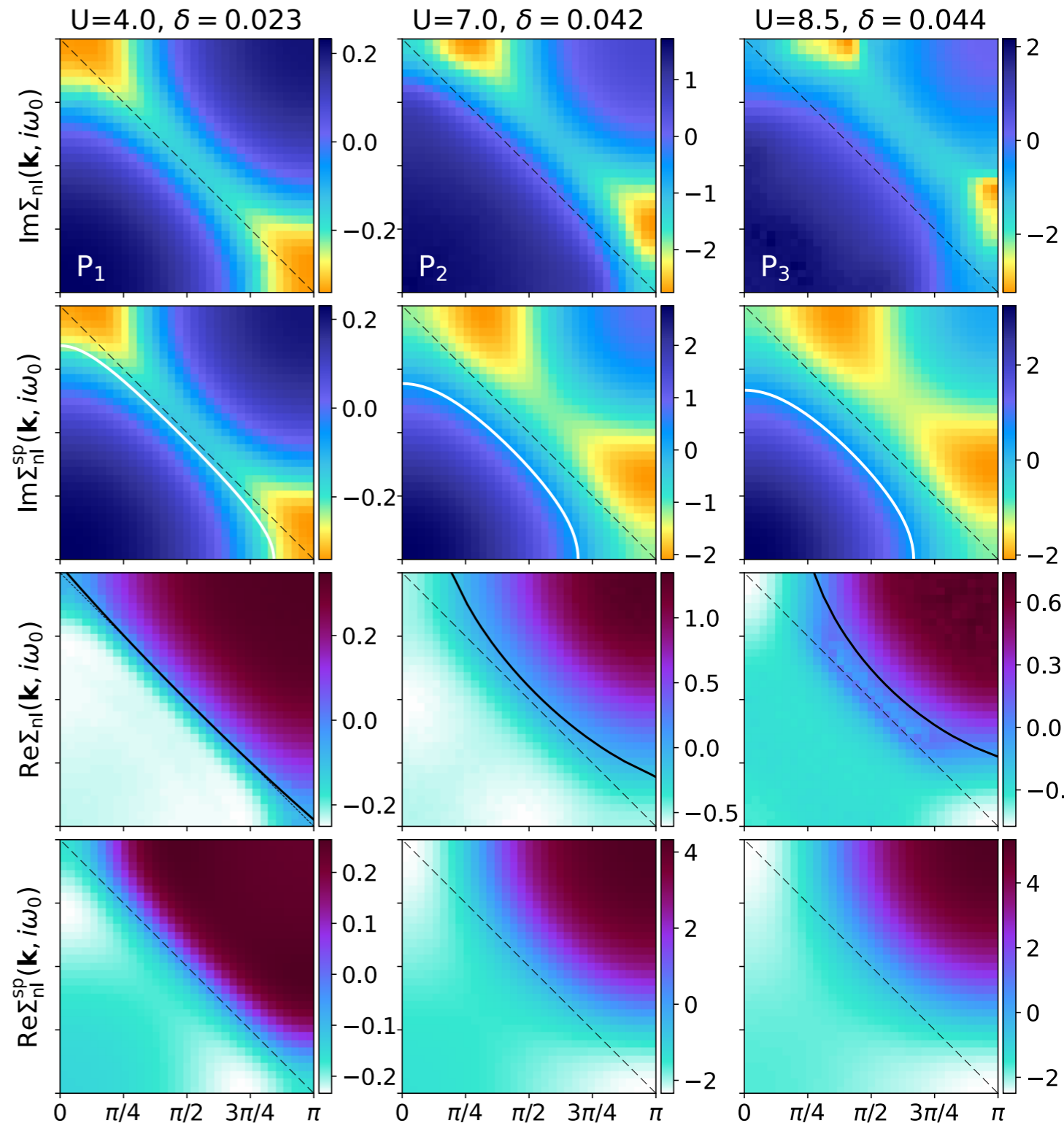
Modified Spin Fluctuation Theory

$$\Sigma = \Sigma_{\text{loc}} + \Sigma_{\text{nl}}$$

$$\Sigma = -G \star W \star \Gamma \quad W = U - U^2 \chi_{\text{sp}} \quad \Gamma = \bar{\gamma}$$

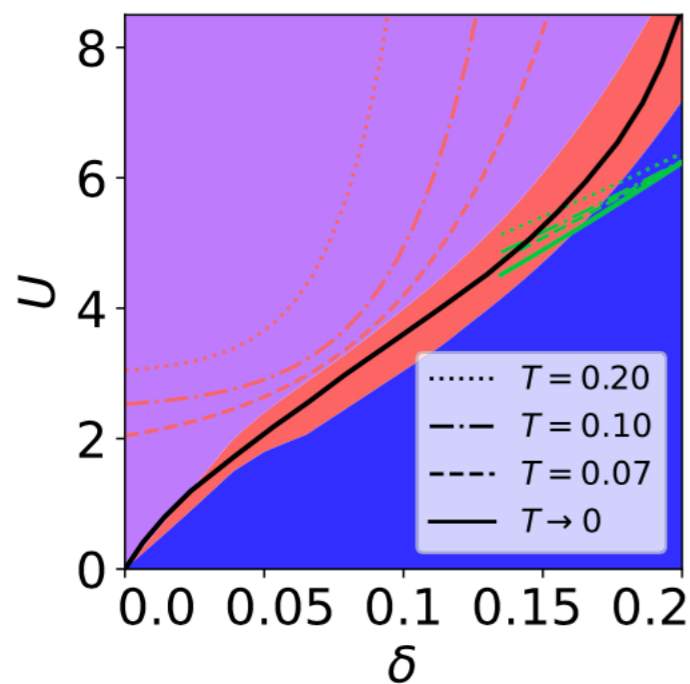
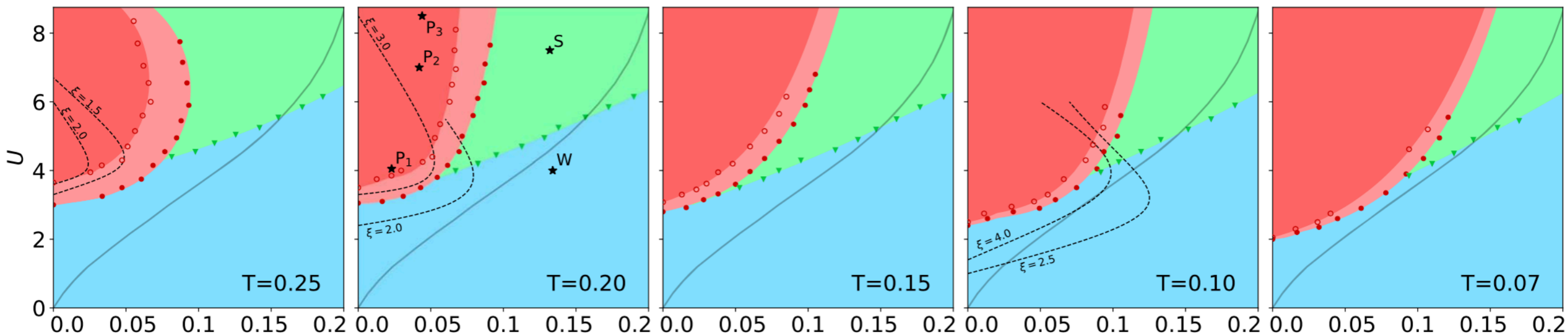
$$\Sigma_{\text{nl}}^{\text{sp}}(\mathbf{k}, i\omega_0) = \bar{\gamma} U^2 T \frac{1}{N} \sum_{\mathbf{q}} \frac{G_0(\mathbf{k} + \mathbf{q}, i\omega_0, \bar{\mu})}{|\mathbf{Q} - \mathbf{q}|^2 + \xi^{-2}}$$

Quantitative agreement not possible
in the strong-coupling pseudogap region

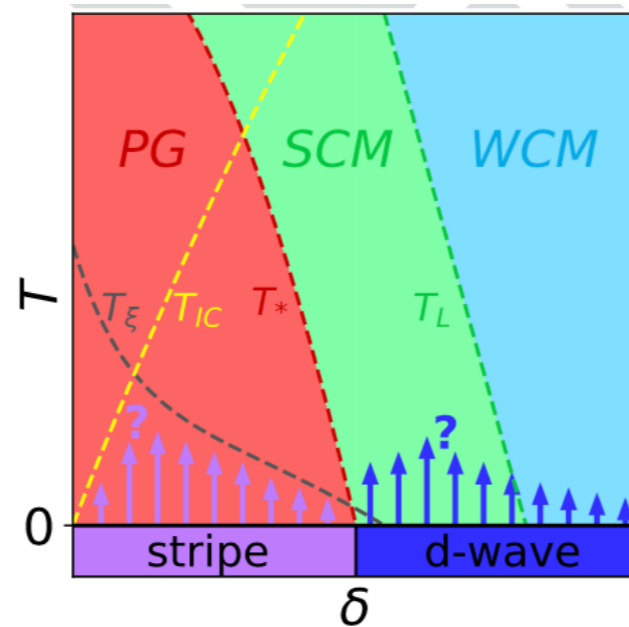


[Šimkovic, RR, Georges, Ferrero, Science 2024]

Pseudogap: Fate?



[Xu et al, PRR 2022]



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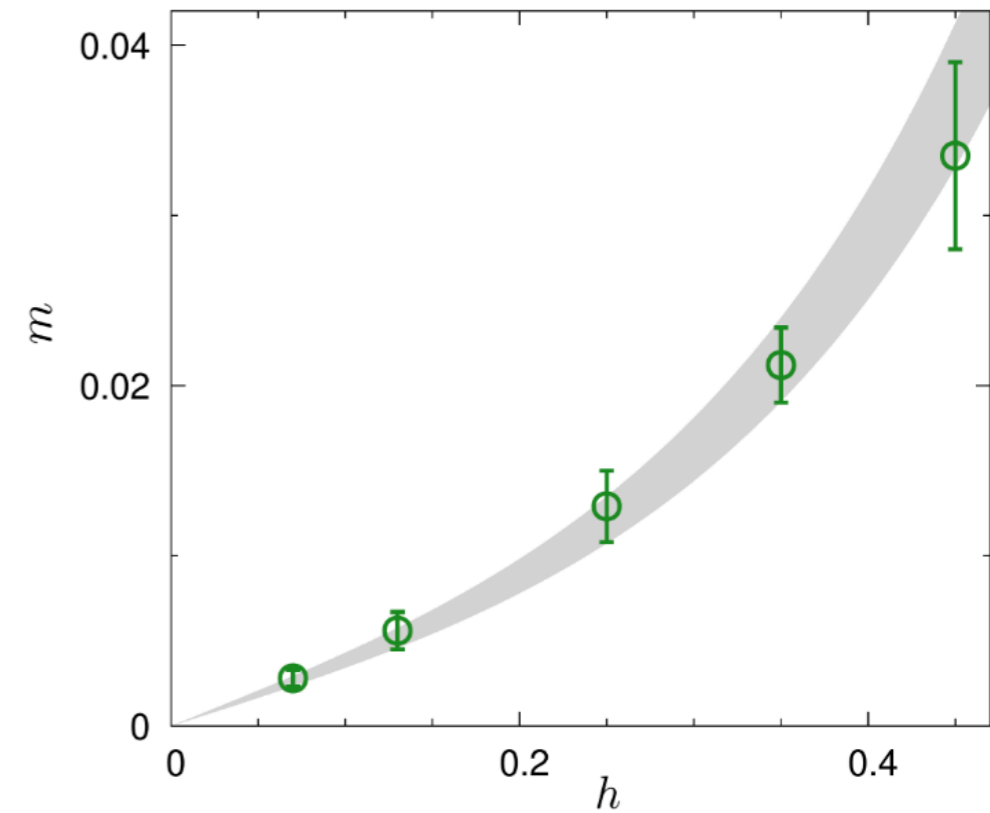
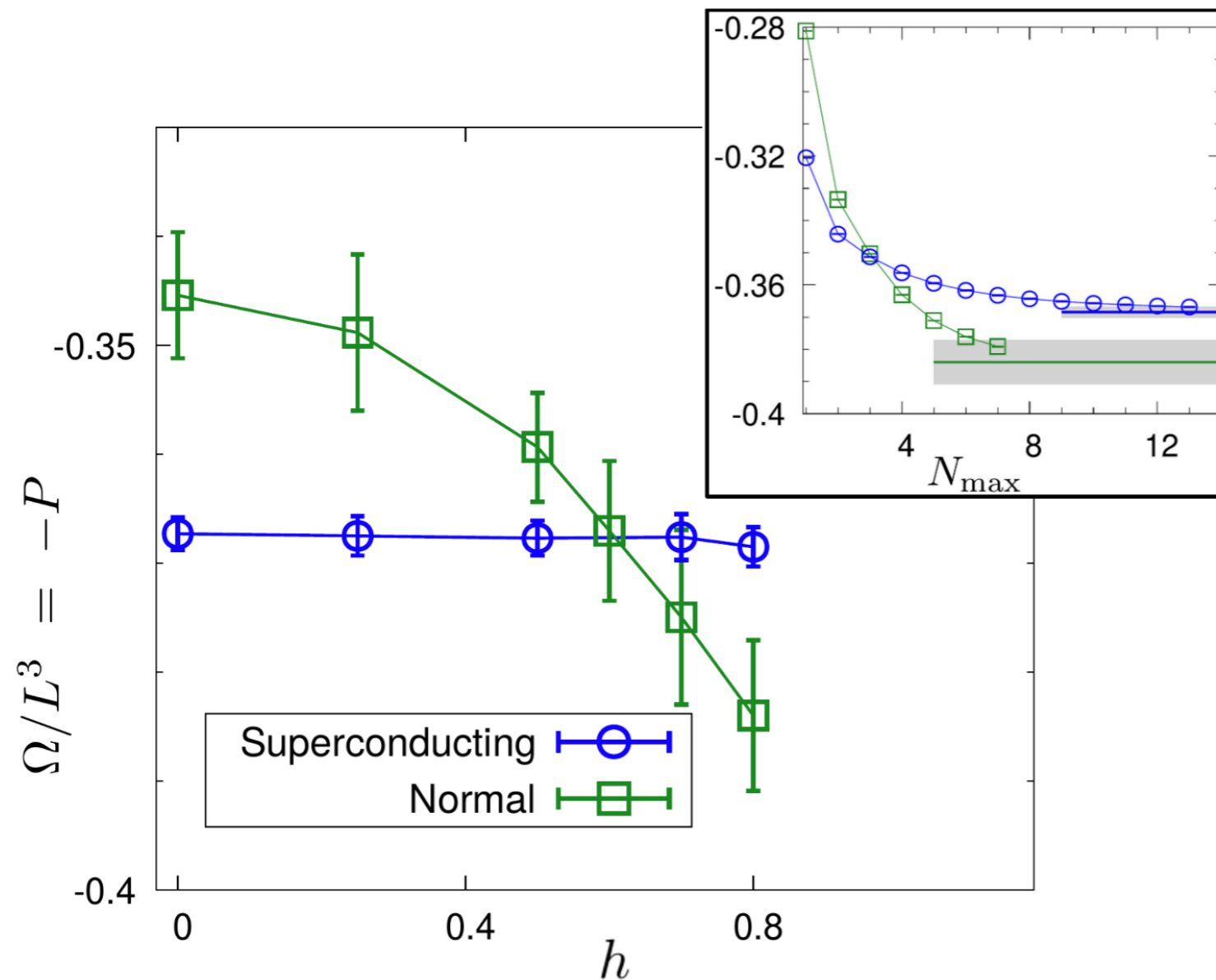
Pseudogap phase
exactly
extrapolates to
stripe ground
state!

[Šimkovic, RR, Georges, Ferrero, Science 2024]

Conclusion

- We can compute unprecedented number of Feynman diagrams, and it is useful for fermions (poly-time method)
- AFM fluctuations in the HM: Hotspots, \sim nesting
- No stripe precursor at intermediate temperature
- Origin and fate of the pseudogap in the HM

Outlook



- d-wave and stripe-order SB high-order expansions for the 2D model
- Renormalized expansions to reach lower T
- C++ code at <https://github.com/FastFeynmanDiagrammatics>, Julia code under development

[Spada, RR, Šimkovic, Garioud, Ferrero, Van Houcke, Werner, arXiv:2021]

[Garioud, Šimkovic, RR, Spada, Schäfer, Werner, Ferrero, PRL 2022]