

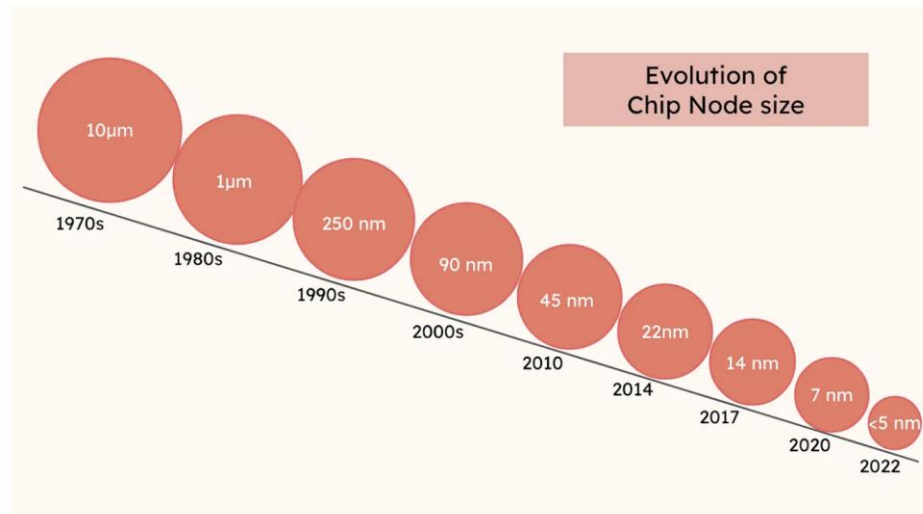
An application of the shift-invert Lanczos algorithm to the non-equilibrium Green function method

Kotaro Uzawa (Kyoto U.)

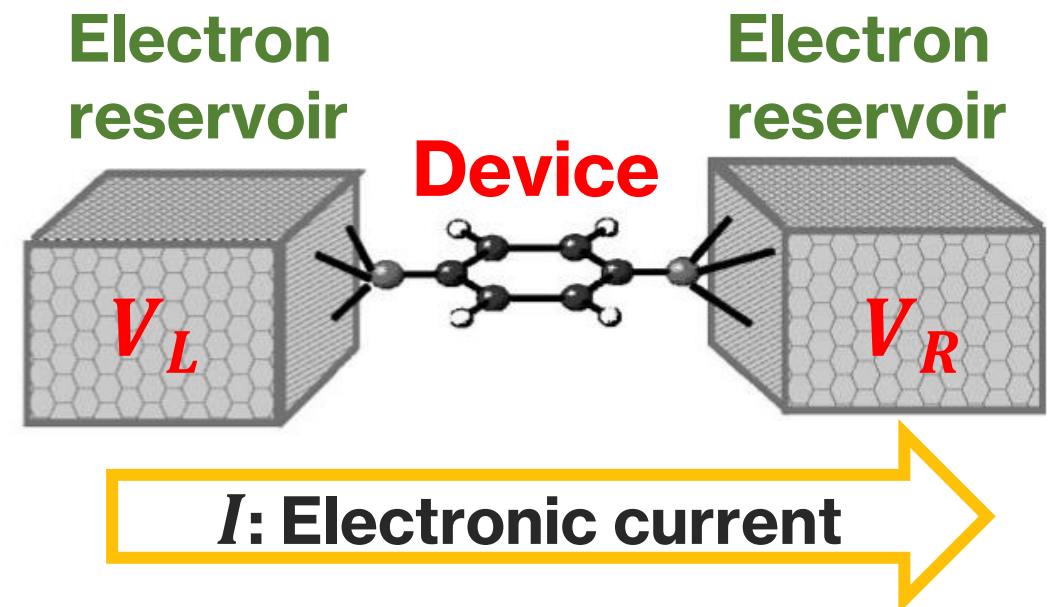
K. Uzawa and K. Hagino, arXiv:2408.06554

Non-Equilibrium Green Function (NEGF) method

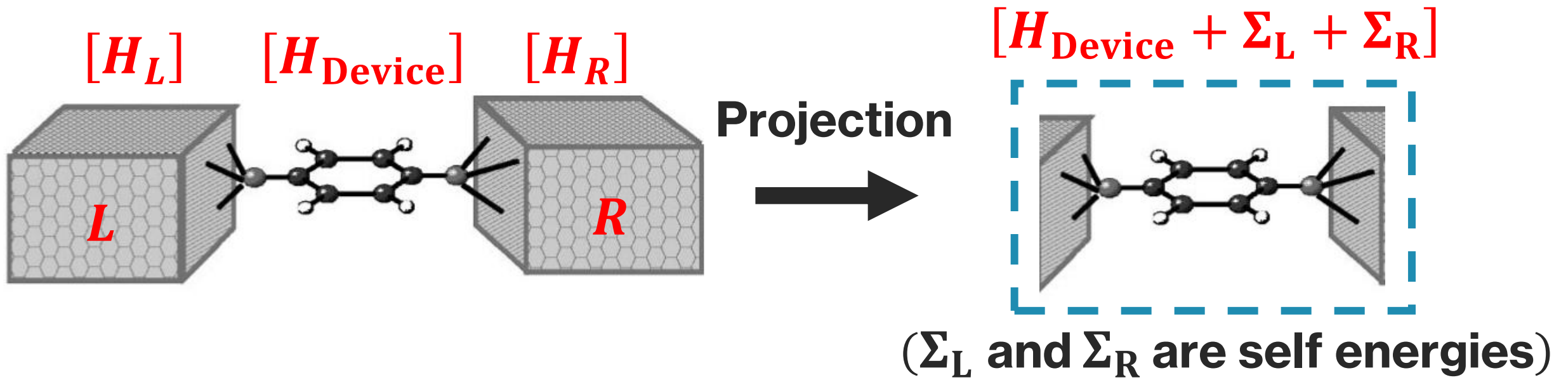
- ✓ With the time, the size of semiconductors continues to decrease
⇒ **Quantum effects become important** in the semiconductor design
- ✓ **NEGF method describes electron transports in nano-devices based on its microscopic Hamiltonian!**



<https://abhisheksingh-4899.medium.com>



P. S. Damle et al., Phys. Rev. B 64, 201403.



- Introduce Green's functions $G(E)$

$$G(E) = (E - H_{\text{Device}} - \Sigma_L - \Sigma_R)^{-1}$$

- Calculate transmission coefficient $T_{L \rightarrow R}$

$$T_{L \rightarrow R}(E) = \text{Tr}[\Gamma_L G(E) \Gamma_R G(E)^\dagger]$$

\propto (conductance)

$$\Sigma_L = \Delta_L - \frac{i}{2} \Gamma_L$$

$$\Sigma_R = \Delta_R - \frac{i}{2} \Gamma_R$$

Nuclear Fission

Michael Bender et al., J. Phys.

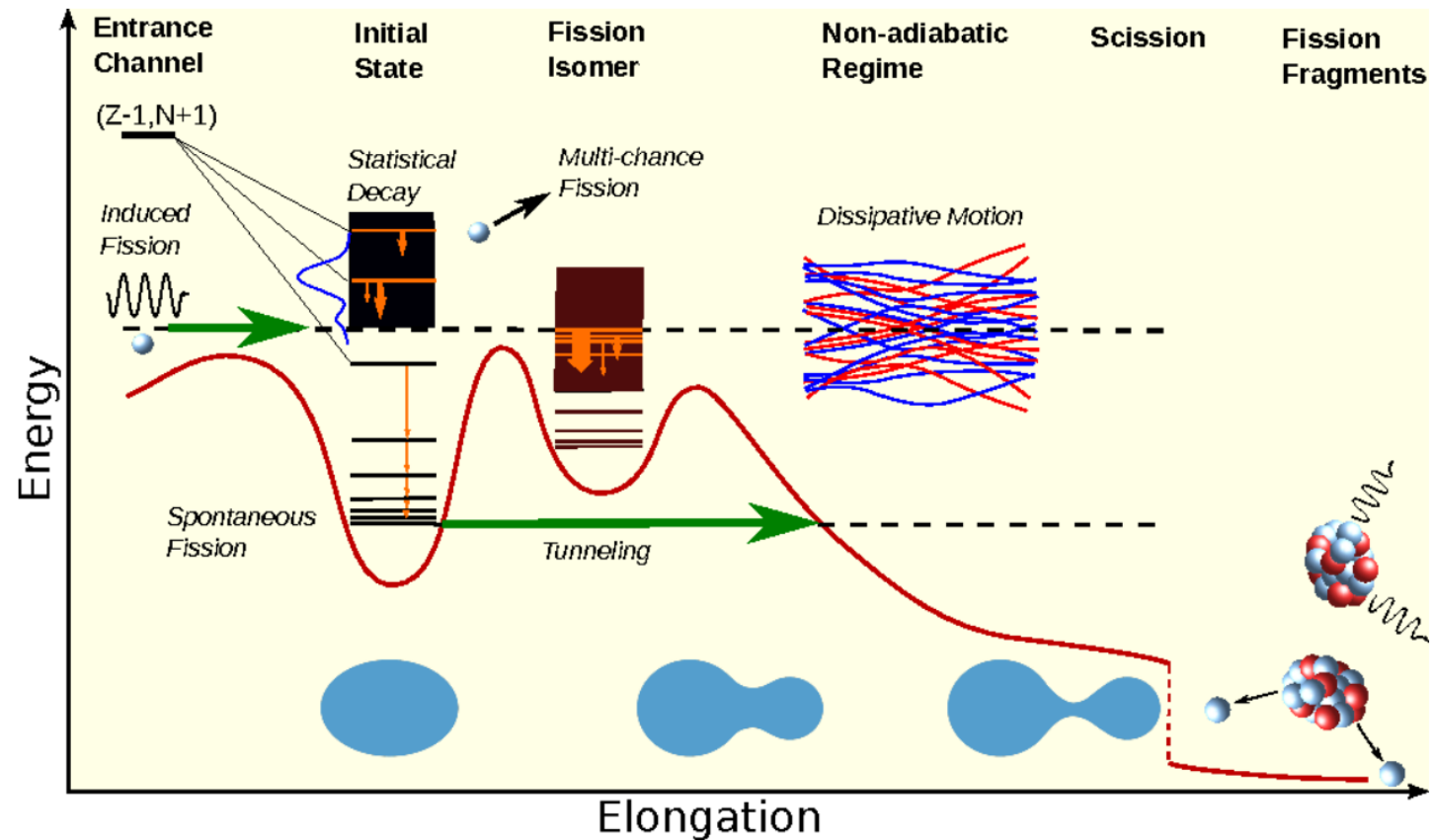
G: Nucl. Part. Phys. 47 113002

✓ A process of a heavy nucleus splitting into two smaller nuclei

✓ It plays an important role in

- nuclear energy
- nucleosynthesis
- RI beam production ...

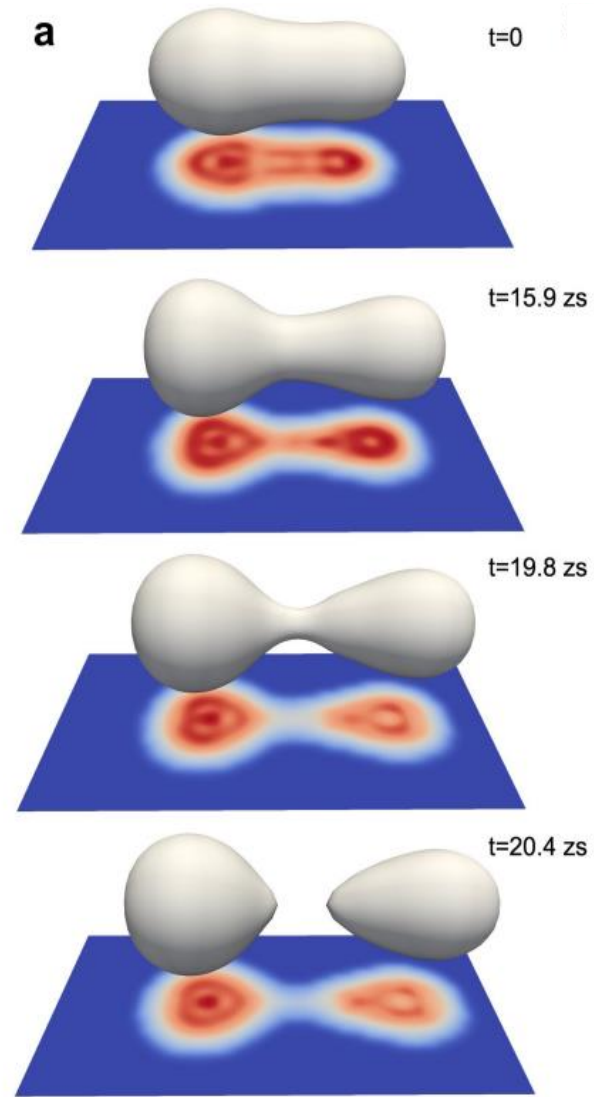
✓ Its microscopic description is extremely difficult !



Difficulties of microscopic description of fission come from ...

① **Fission is a large-amplitude motion**

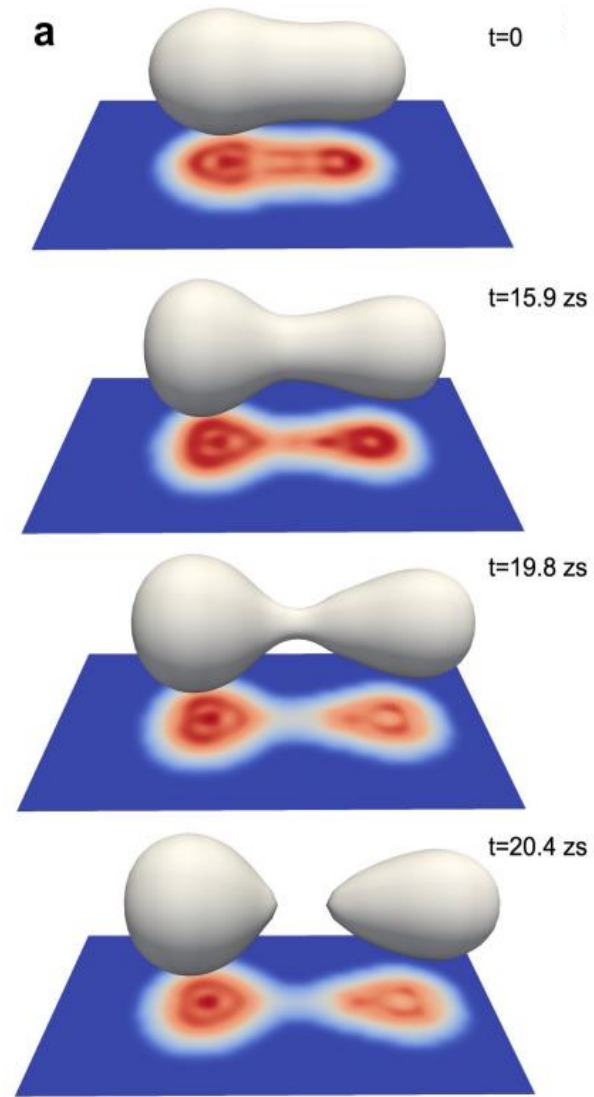
(**×** Perturbation or Linear response theory)



G. Scamps and C. Simenel,
Nature 564 382 (2018).

Difficulties of microscopic description of fission come from ...

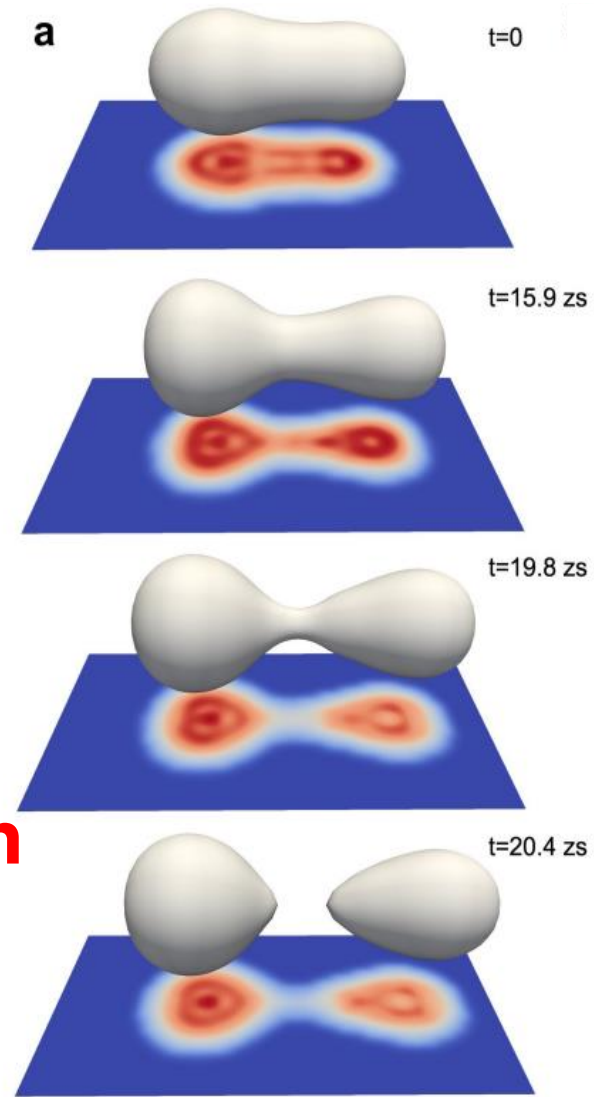
- ① Fission is a large-amplitude motion
(\times Perturbation or Linear response theory)
- ② **Interaction between collective motion
and single-particle motions is complex**



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- ① Fission is a large-amplitude motion
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- ③ **Connection between nuclear structure calculation
and reaction calculation is not straightforward**

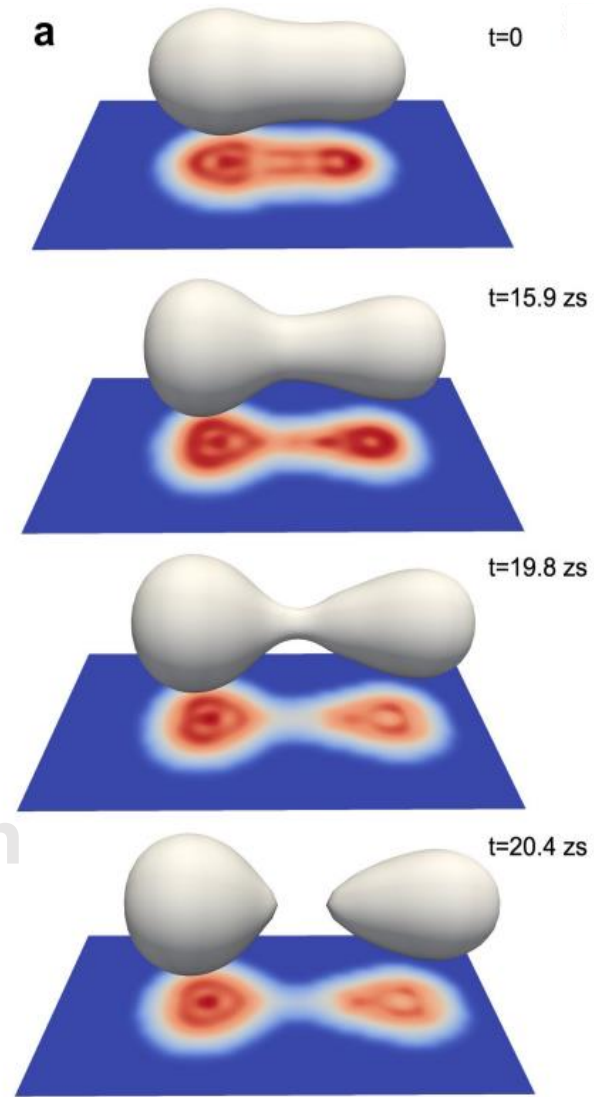


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**We approach nuclear fission using
Non-Equilibrium Green Function method !**

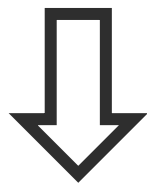


G. Scamps and C. Simenel,
Nature 564 382 (2018).

Theoretical Formulation

First, we prepare many-body basis

In nuclear fission, both **excitation** and **deformation** are important

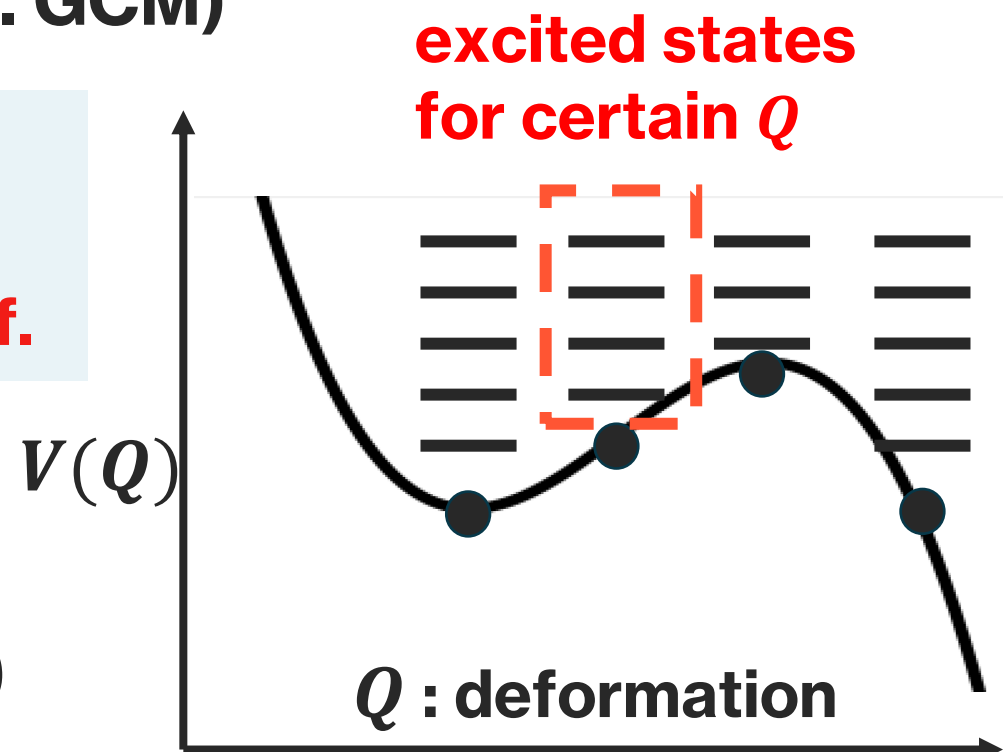


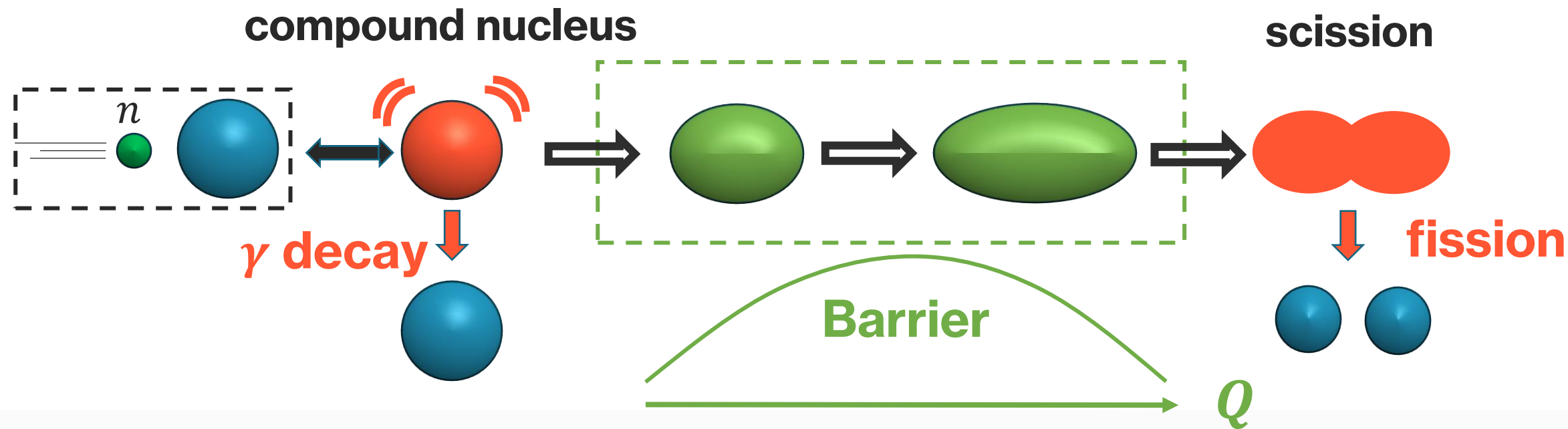
Superpose Hartree-Fock w.f. $|HF(Q, E_i)\rangle$ (i.e. GCM)

$$|\Psi\rangle_{\text{GCM}} = \sum_i \int dQ f(Q, E_i) \underbrace{|HF(Q, E_i)\rangle}_{\text{Hartree-Fock w.f.}}$$

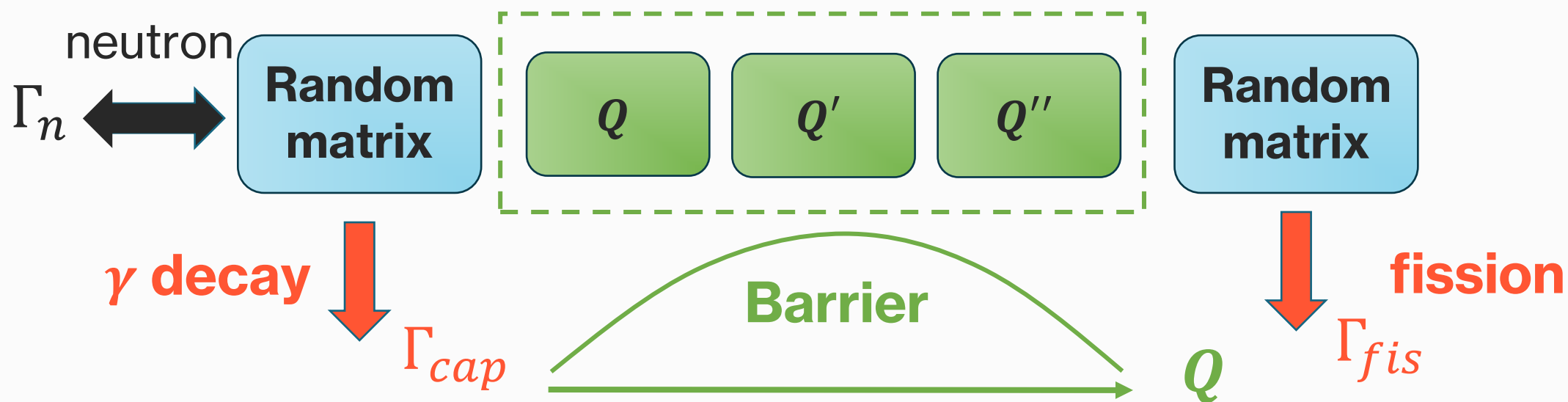
(Q : **deformation parameter**)

(E_i : **particle-hole excitation energy**)



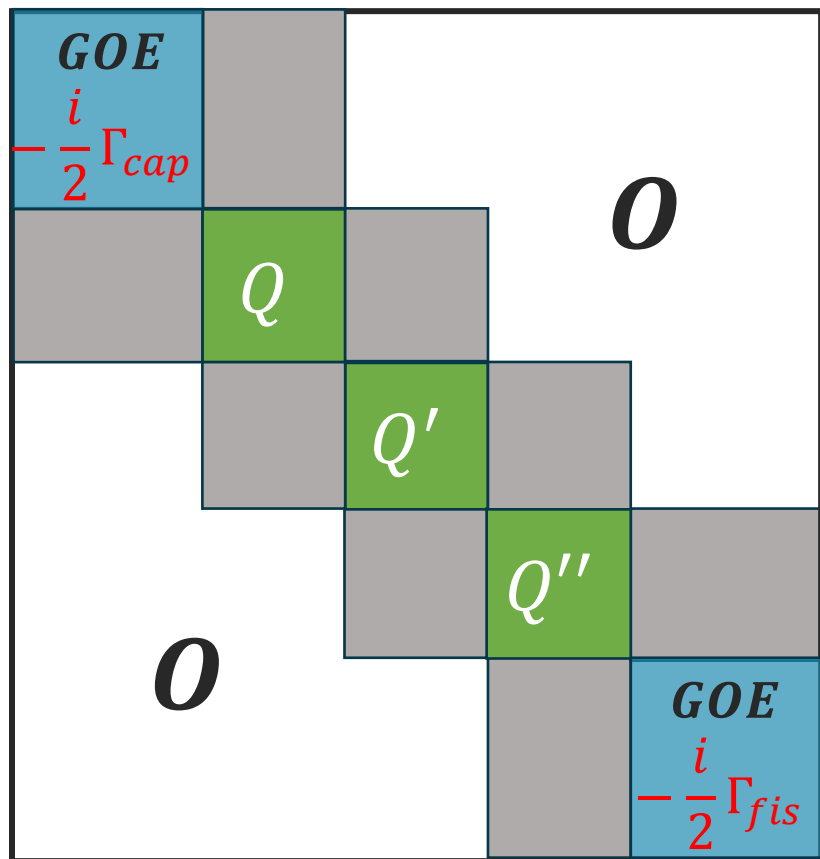


(Q is submatrix in the Hamiltonian)



Non-equilibrium Green's function method

H : Hamiltonian Matrix



**GOE means
Gaussian Orthogonal Ensemble**

&

Green's function

$$G(E) = \left[EN - H + \frac{i}{2}\Gamma \right]^{-1}$$

**(E : excitation energy)
(N : overlap matrix)**



$$T_{n,fis} = \text{Tr}[\Gamma_n G \Gamma_{fis} G^\dagger]$$
$$\sigma_{n,fis}(E) = \frac{\pi}{k_n^2} T_{n,fis}(E)$$

G. F. Bertsch and K. Hagino, PRC 107 044615 (2023).

K. Uzawa and K. Hagino, PRC 110 014321 (2024).

However, the dimension of the Hamiltonian matrix is extremely large

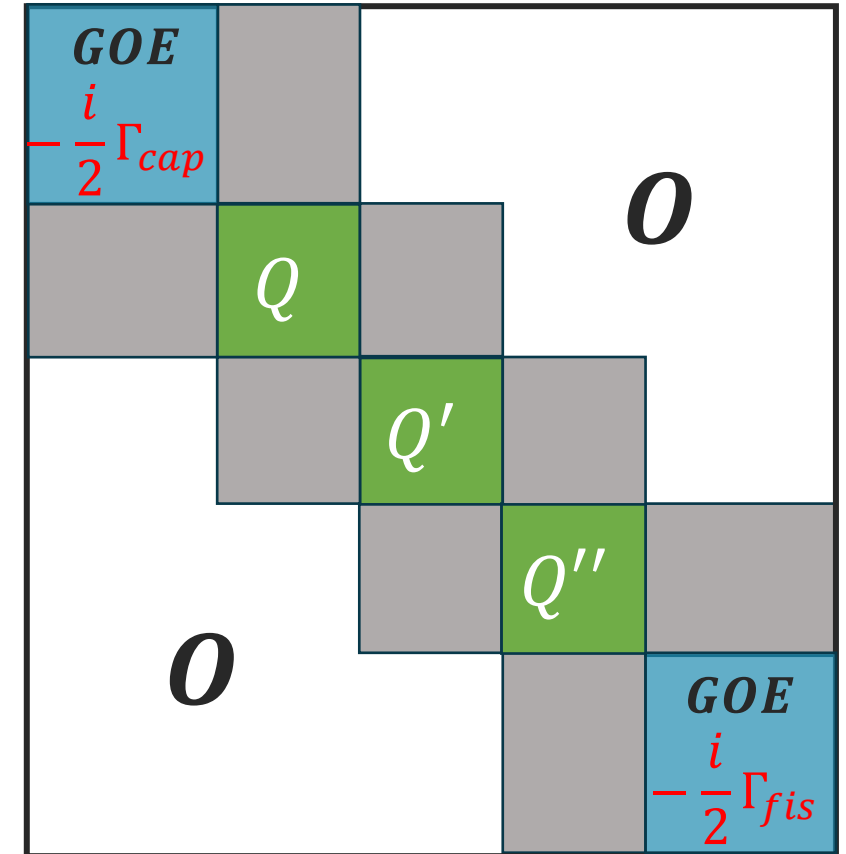
(In the case of $^{235}\text{U}(n, fis)$ reaction,
the dimension is $O(10^5) - O(10^6)$)

An inversion of the Green function matrix
becomes expensive

Green's function

$$G(E) = \left[EN - H + \frac{i}{2} \Gamma \right]^{-1}$$

H : Hamiltonian Matrix



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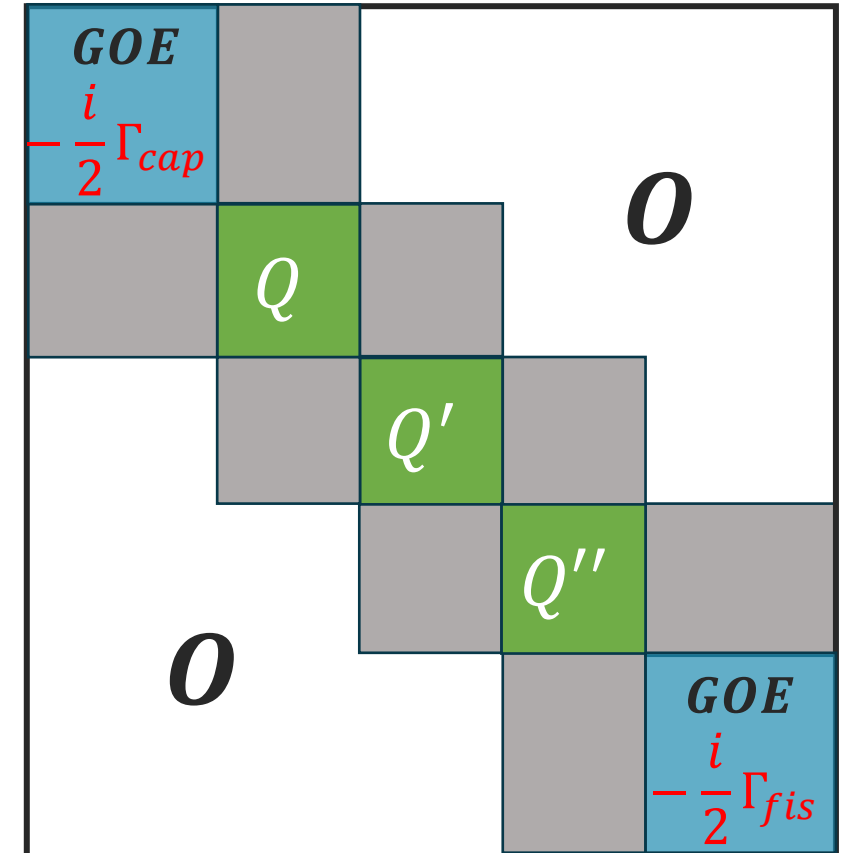
An inversion of the Green function matrix
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Green's function

$$G(E) = \left[EN - H + \frac{i}{2} \Gamma \right]^{-1}$$

⇒ Application of the Lanczos method !

H : Hamiltonian Matrix



The Lanczos method

Solve eigenvalue eq. of matrix A

$$A\vec{x} = \lambda\vec{x} \quad (A \text{ is a symmetric } N \times N \text{ matrix})$$

We introduce a subspace called **Krylov subspace** (\vec{q} is arbitrary vector)

$$\text{span}\{\vec{q}, A\vec{q}, A^2\vec{q}, \dots, A^{n-1}\vec{q}\}$$

$Ax = \lambda x$ is solved in the Krylov subspace,

and **few specific eigenstates** are calculated with $O(n^2 N)$

(direct diagonalization requires $O(N^3)$)

Spectral decomposition of Green's function

$$G(E)_{\mu\mu'} = \sum_{\lambda} f_{\lambda}^*(\mu') \frac{1}{E - E_{\lambda}} f_{\lambda}(\mu)$$

$$(Hf_{\lambda} = E_{\lambda}Nf_{\lambda})$$

$$\left(\begin{array}{l} E_{\lambda} : \text{eigenenergy} \\ f_{\lambda} : \text{eigenvector} \end{array} \right)$$

Substituting it, fission probability $T_{n,fis}(E)$ becomes

$$T_{n,fis}(E) = \text{Tr}[\Gamma_n G(E) \Gamma_{fis} G^{\dagger}(E)] = \gamma_n \gamma_{fis} \sum_{\lambda\lambda'} \frac{f_{\lambda}^{(a)} f_{\lambda'}^{(a)*} f_{\lambda}^{(b*)} f_{\lambda'}^{(b)}}{(E - E_{\lambda})(E - E_{\lambda'})^*}$$

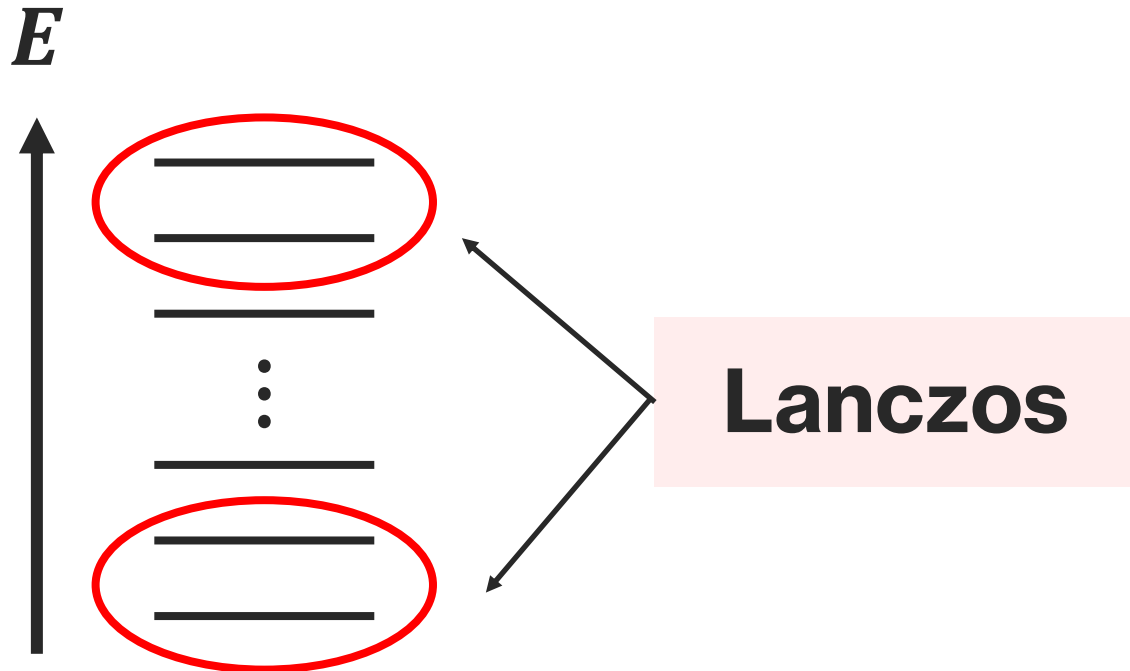
Eigenstates with $E_{\lambda} \simeq E$ are dominant !

(due to $\frac{1}{(E - E_{\lambda})(E - E_{\lambda'})^*}$ factor)

◆ Our approach

Calculate eigenstates $|\lambda\rangle$ with $E_\lambda \simeq E$ using the Lanczos method

However, the Lanczos method gives eigenstates around the ground state

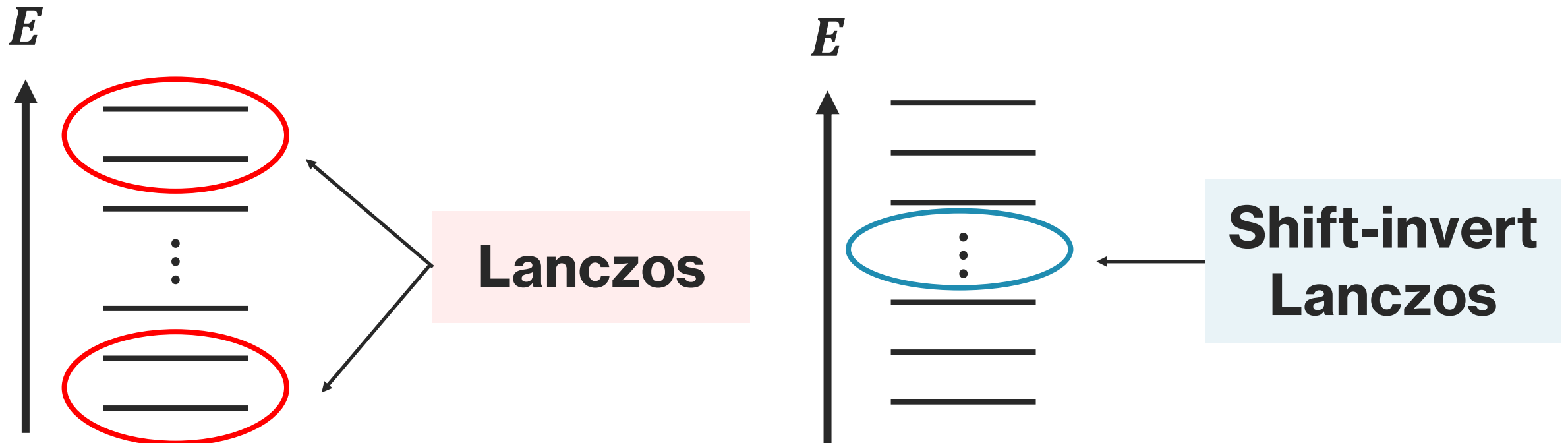


◆ Our approach

Calculate eigenstates $|\lambda\rangle$ with $E_\lambda \simeq E$ using the Lanczos method

However, the Lanczos method gives eigenstates around the ground state

⇒ Shift-invert Lanczos method



Shift-invert transformation

$$Hf = ENf$$

$$\downarrow -\sigma N$$

$$(H - \sigma N)f = (E - \sigma)Nf,$$

$$\downarrow \times (E - \sigma)^{-1} (H - \sigma N)^{-1}$$

$$(H - \sigma N)^{-1} Nf = (E - \sigma)^{-1} f.$$

Applying the Lanczos method to the last eq.

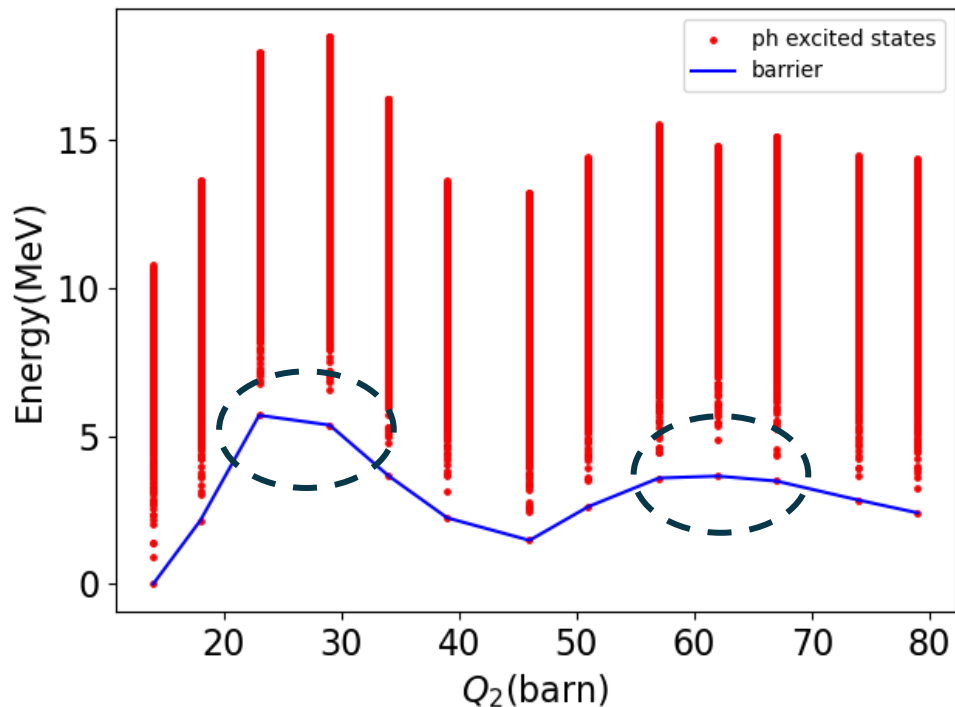
- ✓ **Eigenvalue E is transformed into $(E - \sigma)^{-1}$**
- ✓ **Lanczos method calculates eigenstates with large eigenvalues**
 \Rightarrow the eigenvalue $E \simeq \sigma$ is obtained!

Application to $^{235}\text{U}(n, f)$

Solve Skyrme-Hartree-Fock equation for ^{236}U

⇒ Calculate **fission barrier** and **particle-hole excited states**

fission barrier & **ph excited states**



$H =$

$$\begin{pmatrix} H_{\text{GOE}}^{(L)} & V_L & & & \\ V_L^T & H_1 & V_{1,2} & & O \\ & V_{2,1} & H_2 & V_{2,3} & \\ & & \ddots & & \\ O & & & V_{11,12} & H_{12} & V_R \\ & & & & V_R^T & H_{\text{GOE}}^{(R)} \end{pmatrix}$$

Hamiltonian matrix
(its dimension is 66103)

GCM calculation

- ✓ Calculate Hamiltonian matrix H_{ij} and overlap matrix N_{ij} (GCM)

$$H_{ij} = \langle Q, E | \hat{H} | Q', E' \rangle \text{ and } N_{ij} = \langle Q, E | Q', E' \rangle$$

- ✓ As residual interactions, we apply

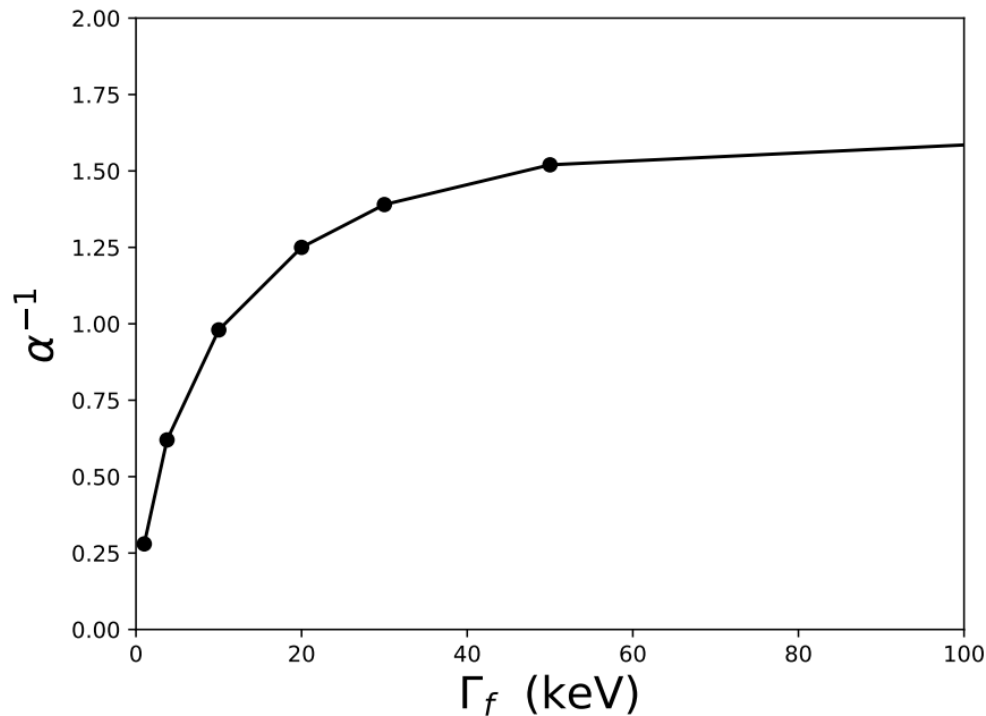
constant pairing interaction and **diabatic** interaction

$$H_{pair} = -GP^\dagger P \quad (P = \sum_\nu a_{\bar{\nu}} a_\nu)$$

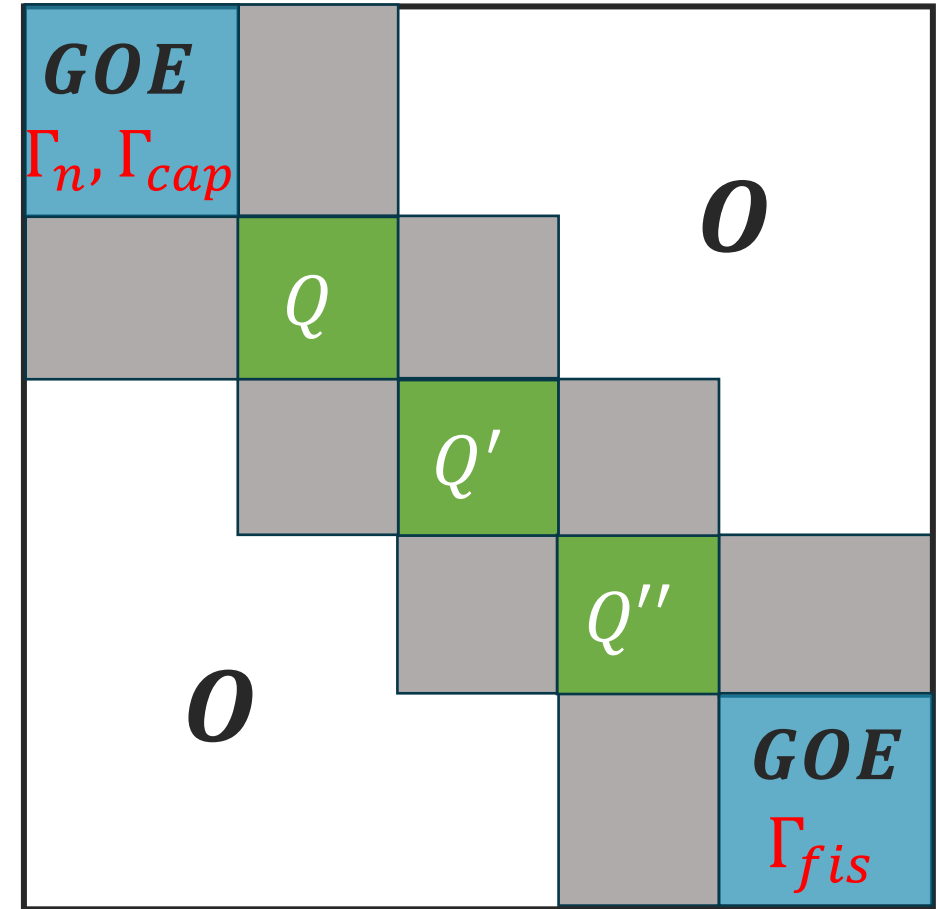
$$\frac{\langle Q, E_\mu | v_{db} | Q', E_{\mu'} \rangle}{\langle Q, E_\mu | Q', E_{\mu'} \rangle} = \frac{E(Q, E_\mu) + E(Q', E_{\mu'})}{2} + h_2 \ln(\langle Q, E_\mu | Q', E_{\mu'} \rangle).$$

Decay widths

- ✓ Γ_n and Γ_{cap} are fitted to empirical values in the RIPL library
- ✓ σ_{fis} is known to be insensitive to the fission width Γ_{fis} (fixed to 125 keV)



H : Hamiltonian Matrix



G. F. Bertsch and K. Hagino, PRC 107 044615 (2023).

K. Uzawa and K. Hagino, PRC 108 024319 (2024).

Calculate fission probability $T_{n,fis}$

$$T_{n,fis}(E) = \gamma_n \gamma_{fis} \sum_{\lambda\lambda'} \frac{f_{\lambda}^{(a)} f_{\lambda'}^{(a)*} f_{\lambda}^{(b*)} f_{\lambda'}^{(b)}}{(E - E_{\lambda})(E - E_{\lambda'})^*}$$

$(E_{\lambda} : \text{eigenenergy})$
 $(f_{\lambda} : \text{eigenvector})$

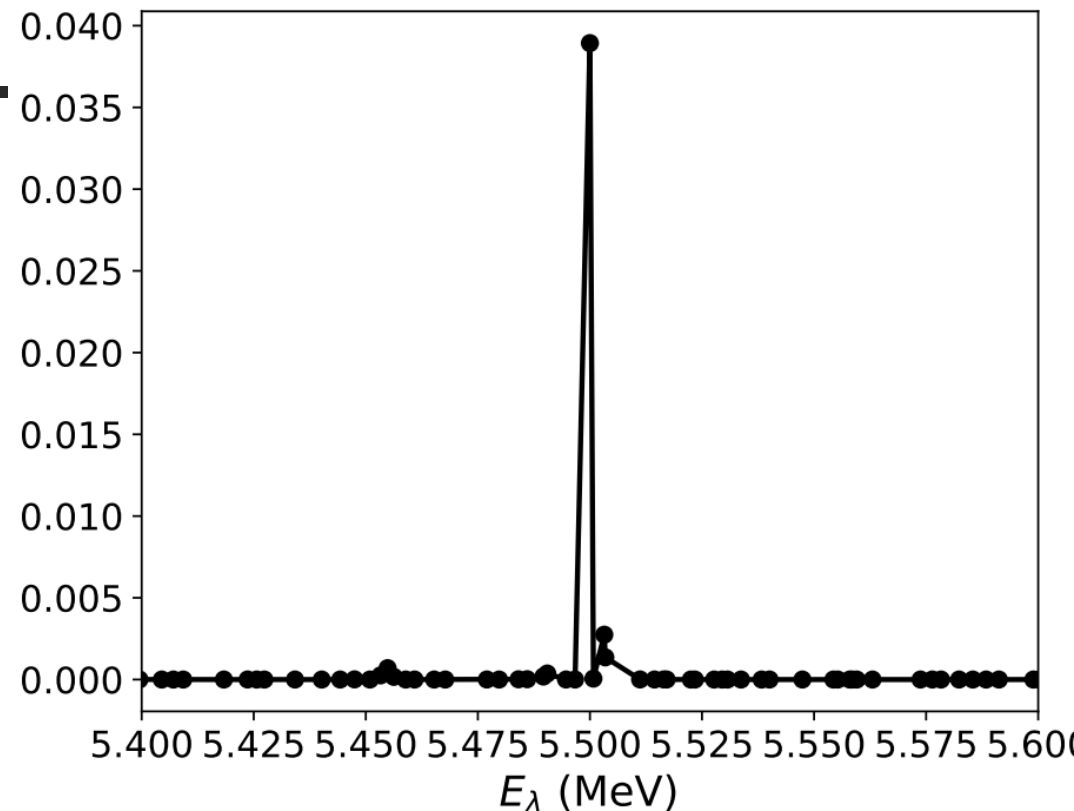
The summation is $\sum_{\lambda=1}^{66103}$.

But, λ with $E_{\lambda} \simeq E$ is dominant (right figure).

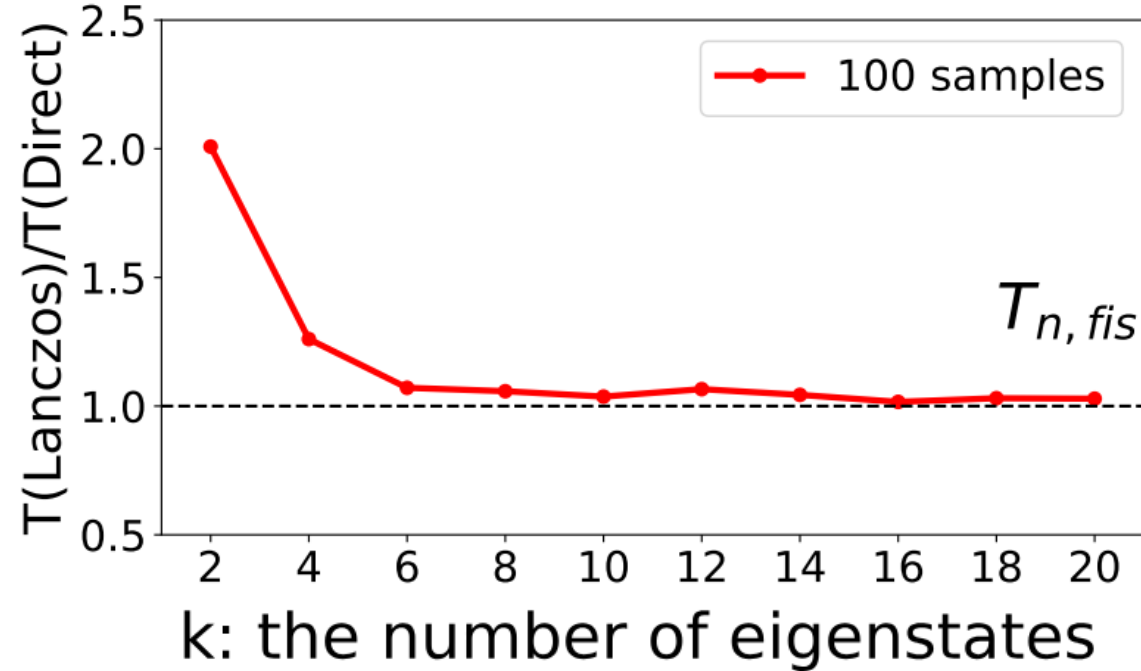
Reduce $\sum_{\lambda=1}^{66103} \rightarrow \sum_{\lambda=1}^k$.

Here k is the number of eigenstates
calculated by shift-invert Lanczos

$$\sum_{\lambda} \frac{|f_{\lambda}^{(a)}|^2 |f_{\lambda}^{(b)}|^2}{|E - E_{\lambda}|^2} \quad (E = 5.5 \text{ MeV})$$



The ratio $\frac{T_{n,\text{fis}}(\text{Lanczos})}{T_{n,\text{fis}}(\text{Direct})}$ as a function of k (# of eigenstates)

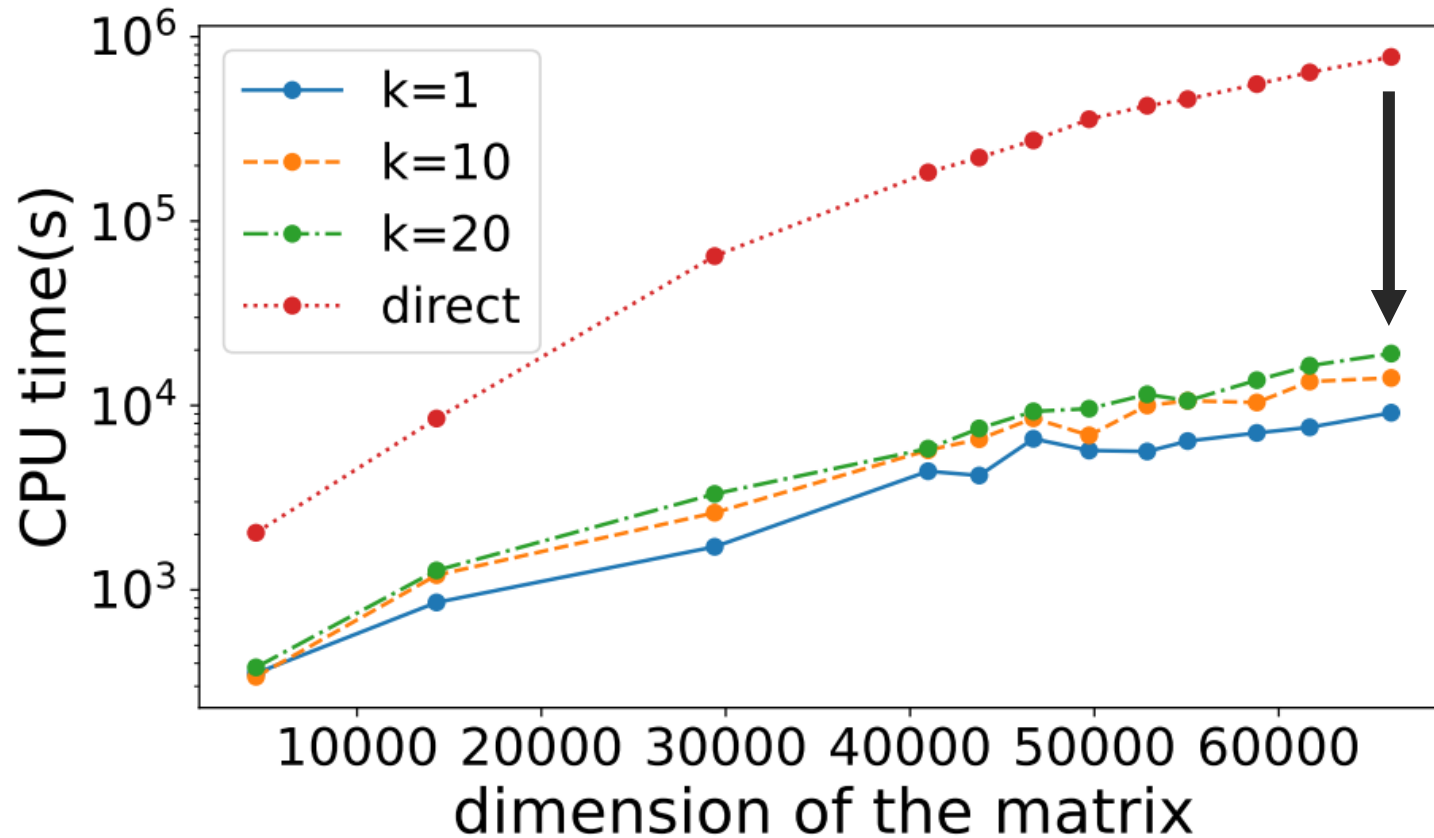


The ratio is close to one with $k \geq 6$, and the error is about 1%

The CPU time with different matrix dimensions

Red line : direct matrix inversion $G(E) = (E - H)^{-1}$ with LAPACK

Other lines : the shift-invert Lanczos calculations with ARPACK



**30-40 time faster
than direct calculations**

Summary

- ◆ **NEGF method is a promising way to describe fission microscopically, but requires huge numerical costs.**
- ◆ **We propose a novel method of calculating $T_{n,fis}(E)$ based on the shift-invert Lanczos + spectral decomposition of $G(E)$**
- ◆ **The error is about 1% and CPU time is about 30-40 time faster.**

Shift-Invert Lanczos method is implemented in ARPACK library or SciPy library (scipy.sparse.linalg.eigsh).

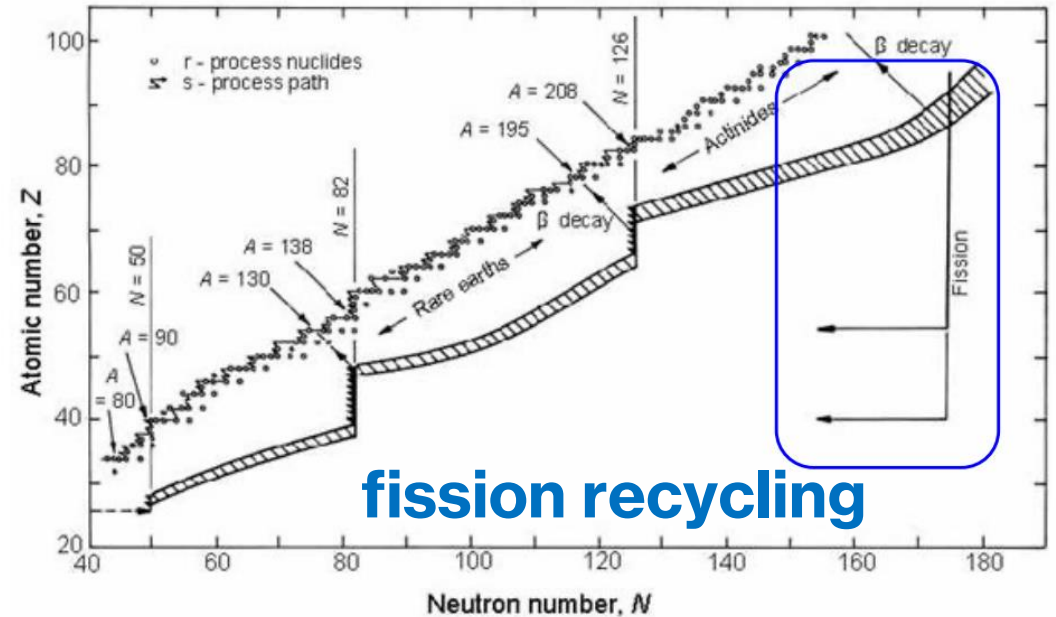
Why we need microscopic theory of fission ?

- In r-process, nuclear fission plays an important role (fission recycling)



fission of neutron-rich nuclei

\Rightarrow low S_n , low E^* , and low $\rho(E^*)$



Hauser-Feshbach theory or Langevin eq. may not be applicable...

\Rightarrow Microscopic models without phenomenological assumptions

We apply **the non-equilibrium Green function method to nuclear fission!**

