Tensor network toward the lattice QCD

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2024.9.25 @ RPMBT22

50th anniversary of lattice QCD

• "Confinement of quarks" by K. G. Wilson in 1974

Kenneth G. Wilson

Laboratory of Nuclear Studies, Cornell University, Ithaca, New York 14850 (Received 12 June 1974)

A mechanism for total confinement of quarks, similar to that of Schwinger, is defined which requires the existence of Abelian or non-Abelian gauge fields. It is shown how to quantize a gauge field theory on a discrete lattice in Euclidean space-time, preserving exact gauge invariance and treating the gauge fields as angular variables (which makes a gauge-fixing term unnecessary). The lattice gauge theory has a computable strong-coupling limit; in this limit the binding mechanism applies and there are up free quarks. There is unfortunately no Lorentz (or Euclidean) invariance in the strong-coupling limit. The strong-coupling expansion involves sums over all quark paths and sums over all surfaces (on the lattice) joining quark paths. This structure is reminiscent of relativistic string models of hadrons.

I, INTRODUCTION

The success of the quark-constituent picture

Successes over the past 50 years

- Hadron mass, matrix elements
- Nuclear force, ...
- Monte Carlo simulations on supercomputers
- Great development of algorithms and supercomputers

particles over short times and short distances. The polarization effects which prevent the ap- Wilson, PRD10(1974)2445 pearance of electrons in the final state take place



Can we solve the QCD using only the MC?

- No, unfortunately
 - MC has the **sign problem**, that has prevented numerous studies
 - QCD at finite density
 - Real time evolution
 - Strong CP problem
 - Chiral gauge theory
 - Other finite-density systems
 - Supersymmetry



Fukushima-Hatsuda, Rept. Prog. Phys. 74(2011)014001

- ...
- Many attempts have been made to overcome the sign problem in the computational physics

Cf. Talk by Siu Chin on Tue, 24/9

Cf. Talk by Alexander Lichtenstein on Wed, 25/9

Tensor Network?

 Theoretical or numerical methods based on representing quantum manybody systems as a network of numerous tensors (multi-index objects)

Cf. Talk by Garnet Kin-Lic Chan on Wed 25/9

- Originated in statistical physics
- Extremely successful for 1D quantum systems
 - Density Matrix RG
 - Matrix Product State (MPS)

White, Phys. Rev. Lett. 69 (1992) 2863 Feenberg Memorial Medal in 2019 (RPMBT-20)

Schollwöck, Rev. Mod. Phys. 77 (2005) 259-315

- Accurate information compression with finite bond dimensions is available when the entanglement of the system is small
- No sign problem



Baxter, J. Math. Phys. 9 (1968) 650

How to use TN for many-body problems

• Within the Hamiltonian formalism

- Cf. Poster by Ryo Watanabe [Board:13]
- TN as a variational ansatz for the **many-body state** $|\Psi\rangle$

 $|\Psi\rangle = \langle \Psi|$

- Determines the ground state and excited states
- Variational optimization of the TN

White, Phys. Rev. Lett. 69 (1992) 2863 Nishino-Hieida-Okunishi-Maeshima-Akutsu-Gendiar, Prog. Theor. Phys. 105 (2001) 409-17 Vidal, Phys. Rev. Lett. 91 (2003) 147902 Verstraete-Cirac, arXiv:cond-mat/0407066

- Within the Lagrangian formalism
 - Path integral Z is described as a TN
 - Contraction of the TN
 - Real-space RG on the TN



Nishino-Okunishi, J. Phys. Soc. Japan 65 (1996) 891 Levin-Nave, Phys. Rev. Lett. 99 (2007) 120601 Gu-Wen, Phys. Rev. B80(2009)155131 Evenbly-Vidal, Phys. Rev. Lett. 115(2015)180405

TN toward the lattice QCD

- "Kogut ladder" can be a good roadmap
 - Gradually try to increase symmetry and dimensionality
- Advantage of TN over MC
 - No sign problem
 - Can directly deal with **fermions**
 - Thermodynamic limit can be handled directly when the system has the translational symmetry
- What should be addressed
 - How to regularize **bosons**
 - Accuracy vs computational cost
- Cf. "A quantum-simulation program for QCD?" by Zohreh Davoudi in RPMBT-21

II. LATTICE FIELD THEORY

A. The Kogut sequence: From Ising to QCD

In the early 1970s, QCD appeared to be a strong candidate for a theory of strong interactions involving quarks and gluons. However, the perturbative methods that provided satisfactory ways to handle the electroweak interactions of leptons failed to explain confinement, mass gaps, and chiral symmetry breaking. A nonperturbative definition of QCD was needed. In 1974, Wilson proposed (Wilson, 1974) a lattice formulation of QCD where the SU(3) local symmetry is exact. As this four-dimensional model is fairly difficult to handle numerically, a certain number of research groups started considering simpler lattice models in lower dimensions and then increased symmetry and dimensionality. This led to a sequence of models, sometimes called the "Kogut ladder," that appears in the reviews of Kogut (1979, 1983) and was later addressed with small modifications by Polyakov (1987) and Itzykson and Drouffe (1991).

The sequence is approximately the following:

- (1) D = 2 Ising model
- (2) D = 3 Ising model and its gauge dual
- (3) D = 2 O(2) spin and Abelian Higgs models
- (4) D = 2 fermions and the Schwinger model
- (5) D = 3 and 4U(1) gauge theory
- (6) D = 3 and 4 non-Abelian gauge theories
- (7) D = 4 lattice fermions
- (8) D = 4 QCD

Starting from 2D, (1+1)D systems

- Schwinger model (= 2D QED)
 - The most QCD-like theory in 2D
 - Chiral symmetry breaking
 - Confinement
 - Topological term
 - The first application of TN to the lattice gauge theory
 - Critical property at $\theta = \pi$
 - 2D Ising universality

$$\left(\frac{m}{g}\right)_c = 0.3335(2).$$
 $\nu = 1.01(1)$
 $\beta/\nu = 0.125(5).$

Bañuls-Cichy-Jansen-Cirac, JHEP13(2013)158 Byrnes-Haegeman-Acoleyen-Verschelde-Verstraete, PRL113(2014)091601 Shimizu-Kuramashi, PRD90(2014)014508, 074503 Bañuls-Cichy-Cirac-Jansen-Kühn, PRL118(2017)071601 Pichler-Dalmonte-Rico-Zoller-Montangero, PRX6(2016)011023 Saito-Bañuls-Cichy-Cirac-Jansen, Lattice2014, Lattice2015 Zapp-Orús, PRD95(2017)114508 Shimizu-Kuramashi, PRD97(2018)034502 Magnifico-Vodola-Ercolessi-Kumar-Müller-Bermudez, PRD99(2019)014503 Magnifico-Vodola-Ercolessi-Kumar-Müller-Bermudez, PRB100(2019)115152 Funcke-Janse-Kühn, PRD101(2020)054507 Butt-Catterall-Meurice-Sakai-Unmuth-Yockey, PRD101(2020)094509 Honda-Itou-Tanizaki, JHEP11(2022)141 Angelides-Funcke-Janse-Küh, PRD108(2023)0145156 Dempsey-Klebanov-Benjamin-Søggard-Zan, PRL132(2024)031603 Yosprakob-Nishimura-Okunishi, JHEP11(2023)187 Itou-Matsumoto-Tanizaki, JHEP11(2023)231, arXiv:2407.11391 Kanno-SA-Murakami-Takeda, Lattice2024



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Starting from 2D, (1+1)D systems

• 2D QCD

- Testing ground for non-Abelian gauge theories
 - SU(N) DoFs have to be regularized
 - Removing the gauge DoFs
 - Irreducible representation
 - Quantum link formulation
 - Qubit regularization
 - Subgroups, Sampling, ...
 - Regularization is also necessary in quantum computing

Kühn-Zohar-Cirac-Bañuls, JHEP07(2015)130 Silvi-Rico-Dalmonte-Tschirsich-Montangero, Quantum1(2017)9 Bañuls-Cichy-Cirac-Jansen-Kühn, PRX7(2017)041046 Sala-Shi-Kühn-Bañuls-Demler-Cirac, PRD98(2018)034505 Silvi-Sauer-Tschirsich-Montangero, PRD100(2019)074512 Bazavov-Catterall-Jha-Unmuth-Yockey, PRD99(2019)114507 Fukuma-Kadoh-Matsumoto, PTEP2021(2021)123B03 Hirasawa-Matsumoto-Nishimura-Yosprakob, JHEP12(2021)011 Rigobello-Magnifico-Silvi-Montangero, arXiv:2308.04488 Bloch-Lohmayer, NPB986(2023)116032 Liu-Bhattacharya-Chandrasekharan-Gupta, arXiv:2312.17734 Hayata-Hidaka-Nishimura, JHEP07(2024)106 Asaduzzaman-Catterall-Meurice-Sakai-Toga, JHEP05(2024)195 Samberger-Bloch-Lohmayer, Lattice2024

> Horn, Phys. Lett. B100(1981)149-151 Orland-Rohrlich, NPB338(1990)647-672 Chandrasekharan-Wiese, NPB492(1997)455-471 Brower- Chandrasekharan-Wiese, PRD60(1999)094502

> > Singh-Chandrasekharan, PRD100(2019)054505

...

Starting from 2D, (1+1)D systems

• Other models: ϕ^4 theory

RG prescription within TN

Sugihara, JHEP05(2004)007 Milsted-Haegeman-Osborne, PRD88(2013)085030 Gillman-Rajantie, PRD96 (2017)094509, PRD97(2018)094505 Shimizu, Mod. Phys. Lett. A27(2012)1250035 Kadoh-Kuramashi-Nakamura-Sakai-Takeda-Yoshimura, JHEP03(2018)141, JHEP05(2019)184 Zahra-Takeda-Yamazaki, Lattice2024



Delcamp-Tilloy, PRR2(2020)033278

• Critical coupling in the continuum limit by various numerical methods

Method	Critical coupling
MPS Milsted+, PRD88(2013)085030	11.064(20)
Monte Carlo Bronzin+, PRD99(2019)034508	11.055(20)
TRG Kadoh+, JHEP05(2019)184	10.913(56)
Gilt-TNR Delcamp+, PRR2(2020)033278	11.0861(90)
MPS Vanhecke+, arXiv:2104.10564	11.09698(31)

Starting from 2D, (1+1)D systems

- Other models: Four-fermion interactions
 - Toy models for QCD
 - Gross-Neveu model
 - Relation to strongly-correlated systems
 - Aoki phase (parity-symmetry breaking phase)
 - Topological phases





Bermudez+, Ann. Phys. NY399(2018)148

Takeda-Yoshimura, PTEP2015(2015)043B01 Jünemann-Piga-Ran-Lewenstein-Rizzi-Bermudez, PRX7(2017)031057 Bermudez-Tirrito-Rizzi-Lewenstein-Hands, Ann. Phys. NY399(2018)148 Tirrito-Rizzi-Sierra-Lewenstein-Bermudez, PRB99(2019)125106 Bañuls-Cichy-Kao-Lin-Lin-Tan, PRD100(2019)094504 SA, PRD108(2023)034514 Bañuls-Cichy-Hung-Kao-Lin-Singh, arXiv:2407.11295

Aoki, PRD30(1984)2653

Starting from 2D, (1+1)D systems

- And other models
 - Zn and U(1) gauge theories
 - Topological term
 - Toward chiral gauge theories
 - O(N) sigma model, CP(N-1) model
 - Study of entanglement entropy
 - Topological term
 - Haldane's conjecture
 - SUSY on a lattice
 - Euclidean quantum gravity

Sugihara, JHEP07(2005)022 Unmuth–Yockey-Zhang-Bazavov-Meurice-Tsai, PRD98(2018)094511 Kuramashi-Yoshimura, JHEP04(2020)089 SA-Kuramashi, JHEP09(2024)086

Milsted, PRD93(2016)085012 Yang-Liu-Zou-Xie-Meurice, PRE93(2016)012138 Kawauchi-Takeda, PRD93(2016)114503 Bazavov-Meurice-Tsai-Unmuth-Yockey-Yang-Zhang, PRD96(2017)034514 Bruckmann-Jansen-Kühn, PRD99(2019)074501 Nakayama-Funcke-Jansen-Kao-Kühn, PRD105(2022)054507 Luo-Kuramashi, JHEP03(2024)020, arXiv:2406.08865 Aizawa-Takeda, Lattice2024

Kadoh-Kuramashi-Nakamura-Sakai-Takeda-Yoshimura, JHEP03 (2018) 141

Dittrich-Mizera-Steinhaus, New Jour. Phys. 18(2016)053009 Asaduzzaman-Catterall-Unmuth-Yockey, PRD102(2020)054510 Ito-Kadoh-Sato, PRD106(2022)106044

Now, moving on to the higher dimensions

- TN computations in higher dimensions (3D, 4D) are challenging
 - Gauge DoFs cannot be eliminated as in (1+1)D
 - "Curse of dimensionality"
 - As dimensions increase, so do computational costs
 - Computational memory
 - Execution time
 - It seems difficult to extend a method that is efficient in (1+1)D
 - Development of the algorithms specialized for higher dimensions is of essential importance
 - Several promising candidates

Higher-dimensional TN algorithms

Tensor Renormalization Group (TRG)

• Real-space RG on tensor networks, whose accuracy can be systematically improved by increasing the bond dimension χ

Nakayama, arXiv:2307.14191

- Higher-Order TRG (HOTRG)
 - Xie-Chen-Qin-Zhu-Yang-Xiang, PRB86(2012)045139
 - Computationally demanding: $O(\chi^{4d-1})$
- Anisotropic TRG (ATRG) Adachi-Okubo-Todo, PRB102(2020)054432
 - Less demanding than HOTRG: $O(\chi^{2d+1})$
- Triad TRG Kadoh-Nakayama, arXiv:1912.02414
 - Based on three-leg tensors: $O(\chi^{d+3})$
- Minimally-Decomposed TRG
 - Comparable accuracy to HOTRG: $O(\chi^{2d+1})$



Xie+, PRB86(2012)045139

Higher-dimensional TN algorithms

- Tensor Network State (TNS)
 - Variational wave function
 - (infinite) Projected Entangled Pair States (PEPS)

Verstraete-Cirac, arXiv:cond-mat/0407066 Cf. "Tensor Product State" by Nishino+, Prog. Theor. Phys. 105 (2001) 409-17 Corboz-Orús-Bauer-Vidal, PRB81(2010)165104

- The natural generalization of MPS
- Satisfies the area law of entanglement entropy
- Tree Tensor Network (TTN)

Shi-Duan-Vidal, PRA74(2006)022320 Tagliacozzo-Evenbly-Vidal, PRB80(2009)235127 Silvi-Giovannetti-Montangero-Rizzi-Cirac-Fazio, PRA81(2010)062335

Zaletel-Pollmann, PRL124(2020)037201

- Hierarchical tensor network w/o loops
- Similarity with the real-space RG
- Isometric TNS (isoTNS)
 - Generalization of the isometric condition of MPS to higher dimensions



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Felser+, PRX10(2020)041040



Zaletel+, PRL124(2020)037201

Higher-dimensional TN algorithms

- Gauged Gaussian PEPS (GGPEPS)
 - Deals with the gauge DoFs with the MC sampling rather than discretizing them
 - Constructs an ansatz state as PEPS
 - The probability is described by the norms of states



Bender-Emonts-Cirac, PRR5(2023)043128



Emonts- Bañuls-Cirac-Zohar, PRD102(2020)074501

• Sign-problem-free MC simulation

Zohar-Burrello, New J. Phys.18(2016)043008 Zohar-Cirac, PRD97(2018)034510

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 $[Review] \ Kelman-Borla-Gomelski-Elyovich-Roose-Emonts-Zohar, ar Xiv: 2404.13123$



Status of TN in 3D, (2+1)D

• Ising, Potts models [TRG, TNS]

Xie-Chen-Qin-Zhu-Yang-Xiang, PRB86(2012)045139 Dai-Cho-Batchelor-Zhou, PRE89(2014)062142 Wang-Xie-Cheng-Bruce-Xiang, Chin.Phys.Lett.31(2014)070503 Bloch-Lohmayer-Schweiβ-Unmuth-Yockey, Lattice2021 Jha, arXiv:2201.01789 Kadow-Pollmann-Knap, PRB107(2023)205106

SA-Kuramashi-Yoshimura, PRD104(2021)034507

Bloch-Jha-Lohmayer-Meister, PRD104(2021)094517 SA-Jha-Unmuth-Yockey, PRD110(2024)034519

Tagliacozzo-Vidal, PRB83(2011)115127 Tagliacozzo-Celi-Lewenstein, PRX4(2014)041024 Kuramashi-Yoshimura, JHEP08(2019)023 Emonts- Bañuls-Cirac-Zohar, PRD102(2020)074501 Crone-Corboz, PRB101(2020)115143 Robina-Bañuls-Cirac, PRL126(2021)050401 SA-Kuramashi, JHEP05(2022)102 Emonts-Kelman-Borla-Moroz-Gazit-Zohar, PRD107(2023)014505 Nakayama-Schneider, arXiv:2407.14226

Milde-Bloch-Lohmayer, Lattice2021

Zohar-Burrello-Wahl-Cirac, Ann. Phys. 363(2015)84 Zapp-Orús, PRD95(2017)114508 Felser-Silvi-Collura-Montangero, PRX10(2020)041040 Bender-Emonts-Cirac, PRR5(2023)043128

Zohar-Wahl-Burrello-Cirac, Ann. Phys. 374(2016)84 Cataldi-Magnifico-Silvi-Montangero, PRR6(2024)033057 Kuwahara-Tsuchiya, PTEP022(2022)093B02 Yosprakob-Okunishi, arXiv:2406.16763

- ϕ^4 theory [TRG]
- O(N) models [TRG]
- Zn gauge theories [TRG, GGPEPS, TNS]

- Infinite-coupling U(N) model [TRG]
- QED [TNS, GGPEPS]
- SU(N) gauge theories [TRG, TNS]

Z_2 gauge theories w/ TRG

- Second-order phase transition at finite temperature
 - TRG confirmed the 2D Ising universality as conjectured by Svetitsky and Yaffe
- Phase transitions in 3D Z₂ gauge-Higgs model
 - TRG confirmed the self-dual transition line

Balian-Drouffe-Itzykson, PRD11(1975)2098

Svetitsky-Yaffe, NPB210(1982)423



Kuramashi+, JHEP08(2019)023

 Resulting critical endpoint is consistent with the latest MC estimation

Method	Critical endpoint
MC Somoza+, PRX11(2021)041008	0.701
TRG SA+, JHEP05(2022)102	0.70051(7)



SA+, JHEP05(2022)102

Z₃ gauge theory w/ GGPEPS, iPEPS

• Ground state energy and string tension by GGPEPS Emonts+, PRD 102(2020)074501



• Ground state energy and string tension by iPEPS Robina+, PRL126(2021)050401





QED w/ TTN, GGPEPS

- TTN simulation w/ finite magnetic coupling
 - Ground state energy on an L=8 lattice is compared with the perturbation theory
- Confinement-deconfinement phase transition is captured by GGPEPS
- Finite-density computations are also done by GGPEPS, avoiding the sign problem





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Bender+, PRR5(2023)043128

Status of TN in 4D, (3+1)D

Ising model [TRG]

SA-Kuramashi-Yamashita-Yoshimura, PRD100(2019)054510 Sugimoto-Sasaki, Lattice2024

• ϕ^4 theory [TRG]

SA-Kadoh-Kuramashi-Yamashita-Yoshimura, JHEP09(2020)177 SA-Kuramashi-Yoshimura, PRD104(2021)034507

- Zn gauge-Higgs model at finite density [TRG]
- Infinite-coupling U(N) model [TRG]
- Nambu-Jona-Lasinio model at finite density [TRG]

SA-Kuramashi-Yamashita-Yoshimura, JHEP01(2021)121

• QED [TNS]

Magnifico-Felser-Silvi-Montangero, Nat. Commun. 12 (2021) 3600

SA-Kuramashi, JHEP05(2022)102, 10(2023)077

Milde-Bloch-Lohmayer, Lattice2021

Z₃ gauge-Higgs model w/ TRG

0.8

0.6

0.4

0.2

0.0

0.0

0.2

0.4

0.6

ß

Results by the MC and TRG at vanishing chemical potential are consistent





- First-order confinement-Higgs transition line and its endpoint are determined
 - Critical exponents are comparable to the mean-field ones also at finite density



4.0

3.0

2.0

1.0

0.0

0.2

0.4

ß

0.6

 $\langle U \rangle$

Э−Ө <U>

0.8

1.0

SA+, JHEP10(2023)077

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 $\odot \odot \chi_U$

Xυ

0.8

1.0

Infinite-coupling U(N) model w/ TRG

- Investigation of chiral symmetry breaking
 - Staggered fermion theory is reformulated as a dimer-monomer system
 - Infinite-volume limit and chiral limit
 - Chiral symmetry seems to be spontaneously broken in these limits
- Comparison with the MC simulation
 - Chiral condensate in the large-N limit
 - Difference is less than 1% btw TRG and MC

	$\lim_{N \to \infty} \frac{\langle \bar{\psi} \psi \rangle}{N}$
HT-HOTRG (D=20)	0.637(2)
Metropolis	0.64279(1)





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Nambu-Jona-Lasinio model w/ TRG

• NJL model at finite density is an effective theory of the QCD at finite density





Fukushima-Hatsuda, Rept. Prog. Phys. 74(2011)014001

- TRG has confirmed that the chiral symmetry is restored in the cold dense regime, where the MC is hindered by the severe sign problem
 - Pressure and number density have also been obtained



SA+, JHEP01(2021)121

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QED w/ TTN

Magnifico+, Nat.Commun.12(2021)3600

- The first TTN simulation of the ground state of the (3+1)D compact QED
 - Confinement properties are investigated
 - Coulomb potential in the weak coupling regime
 - Linear potential in the strong coupling regime







Status of TN in the higher dimensions

- TRG, TNS, and GGPEPS have been applied to the 3D, (2+1)D systems
 - Starting from the simplest spin models and proceedings to applications to non-Abelian gauge theories
 - Relatively large bond dimensions are available in TRG and TTN
- TRG algorithms are most employed in numerical computations of 4D, (3+1)D systems
 - Since the tensor contraction is the bottleneck in several TRG algorithms, parallelization can reduce the execution time
- Parallel computing methods specialized for individual algorithms are being developed
 SA-Kuramashi-Yamashita-Yoshimura, Lattice2019 Sakurai-Yamashita, CPC278(2022)108423

Sun-Shirakawa-Yunoki, PRB110(2024)08514 Magnifico-Cataldi-Rigobello-Majcen-Jaschke-Silvi-Montangero, arXiv:2407.03058

Future perspectives

- The Hamiltonian- and Lagrangian-based TN algorithms share the problem of how to regularize or deal with the non-Abelian DoFs
 - Not only in classical TN computations but also in quantum computations. There is potential for mutual development
 - GGPEPS is providing a new approach toward this issue
- Need to develop the schemes to analyze numerical results that include
 finite bond dimension effect
 Tagliacozzo-Oliveira-Iblisdir-Latorre, PRB78(200
 - Lessons form conformal field theories
- New TN algorithms by combining other methods
 - TN + Machine Learning

Liao-Liu-Wang-Xiang, PRX9(2019)031041 Chen-Gao-Guo-Liu-Zhao-Liao-Wang-Xiang-Li-Xie, PRB101(2020)220409 Jha-Samlodia, CPC294(2024)108941

- TN + Monte Carlo
 - From the systematic error to statistical one

Ferris, arXiv:1507.00767 Huggins-Freeman-Stoudenmire-Tubman-Whaley, arXiv:1710.03757 Zohar-Cirac, PRD97(2018)034510 Arai-Ohki-Takeda-Tomii, PRD107(2023)114515 Todo, 18th Extreme Universe Colloquium (2023) Chen-Guo-Zhang-Zhang-Deng, arXiv:2409.06538

Tagliacozzo-Oliveira-Iblisdir-Latorre, PRB78(2008)024410 Pollmann-Mukerjee-Turner-Moore, PRL102(2009)255701 Ueda-Oshikawa, PRB108(2023)024413 Huang-Chan-Kao-Chen, PRB107(2023)205123