

# Surrogate models for Quantum many-body systems

Recent Progress in Many-Body Theories(RPMBT22)@Tsukuba, Sept. 23-27, 2024



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Condensed matter physics, Quantum Chemistry, Nuclear Physics, etc.

share issues on...

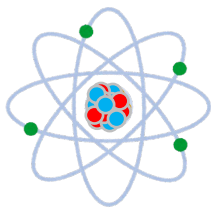
- exponential growth of the size of Hilbert space
- repeating simulations under tons of different params for quantifying uncertainties/inverse problem (e.g. nuclear force)



surrogate models, emulators, reduced order models, ...

you may call them in different ways

Quantum chemistry:



“99 > % of energy of a molecule in equilibrium

is explained within Hartree-Fock level”

(i.e., single Slater determinant)

rest 1 % is called **correlation energy**

Møller – Plesset (MP a.k.a MBPT)

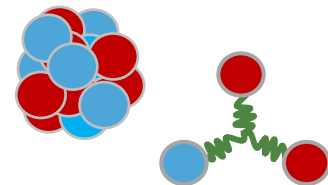
Coupled Cluster Single and Double (CCSD)

CCSD + Triple (CCSDT)

Full Configuration interaction (Full-CI)

accurate but heavy

Nuclear physics:



Interaction is highly non-perturbative & uncertain  
many channels, **three-nucleon force**,...



$^{56}\text{Ni}$  under modern Nuclear Force (Chiral EFT)

$$\text{HF} = -302.716 \text{ MeV}$$

$$\text{HF} + \text{MP2} + \text{MP3} = -473.089 \text{ MeV}$$

(MP2 = -152.533, MP3 = -17.716)

~~How dare people say perturbation theory !!~~

c.f. Energy (Exp.) = -483.996 MeV

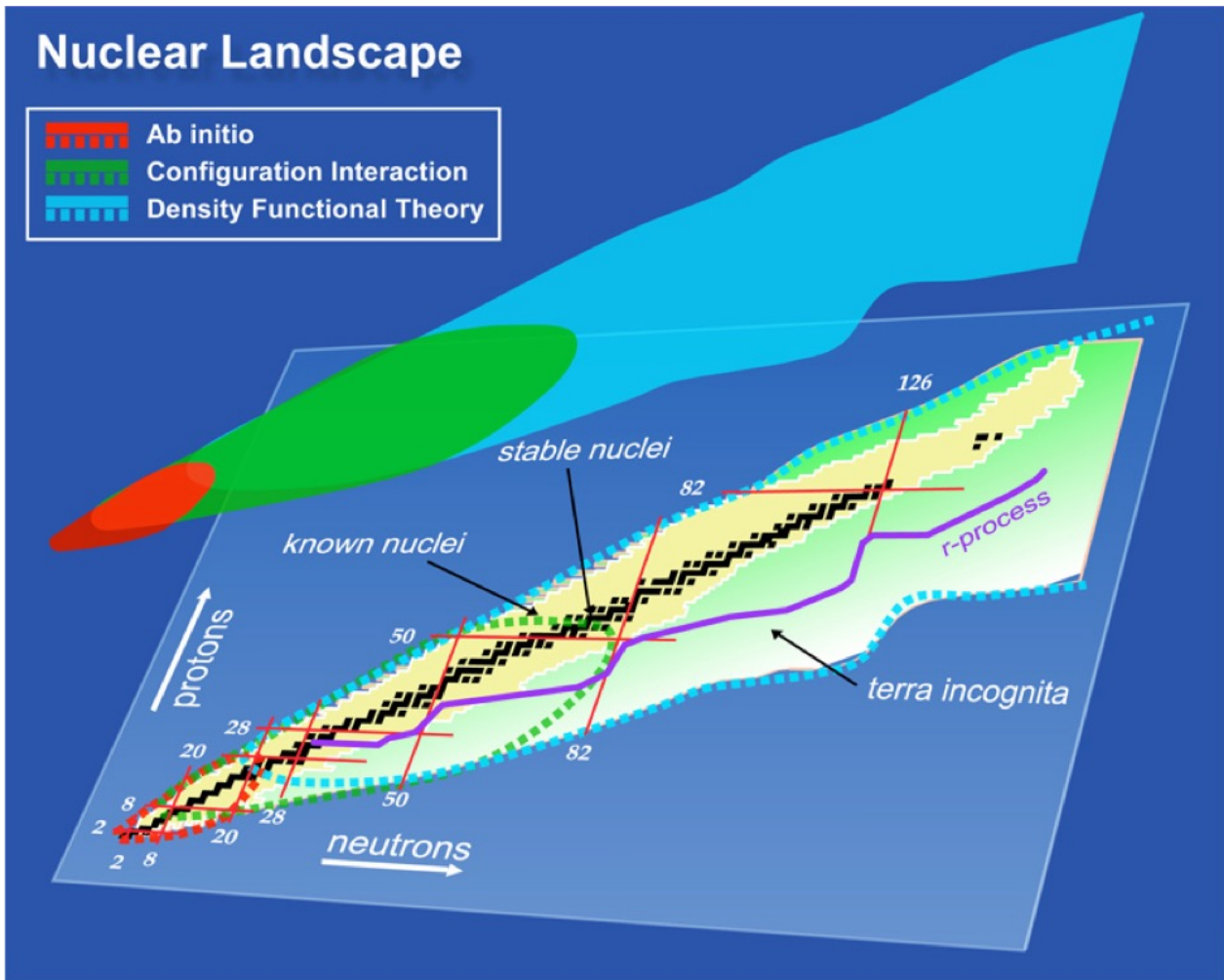


Fig. from UNEDF

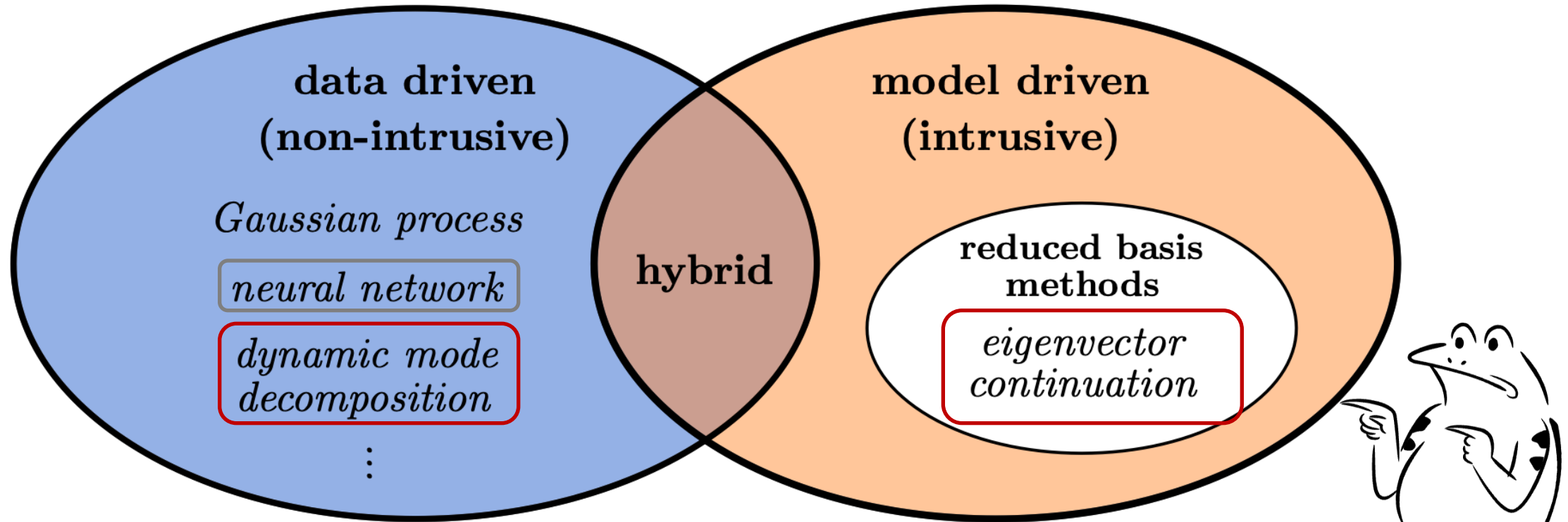
“Terra incognita” in nuclear physics

- unstable nuclei (e.g. r-process path)
- superheavy nuclei
- etc.

The scope of **Ab initio** and **CI** is gradually expanding, but still limited

The connection between **EDF** and nuclear force is unclear

## reduced order models



In 2018, a seminal paper is published: Eigenvector Continuation (EC)

PHYSICAL REVIEW LETTERS **121**, 032501 (2018)

Featured in Physics

**Eigenvector Continuation with Subspace Learning**

Dillon Frame,<sup>1,2</sup> Rongzheng He,<sup>1,2</sup> Ilse Ipsen,<sup>3</sup> Daniel Lee,<sup>4</sup> Dean Lee,<sup>1,2</sup> and Ermal Rrapaj<sup>5</sup>

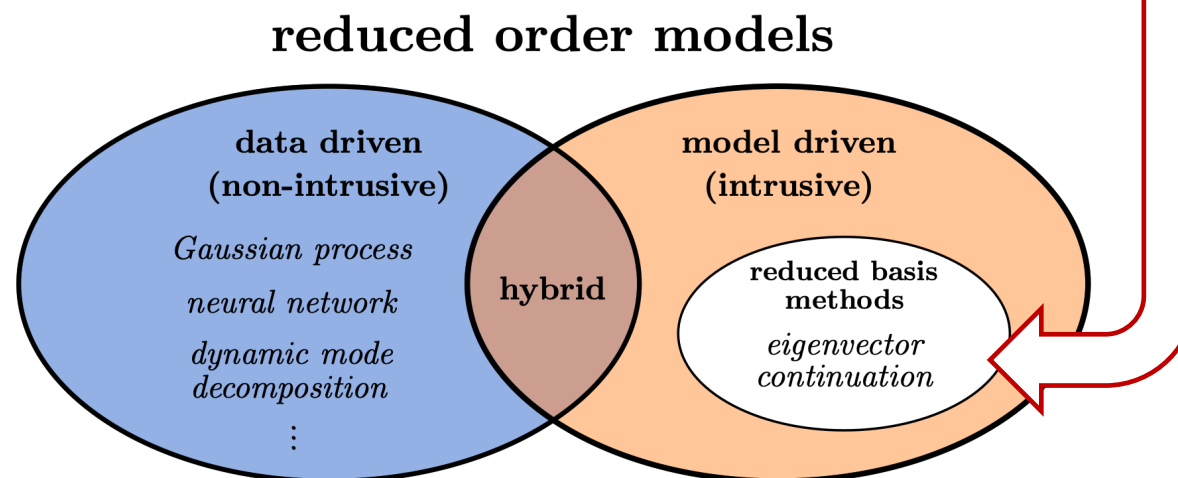
<sup>1</sup>*Facility for Rare Isotope Beams and Department of Physics and Astronomy,  
Michigan State University, East Lansing, Michigan 48824, USA*

<sup>2</sup>*Department of Physics, North Carolina State University, Raleigh, North Carolina 27695, USA*

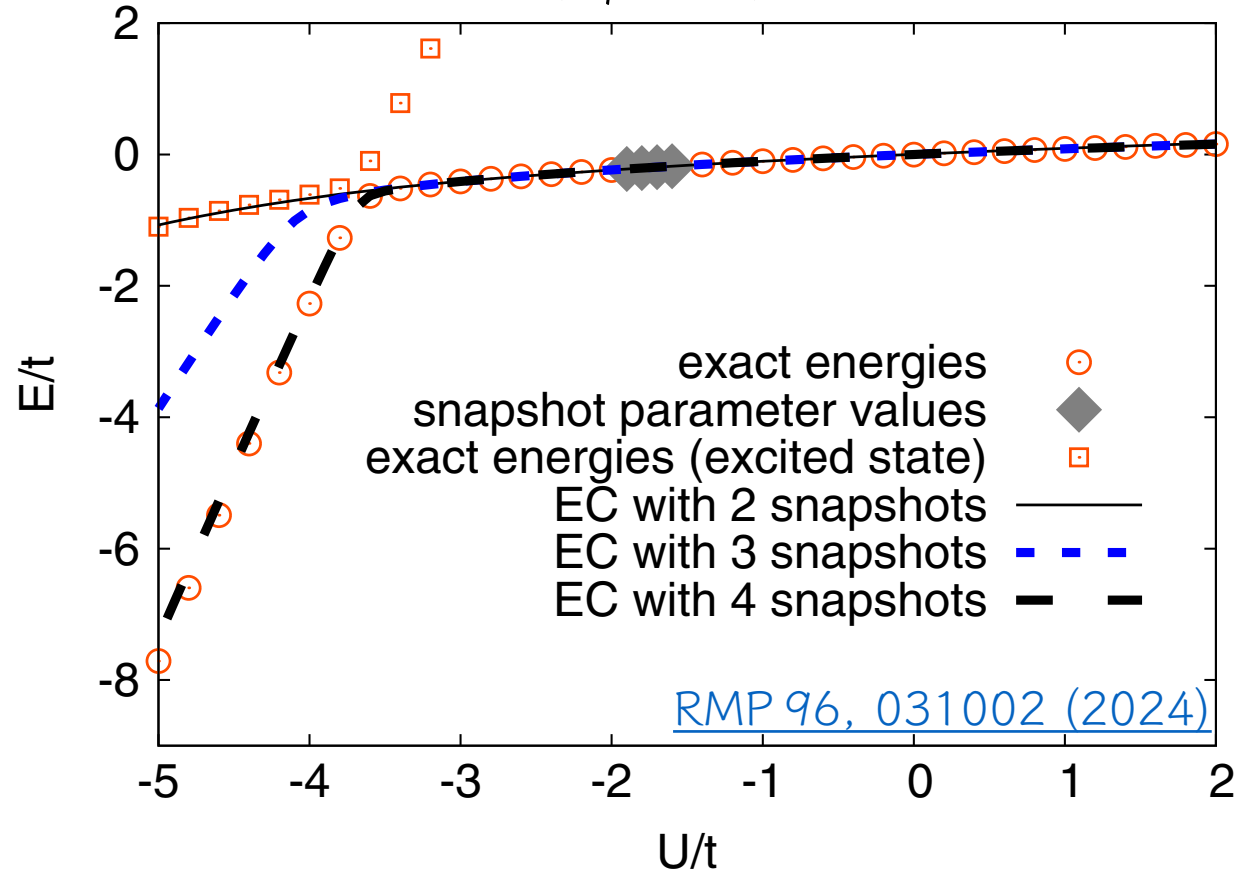
<sup>3</sup>*Department of Mathematics, North Carolina State University, Raleigh, North Carolina 27695, USA*

<sup>4</sup>*School of Engineering and Applied Science, University of Pennsylvania, Philadelphia, Pennsylvania 19104, USA*

<sup>5</sup>*Department of Physics, University of Guelph, Guelph, Ontario N1G 2W1, Canada*



3D Bose-Hubbard model  
 $4 \times 4 \times 4$ ,  $\mu = -6t$ ,  $N=4$



1. Suppose you have exact eigenstates at some points (taking **snapshots**)
2. Span the wavefunction by the samples and solve generalized eigen val. prob.

$$\tilde{H}\vec{v} = \lambda N\vec{v},$$

$$\tilde{H}_{i,j} = \langle \psi(\vec{c}_i) | H(\vec{c}_\odot) | \psi(\vec{c}_j) \rangle,$$

$$N_{i,j} = \langle \psi(\vec{c}_i) | \psi(\vec{c}_j) \rangle.$$

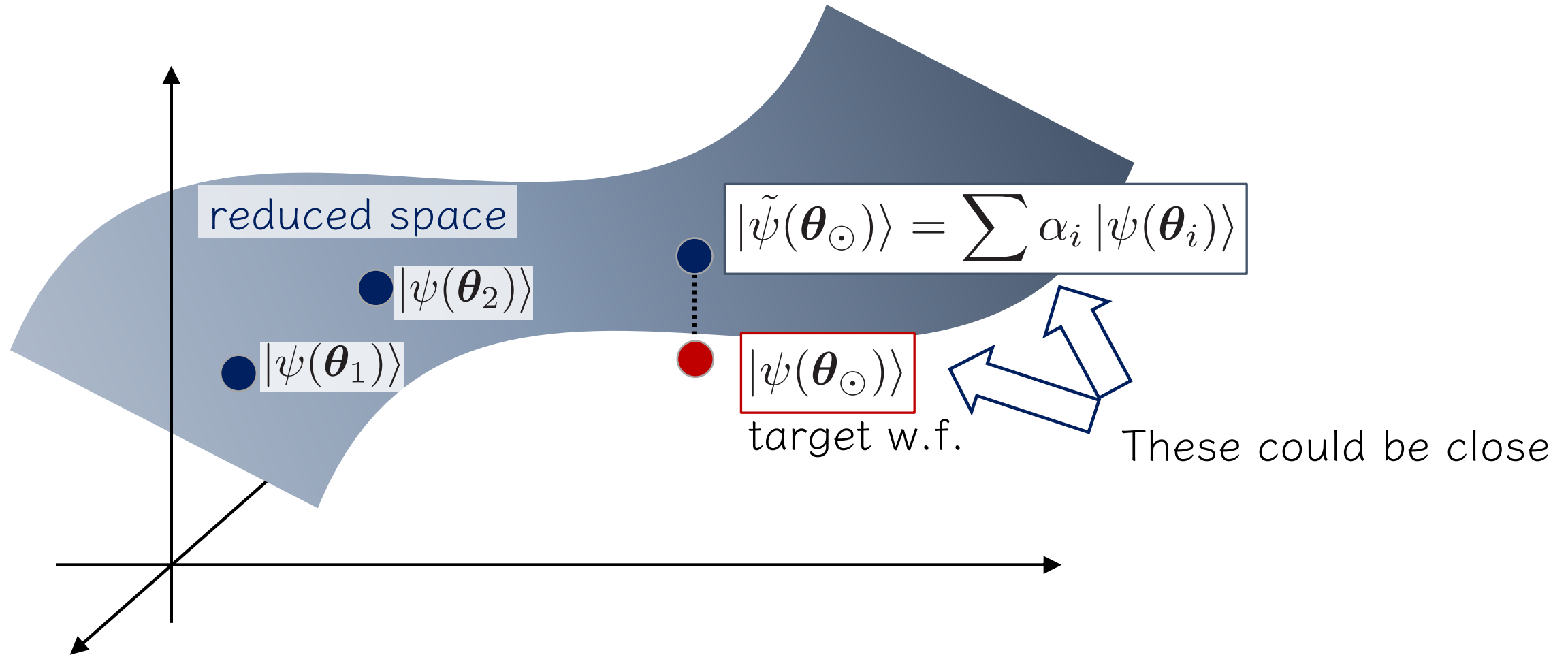
$$E(\vec{c}_\odot) \simeq \lambda, \quad |\psi(\vec{c}_\odot)\rangle \simeq \sum_{i=1}^{N_s} v_i |\psi(\vec{c}_i)\rangle \equiv |\psi_{EC}(\vec{c}_\odot)\rangle.$$

↑ several snapshots are enough to express eigenstates elsewhere

That have been proven in other quantum many-body systems

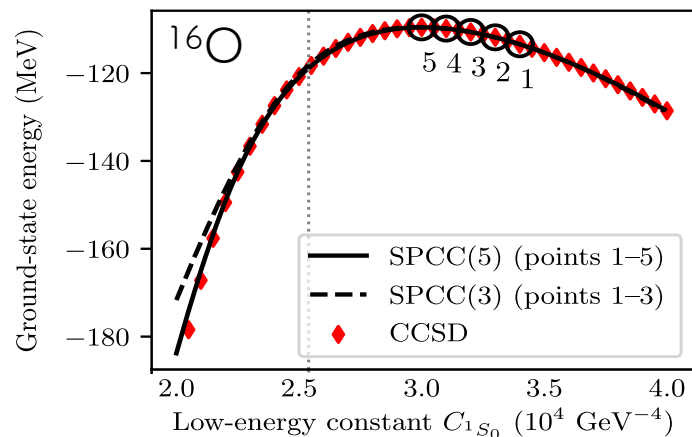
Wavefunctions/eigen states are obviously **not arbitrary** in the Hilbert space

Your snapshots and target eigenstates may live on a certain subspace



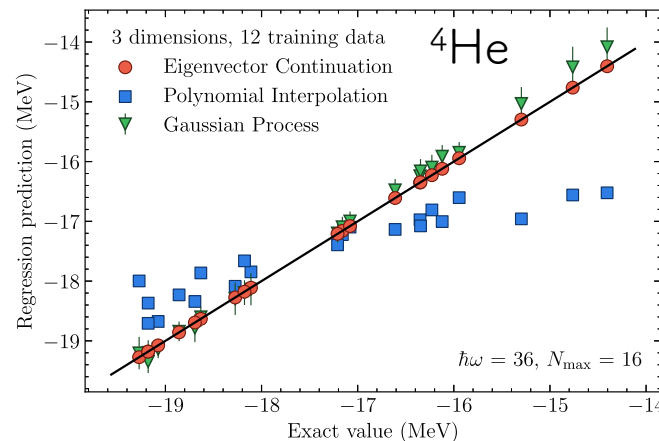


## Coupled cluster



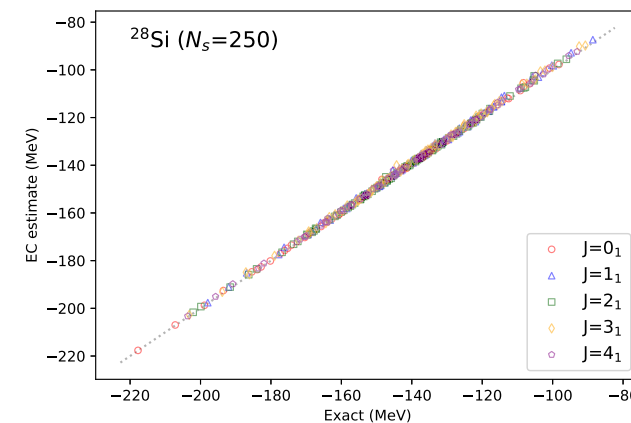
[PRL 123, 252501 \(2019\)](#)

## No-core shell model (Full-CI)



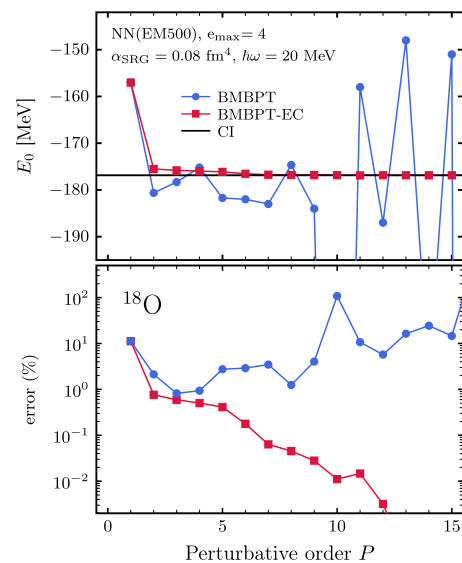
[PLB 810, 135814 \(2020\)](#)

## Valence CI



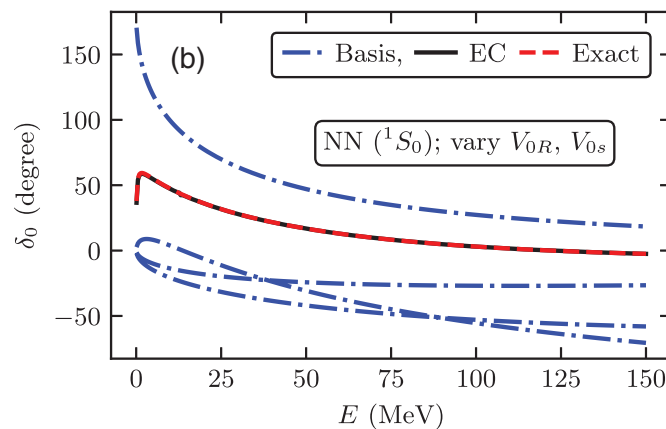
SY and N.Shimizu, [PTEP 2022 053D02 \(2022\)](#)

## Bogoliubov-MBPT



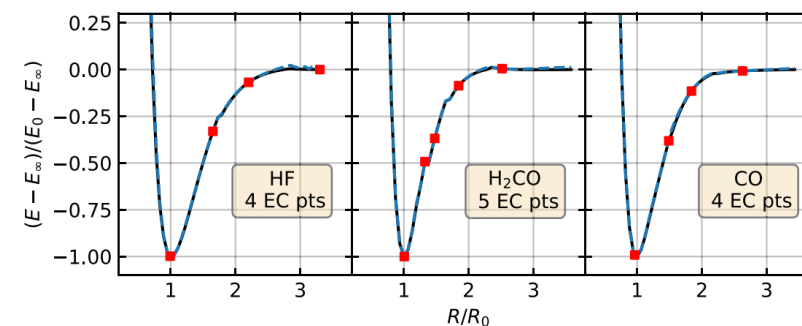
[PRC 101, 041302\(R\)\(2020\)](#)

## Scattering phase shifts



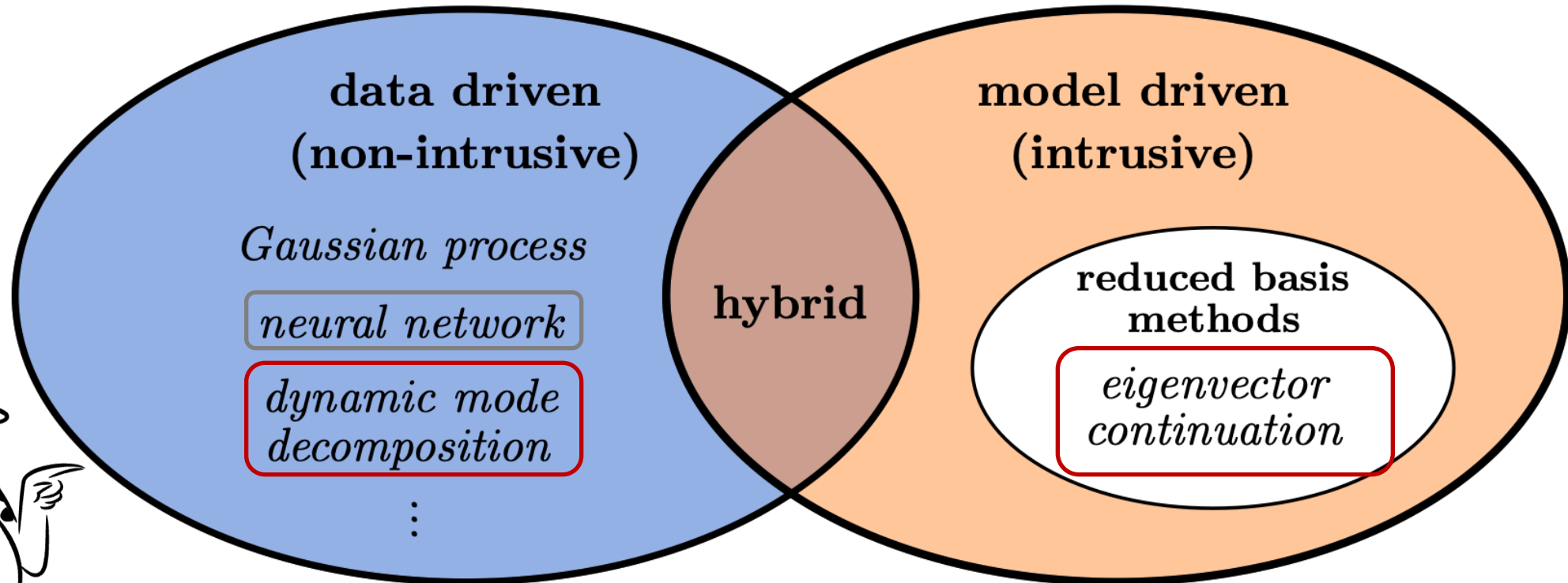
[PLB 809, 135719 \(2020\)](#)

## Chemistry



[Electron. Struct. 5 \(2023\) 045007](#)

## reduced order models



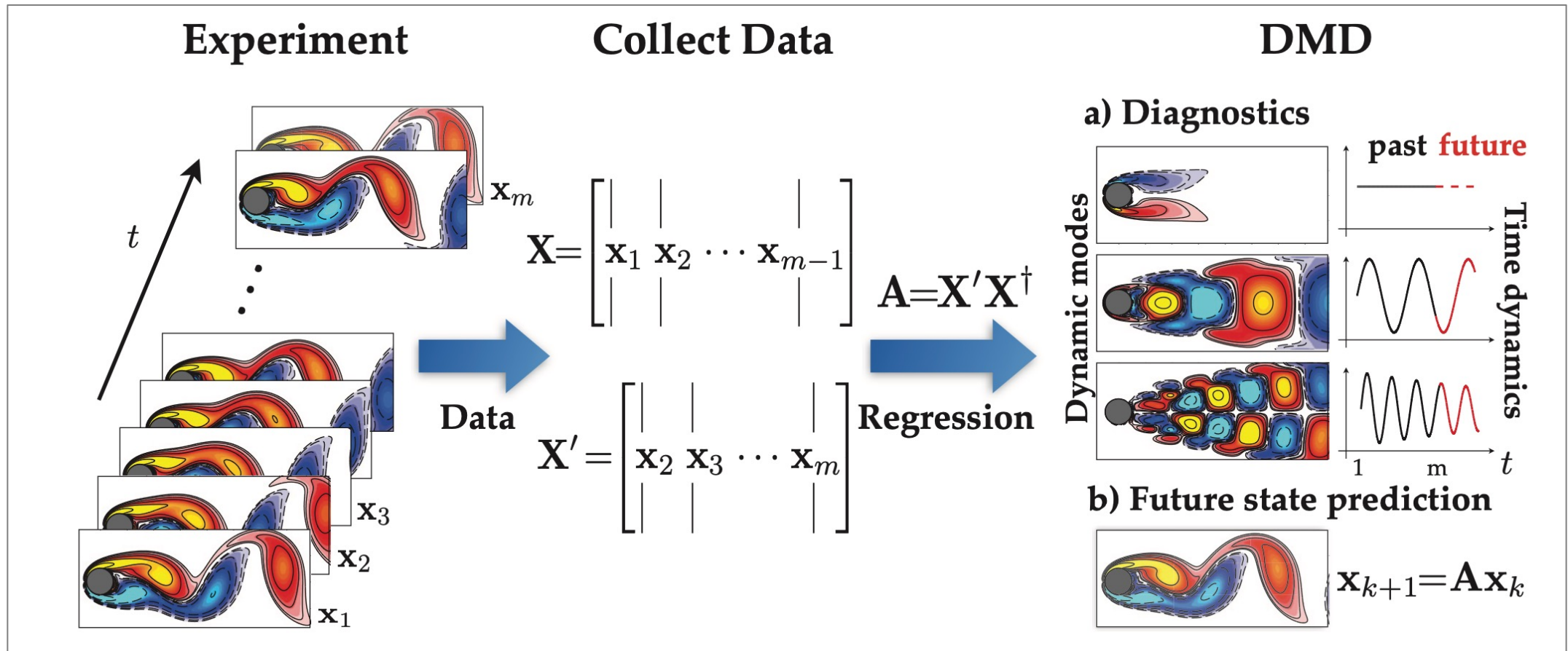


Fig 1.1 from Kutz et al., "[Dynamic Mode Decomposition](#)" SIAM

- various modes are decomposed into "dynamical modes"
- one can reconstruct original snapshots (and make predictions)

$$1. \quad \mathbf{X} \equiv \begin{pmatrix} | & & | \\ \mathbf{x}_1 & \cdots & \mathbf{x}_N \\ | & & | \end{pmatrix}^{\text{D} \times \text{N}}, \mathbf{Y} \equiv \begin{pmatrix} | & & | \\ \mathbf{x}_2 & \cdots & \mathbf{x}_{N+1} \\ | & & | \end{pmatrix}^{\text{D} \times \text{N}} \quad \mathbf{Y} = \mathbf{F}(\mathbf{X}) \Rightarrow \mathbf{Y} \approx \mathbf{A}\mathbf{X}$$

D (dimension of many-body operator)  $> 10^7$     approximating non-linear map F by linear map A  
 N (# of snapshots)  $\sim 10 - 10^3$ ?

$$2. \text{ SVD of } \mathbf{X} \quad \mathbf{X} = \mathbf{U}\mathbf{\Sigma}\mathbf{V}^* \rightarrow \text{truncated SVD } \mathbf{X} \approx \mathbf{U}_r\mathbf{\Sigma}_r\mathbf{V}_r^\dagger$$

3. Compute the matrix A using Moore-Penrose pseudo-inverse of X

$$\mathbf{A} \approx \mathbf{Y}\mathbf{X}^+ = \mathbf{Y} (\mathbf{V}_r\mathbf{\Sigma}_r^{-1}\mathbf{U}_r^\dagger)$$

4. Obtain the time evolution linear map in a latent space

$$\tilde{\mathbf{A}} = \mathbf{U}_r^\dagger \mathbf{A} \mathbf{U}_r \approx \mathbf{U}_r^\dagger \mathbf{Y} \mathbf{V}_r \mathbf{\Sigma}_r^{-1}$$

$$\mathbf{A} = \underbrace{\mathbf{U}_r}_{\text{encoder}} \tilde{\mathbf{A}} \underbrace{\mathbf{U}_r^\dagger}_{\text{decoder}}$$

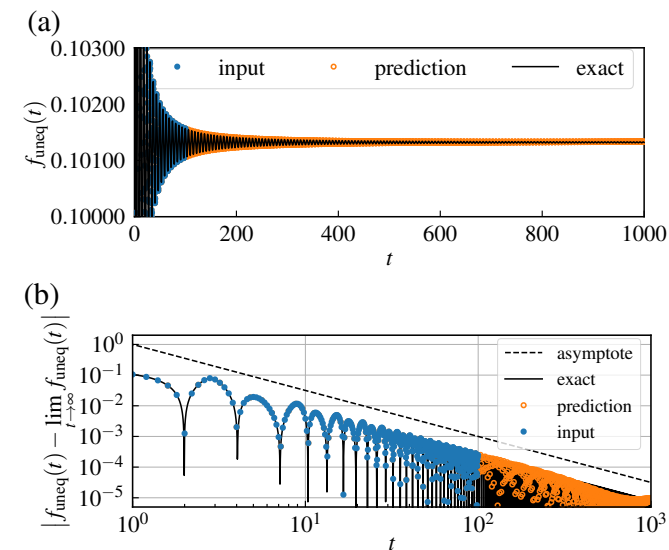
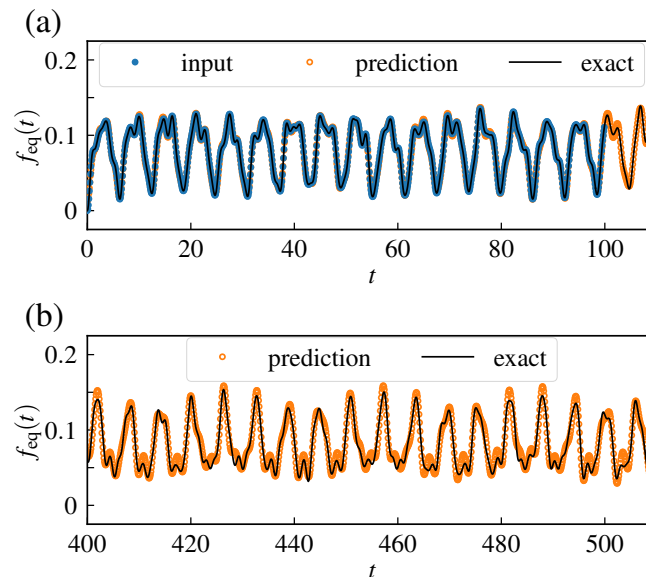
k-time step forward can be done in the latent space

## Spin system

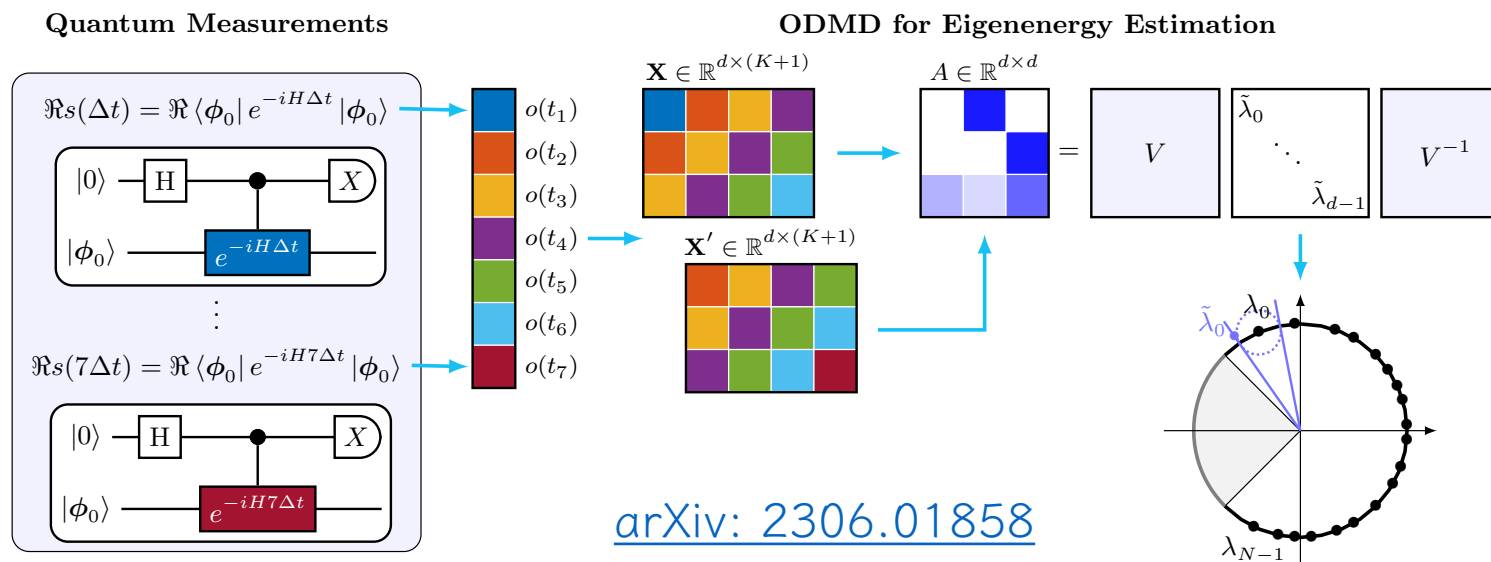
transverse-field Ising model

$$H = -J \sum_{\langle i,j \rangle} S_i^z S_j^z - \Gamma \sum_i S_i^x,$$

[arXiv: 2403.19947](https://arxiv.org/abs/2403.19947)



## Quantum computing



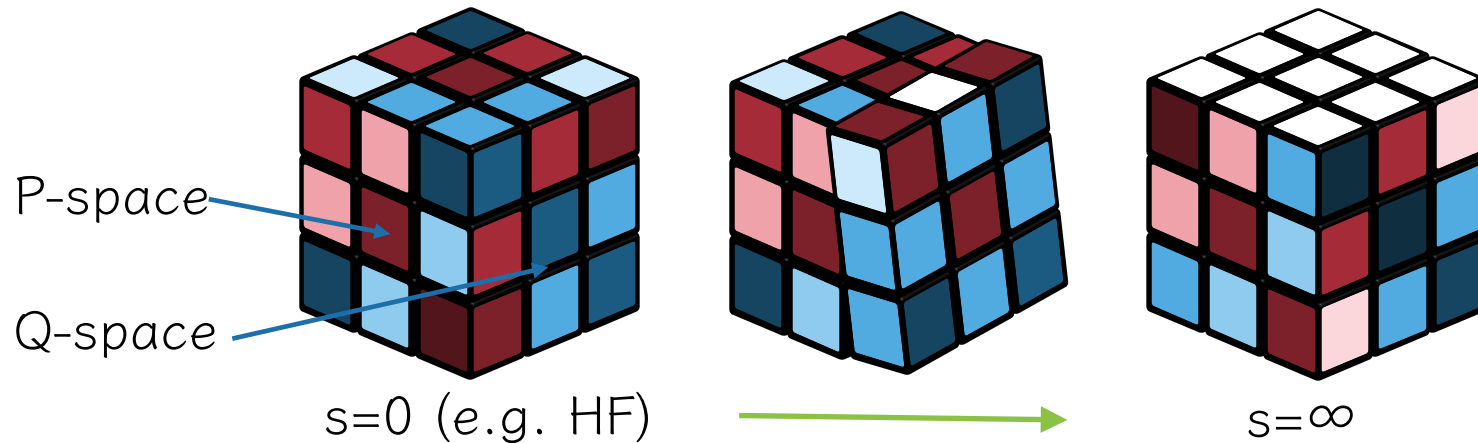
[arXiv: 2306.01858](https://arxiv.org/abs/2306.01858)

## IMSRGflow:

K. Tsukiyama, S. K. Bogner, and A. Schwenk, [PRL 106, 222502 \(2011\)](#), [PRC 85, 061304 \(2012\)](#).  
 S.R.Stroberg et al., [Annu. Rev. Nucl. Part. Sci. 2019. 69:307–362 \(2019\)](#)  
 T.D.Morris et al., [PRC 92, 034331 \(2015\)](#)

$$\frac{dH(s)}{ds} = [\eta(s), H(s)]$$

P-Q coupling



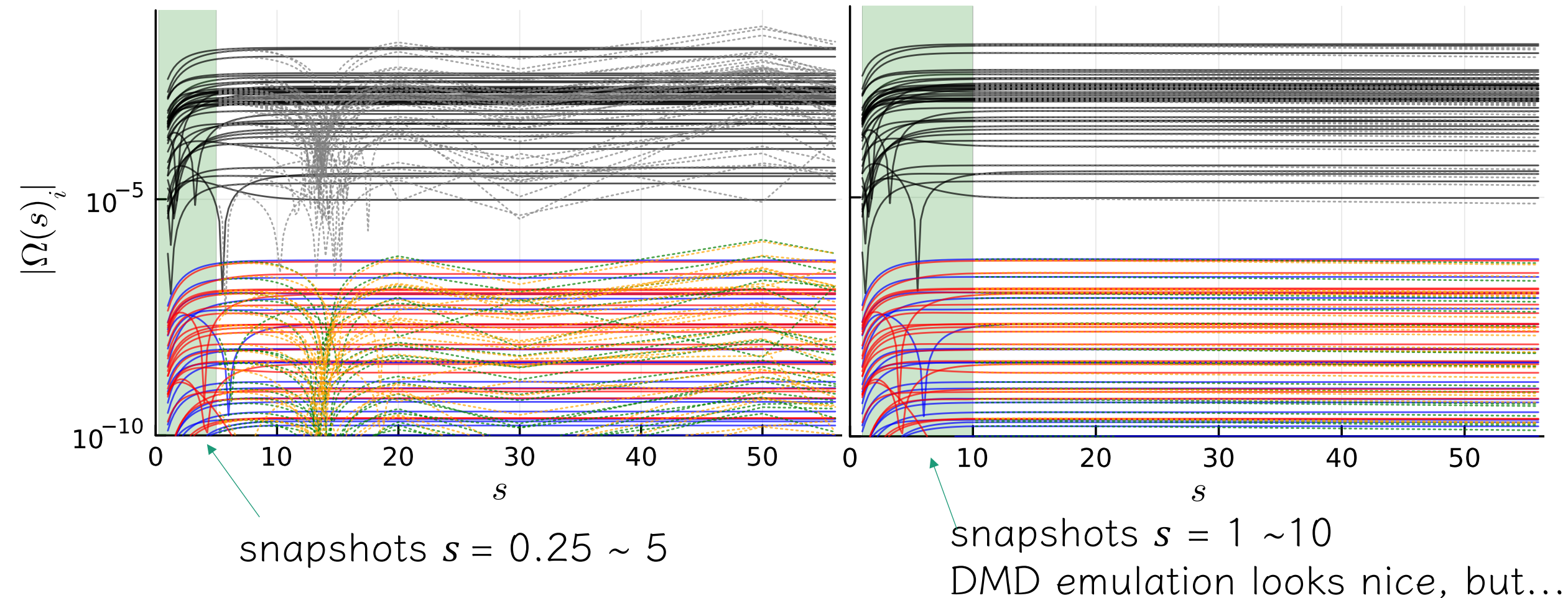
Magnus formulation  $U(s) = e^{\Omega(s)}$

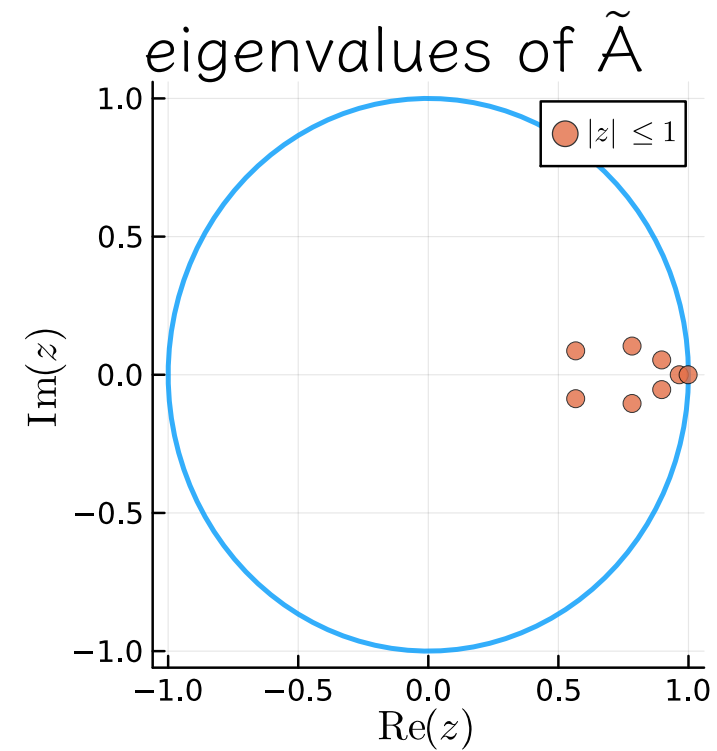
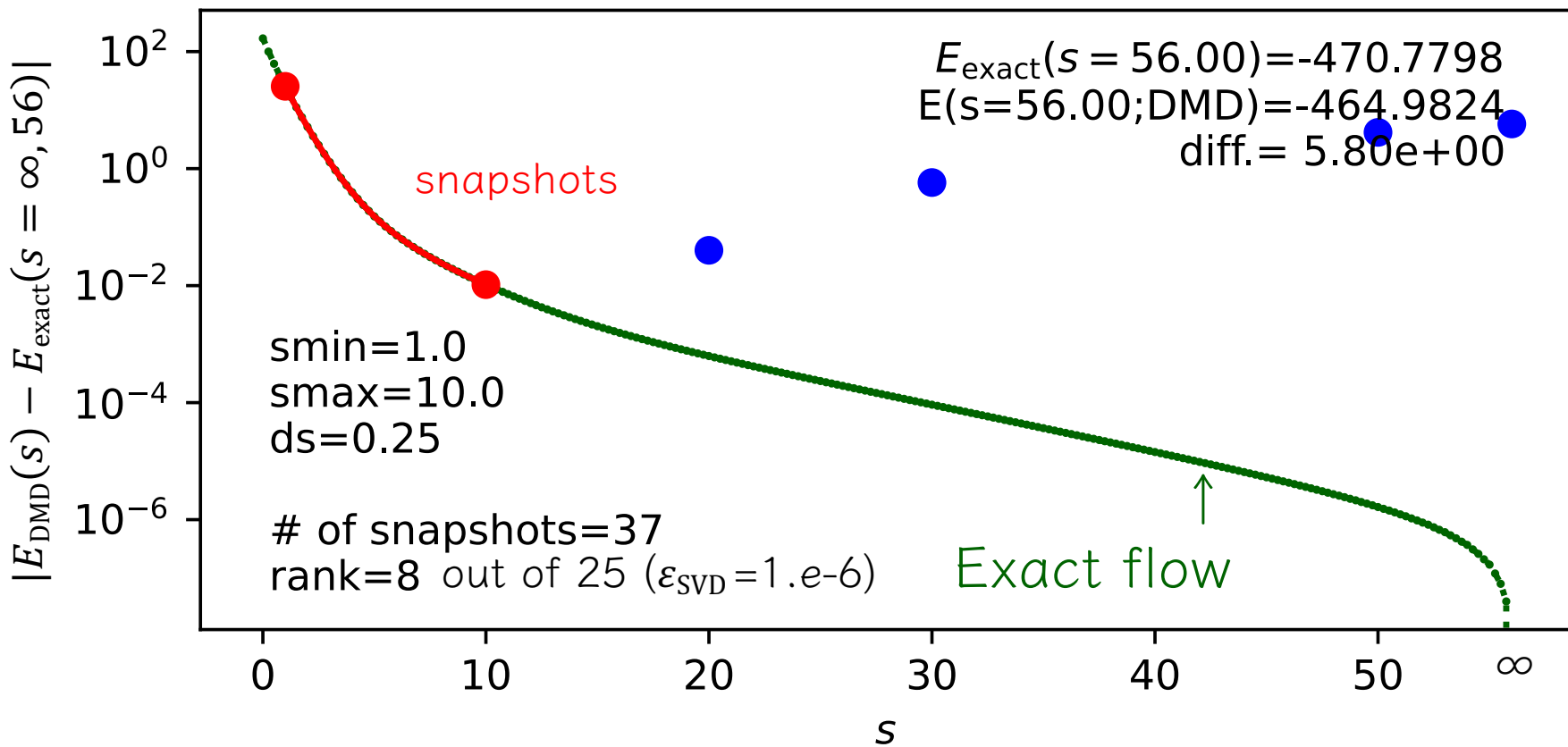
$$O(s) = e^{\Omega(s)} O(0) e^{-\Omega(s)}$$

(P, Q) = (hole, particle), (valence, others), etc.

showing 40 for each (1b, 2bpp, 2bpn)  
out of  $\sim 10^8$  ( $e_{\text{max}}=12$ ) elements

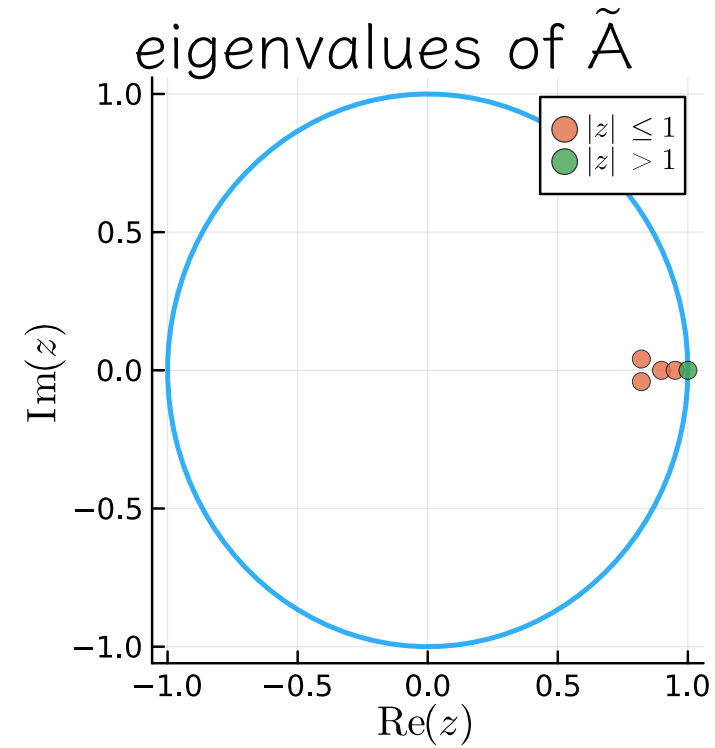
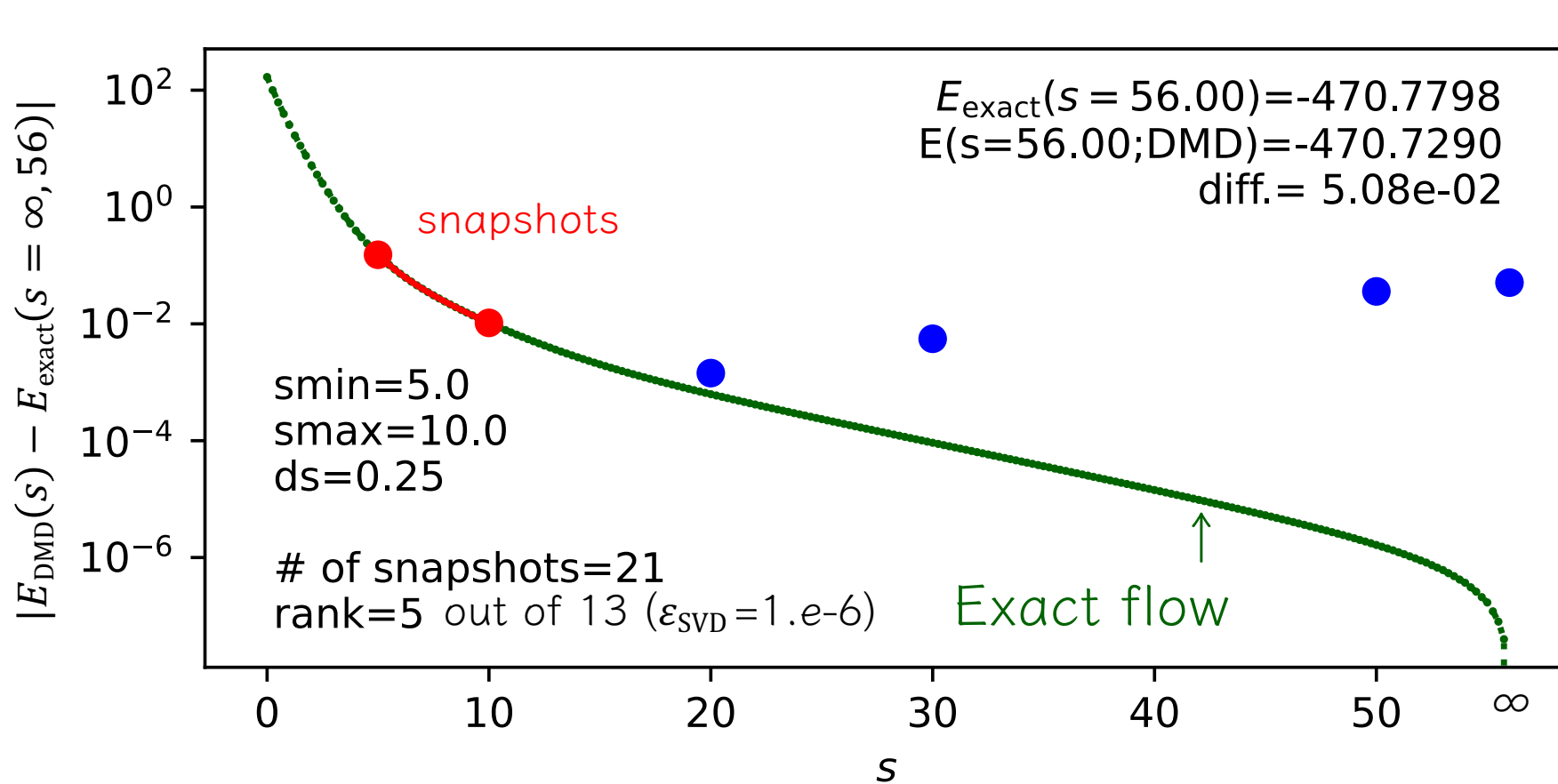
— 1b      — 2bpp      — 2bpn  
 ..... 1b(DMD)      ..... 2bpp(DMD)      ..... 2bpn(DMD)





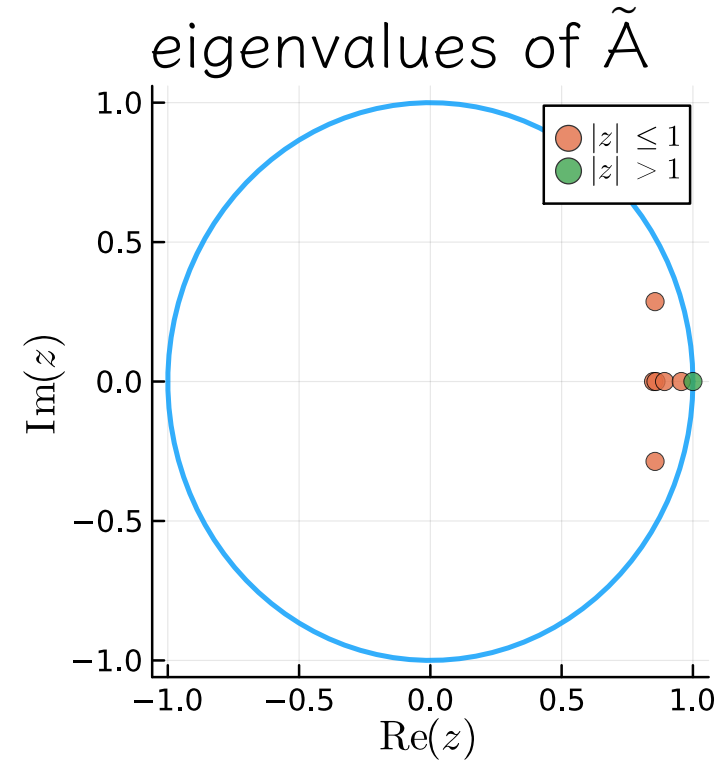
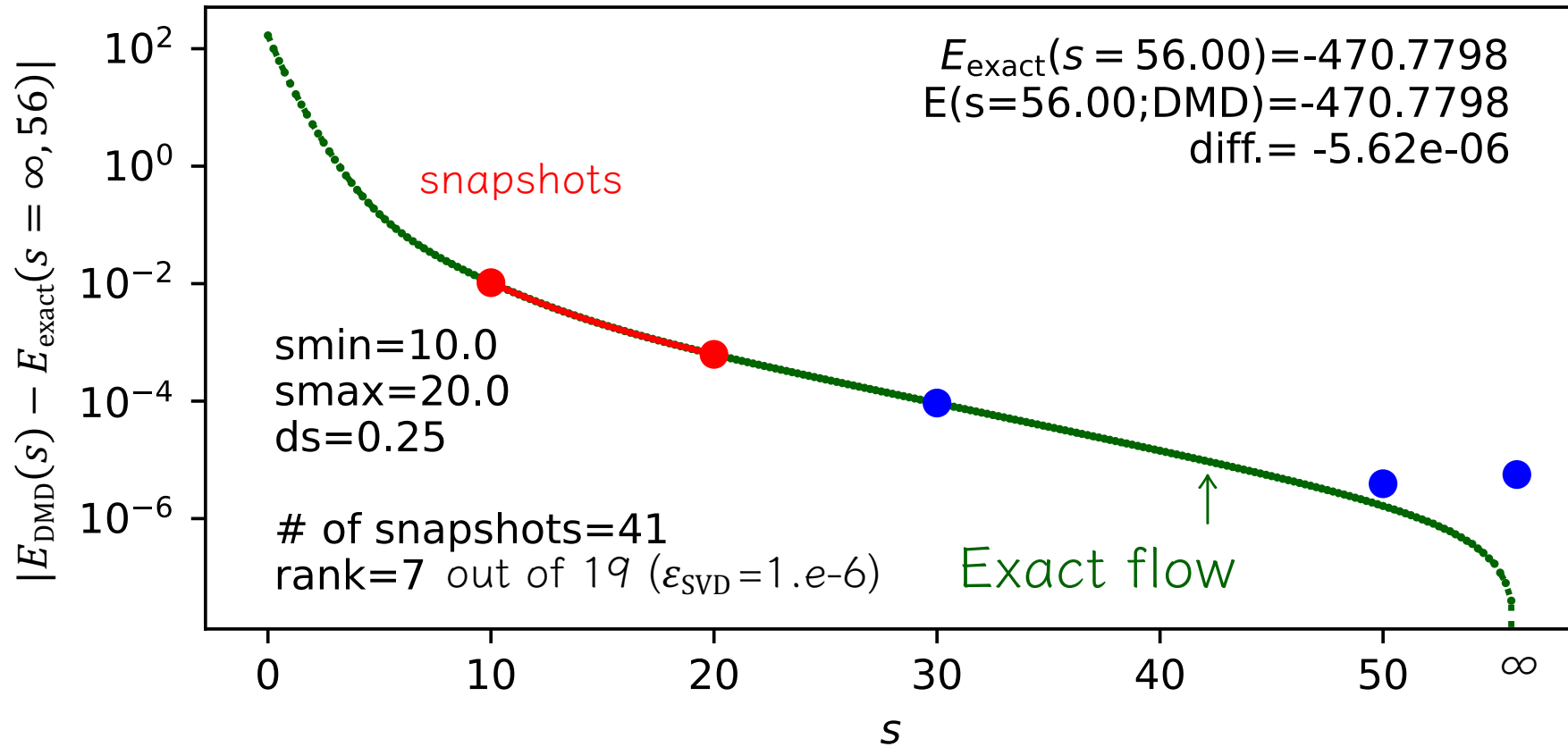
DMD predictions (blue circles ●) are obviously wrong





slightly improved but ...

→ non-linear nature of IM-SRG seen in smaller  $s$  region  
 may not be well captured by DMD (single linear operator)



we got even better results, but...

it is nontrivial to take snapshots from where to where

## PROS:

low-rank nature, cost on emulator is negligible  $\rightarrow$   $\times$  3~5 speedup



DMD can learn (at least) “linear” part of quantum dynamics

## CONS:



accumulation of errors from imperfect linear map

no a priori knowledge where to stop high-fidelity calculations

If an emulator is fast enough, you can make a diagnostic on emulator by e.g., looking at credible intervals of emulator predictions

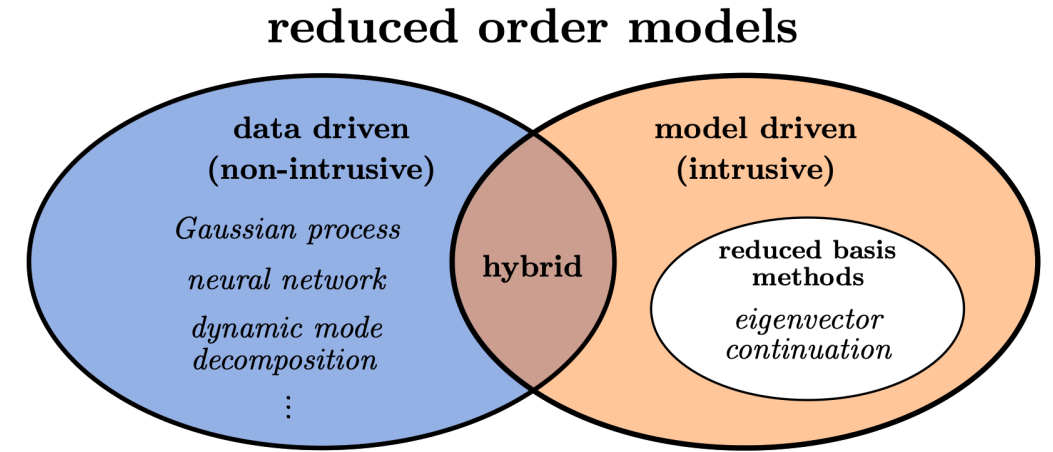
c.f. ML approach, IMSRG-Net(PINNs)  
SY [PhysRevC.108.044303 \(2023\)](https://arxiv.org/abs/2304.04430)

# Summary

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Surrogate models/Emulators matter!!

- Eigenvector continuation
- Dynamic Mode Decomposition



T. Duguet et al., [Rev. Mod. Phys. 96, 031002 \(2024\)](#)

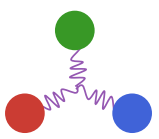
## Emulators for IM-SRG

DMD & Neural Network (For the latter, see appendix)

Both work fine for “linear” part of the IM-SRG flow,  
They are complementary. Next step will be to learn non-linear part

Related talks on emulators in nuclear physics :

Hinohara-san@Tsukuba, Xilin Zhang@MSU/FRIB



v.0.4.2 (Mar. 2024)

- ChiEFTint ~ 9,000 lines

- ◆ NN potential, Entem-Machleidt(N3LO), EMN(N4LO)
- ◆ SRG in momentum space (NN-only)
- ◆ read 3NFs by NuHamil code (T. Miyagi@Tsukuba)
- ◆ etc.

- HartreeFock ~ 3,500 lines

- ◆ spherical HF
- ◆ Møller-Plesset (a.k.a MBPT)
- ◆ Normal ordering (including TNO/ENO)

- IM-SRG ~ 3,000 lines

IMSRG & VS-IMSRG (only scalar ones for now)

emulator for Magnus-IMSRG(2) with Dynamic Mode Decomposition (DMD)

-  ShellModel.jl ~ 8,700 lines

eigenvector continuation SY and N.Shimizu [PTEP 2022 053D02](https://arxiv.org/abs/2205.05302)

# EC: shell model emulator

SY and N.Shimizu, [PTEP 2022 053D02 \(2022\)](#)

Example:

sd-shell ( $^{16}\text{O}$  core + 0d5/0d3/1s1 valence orbits)

parameters: 66 (3 SPEs & 63 TBMEs, w/ isospin)

target nuclei:  $^{25}\text{Mg}$  (vp=4,vn=5),  $^{28}\text{Si}$  (vp=vn=6, dim.  $\sim 90,000$ )

sampling 5 states for given total J at 50 (random) different points ( $5 \times 50 = 250$  samples) around USDB

$$\tilde{H}\vec{v} = \lambda N\vec{v},$$

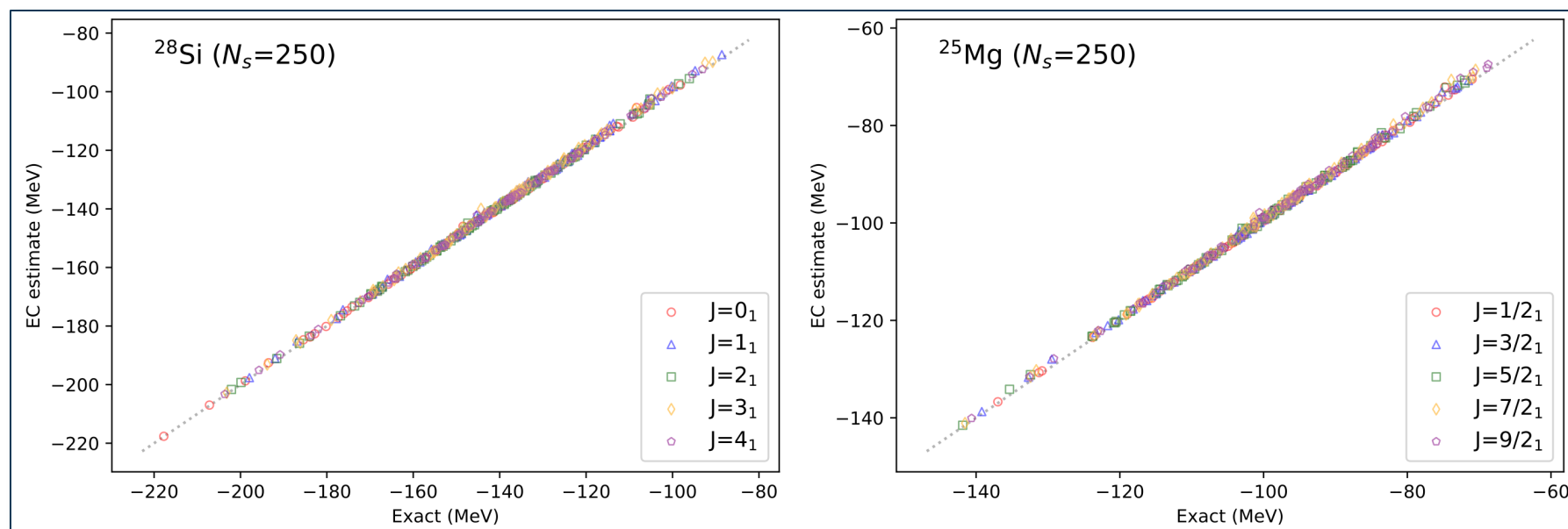
$$\tilde{H}_{i,j} = \langle \psi(\vec{c}_i) | H(\vec{c}_0) | \psi(\vec{c}_j) \rangle,$$

$$N_{i,j} = \langle \psi(\vec{c}_i) | \psi(\vec{c}_j) \rangle.$$

$$E(\vec{c}_0) \simeq \lambda,$$

$$|\psi(\vec{c}_0)\rangle \simeq \sum_{i=1}^{N_s} v_i |\psi(\vec{c}_i)\rangle \equiv |\psi_{EC}(\vec{c}_0)\rangle.$$

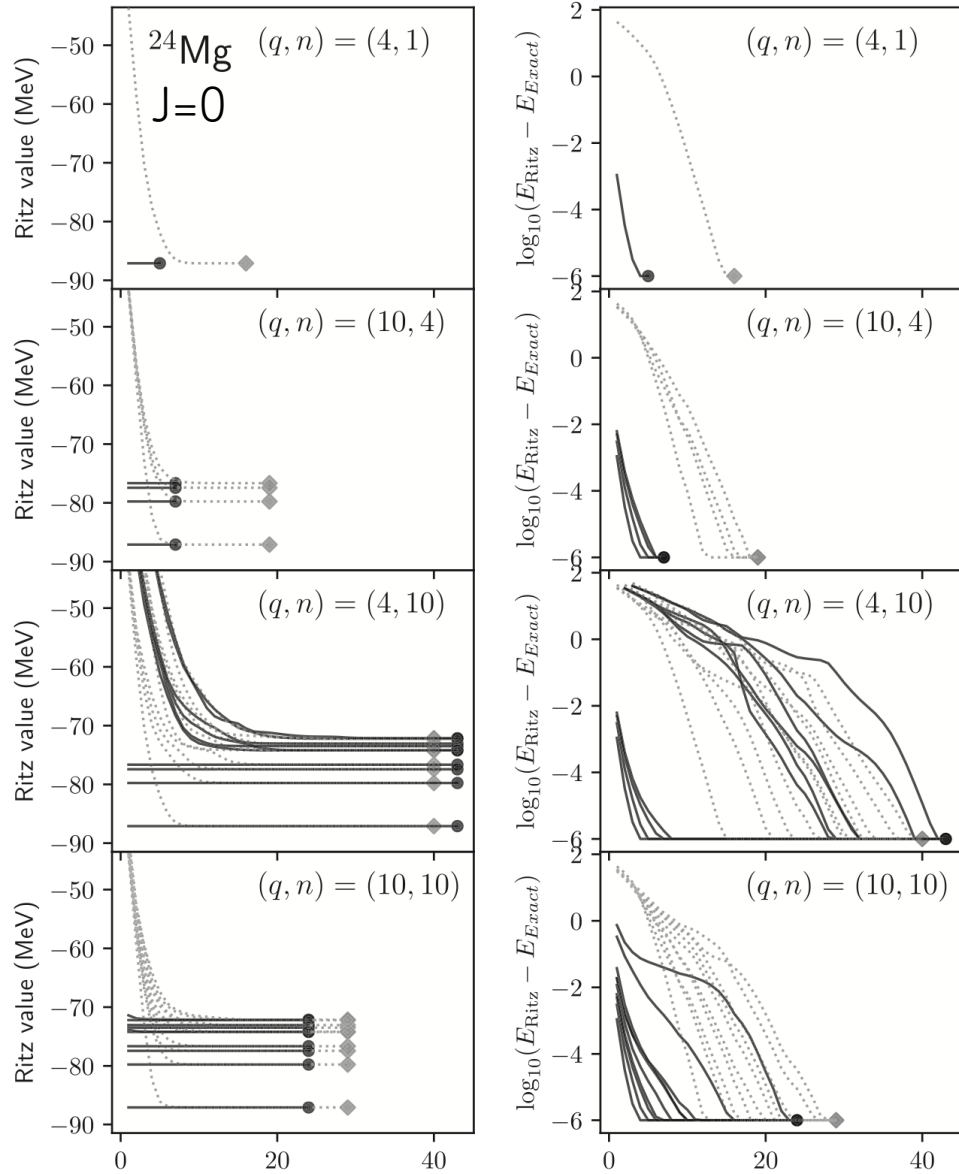
“validation” for 100 random parameters



EC approximates energies within a few percent accuracy

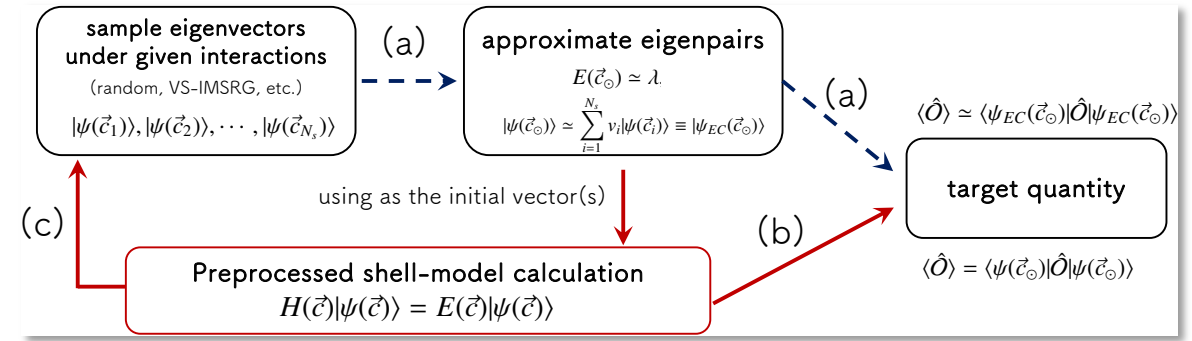
# EC: As a Lanczos preprocessor

SY and N.Shimizu, [PTEP 2022 053D02 \(2022\)](#)



Number of H operation during the (block) Lanczos method

converged results are obtained at ●/◆



$q$ : size of initial “block” vector

$n$ : # of excited states of interest

dotted: initialized by random vectors

solid: initialized by EC eigenvectors

Starting from better initial guess,  
# of manipulation could be reduced!!

Exception =>  $(q, n) = (4, 10)$

since the emulator is trained with 5 lowest states,  
such emulator do not have much info. on higher states

# EC: To feed more samples...

Sampling itself is not easy ...

$$\tilde{H}\vec{v} = \lambda N\vec{v},$$

$$\tilde{H}_{i,j} = \langle \psi(\vec{c}_i) | H(\vec{c}_\odot) | \psi(\vec{c}_j) \rangle, \quad \leftarrow \text{most time-consuming part}$$

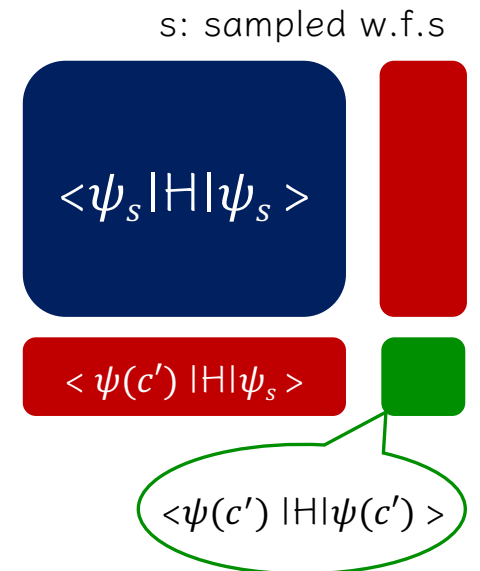
$$N_{i,j} = \langle \psi(\vec{c}_i) | \psi(\vec{c}_j) \rangle.$$

- You don't need to explicitly calculate  $H(\vec{c}_\odot) | \psi(\vec{c}_j) \rangle$  for each parameter  $\vec{c}_\odot$  to evaluate  $\tilde{H}$  above:

all you need is 1&2-body transition densities

$$\tilde{H}_{i,j} = \sum_k \underbrace{h_k^{(1)}}_{\text{SPEs}} \times \underbrace{\overline{\text{OBTD}}_{i,j,k}}_{\text{OBTD}} + \sum_k \underbrace{V_J(abcd)_k}_{\text{TBMEs}} \times \underbrace{\overline{\text{TBTD}}_{i,j,k}}_{\text{TBTD}}$$

- If you want to increase sample number (for better accuracy), prepare new sample (green) and calc. overlap (transition densities) between new w.f. and previous samples (red)





# First attempt to apply DMD to IMSRG

Jacob Davison, Ph.D. dissertation, MSU, 2023:

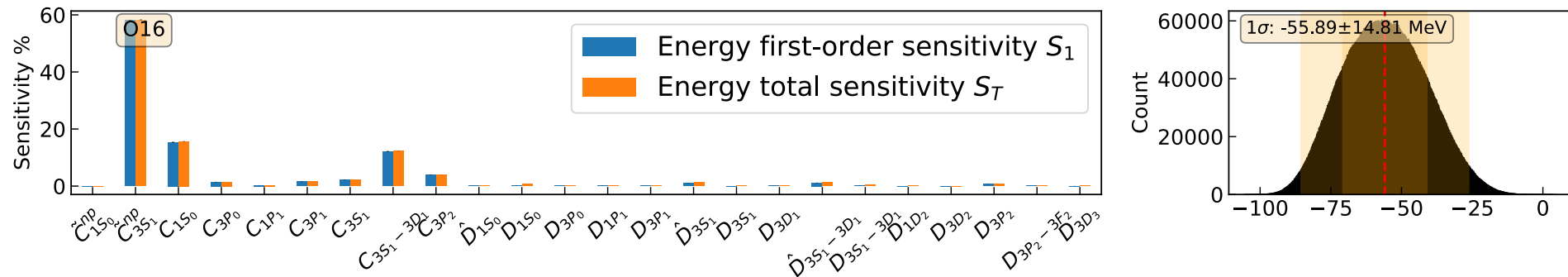
“Theoretical and computational improvements to the in-medium similarity renormalization group”

$$H(s) = \mathcal{U}(s)H(0) = \{\Phi \exp(\Omega t)\Phi^\dagger\}H(0)$$

“DMD emulation will treat  $H(s)$  as a one-dimensional vector” chapter7

chapter8

sensitivity analysis

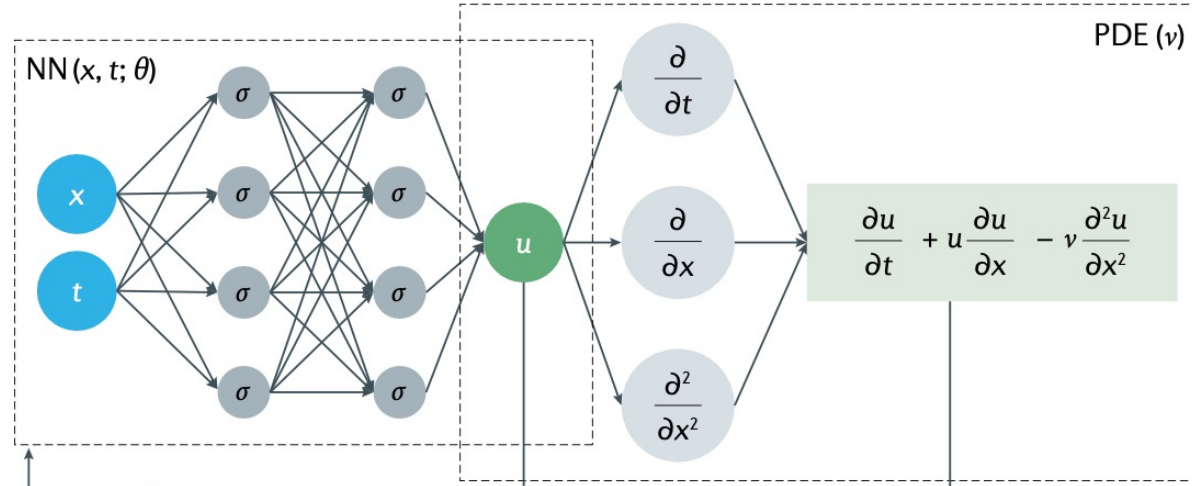


**Figure 8.1** First-order and total sensitivity of the IMSRG(2) ground-state energy for  $^{16}\text{O}$  energy to variation of the LECs around the standard  $\text{N}^3\text{LO}(500)$  interaction of Entem and Machleidt [4]. The left panel contains the sensitivity information per LEC. The right panel plots the total variance in the energy. The shaded region in the variance plot represents the  $1\sigma$  and  $2\sigma$  range.

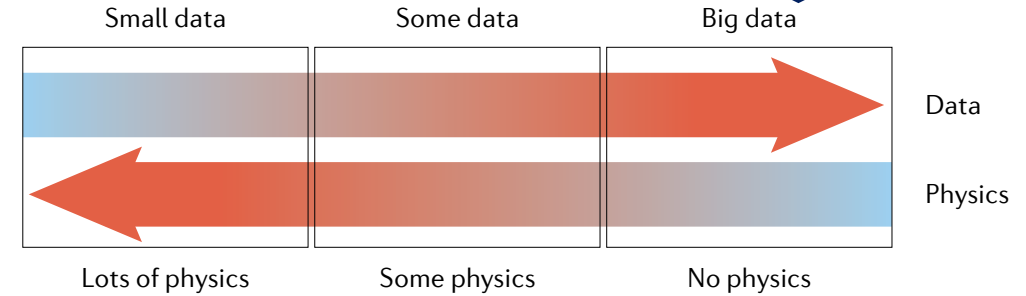
# ML example: Physics-Informed Neural Networks (PINNs)

Burgers eq. : 
$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = \nu \frac{\partial^2 u}{\partial x^2}$$

Ref: [Nature Reviews Physics 3, 422–440 \(2021\)](#)



e.g. Large Language Model



high-fidelity model calculations

$$\mathcal{L} = w_{\text{data}} \mathcal{L}_{\text{data}} + w_{\text{PDE}} \mathcal{L}_{\text{PDE}},$$

where

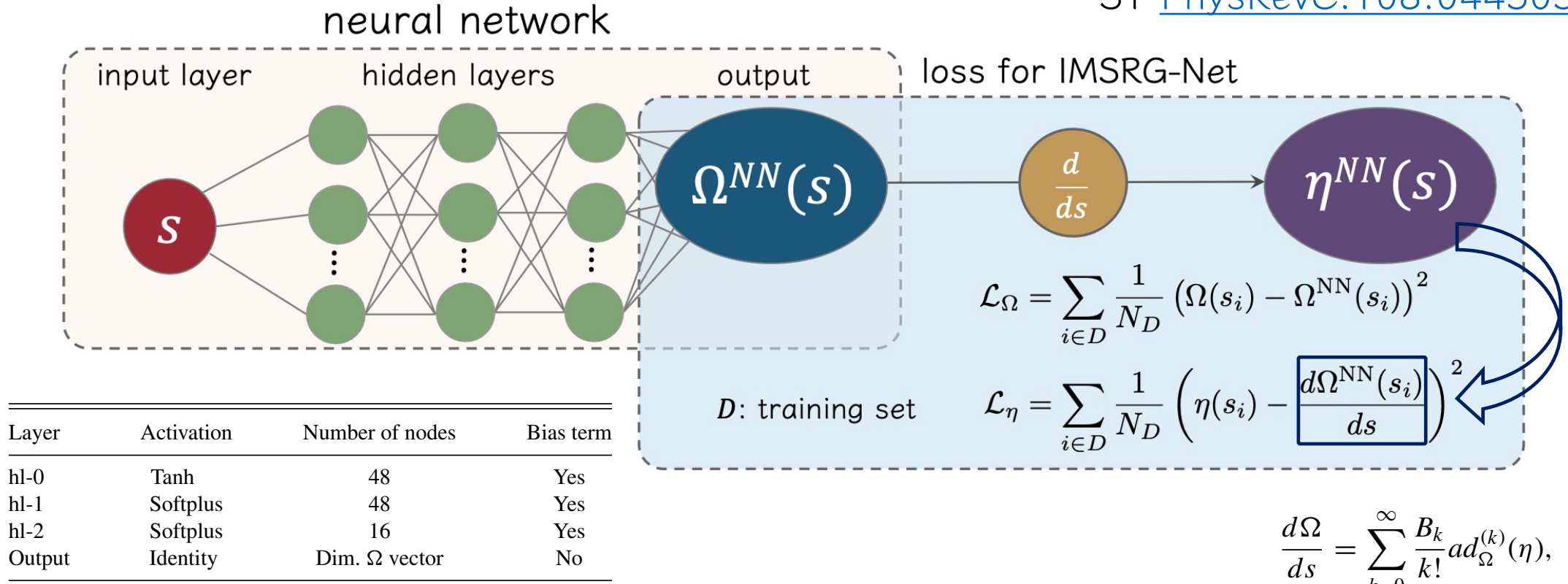
$$\mathcal{L}_{\text{data}} = \frac{1}{N_{\text{data}}} \sum_{i=1}^{N_{\text{data}}} (u(x_i, t_i) - u_i)^2 \quad \text{and}$$

$$\mathcal{L}_{\text{PDE}} = \frac{1}{N_{\text{PDE}}} \sum_{j=1}^{N_{\text{PDE}}} \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} - \nu \frac{\partial^2 u}{\partial x^2} \right)^2 \Big|_{(x_j, t_j)}.$$

“natural” expectation: you want neural networks approximating  $u(x, t)$  respecting underlying equation and boundary conditions, but...

# IMSRG-Net: PINN-based solver for IMSRG

SY [PhysRevC.108.044303 \(2023\)](#)



$$\frac{d\Omega}{ds} = \sum_{k=0}^{\infty} \frac{B_k}{k!} ad_{\Omega}^{(k)}(\eta),$$

$$ad_{\Omega}^{(k)}(\eta) = [\Omega, ad_{\Omega}^{(k-1)}(\eta)],$$

$$\boxed{ad_{\Omega}^{(0)}(\eta) = \eta}, \text{ taking leading term}$$

neural network part is simple Affine layers

and trained to minimize the sum of loss terms  $\mathcal{L} = \mathcal{L}_\Omega + \lambda_\eta \mathcal{L}_\eta$

$$\lambda_\eta = 100$$