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Approximating Many-Electron Wave Functions using Neural Networks

IMPERIAL

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Simulating Many-Electron Systems

The Many-Electron Schrödinger Equation

$$
\left(-\frac{1}{2}\sum_i \nabla_i^2 + \frac{1}{2}\sum_{i,j} \frac{e^2}{|\bm{r}_i-\bm{r}_j|} - \sum_a \sum_{i=1}^N \frac{Z_a e^2}{|\bm{d}_a-\bm{r}_i|} + \frac{1}{2}\sum_{a,b} \frac{Z_a Z_b e^2}{|\bm{d}_a-\bm{d}_b|}\right)\Psi = E\Psi
$$

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$$

The fundamental laws necessary for the mathematical treatment of a large part of physics and the whole of chemistry are thus completely known, and the difficulty lies only in the fact that application of these laws leads to equations that are too complex to be solved.

P.A.M. Dirac, Proc. Roy. Soc. A, **123**, 714 (1929)

Density Functional Theory

- Deal with the electrons one by one.
- **Effects of other electrons approximated by a mean field.**
- $\Phi \Psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N) \longrightarrow \mathcal{A} \phi_1(\mathbf{r}_1) \phi_2(\mathbf{r}_2) \dots \phi_N(\mathbf{r}_N)$

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- There exists a mean field that yields exact ground-state energies and densites.

DFT works remarkably well, but:

- is not accurate enough for room-temperature chemistry;
- cannot tell you about the correlations between electrons.

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(. . . and prove Dirac wrong?)

Neural Wave Functions

Artificial neural networks

are flexible and efficient function approximators in high-dimensional spaces.

$$
f(\mathbb{C}) \to \text{dog}
$$
\n
$$
f(\mathbb{C}) \to \text{cat}
$$

Neural Wave Functions

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A Simple Neural Network

 $\mathbf{x}^{\ell+1} = A(\mathbf{W}^{\ell} \mathbf{x}^{\ell} + \mathbf{b}^{\ell})$

Learn network parameters using automatic differentiation and back propagation.

A Simple Neural Wave Function

$$
\mathbf{x}^{\ell+1} = \mathsf{A}(\mathbf{W}^\ell\mathbf{x}^\ell + \mathbf{b}^\ell)
$$

Learn wavefunction parameters using automatic differentiation and back propagation.

$$
\mathsf{E}(\pmb{\theta}) = \int \ldots \int \, \Psi_{\pmb{\theta}}^* \, \hat{\mathsf{H}} \Psi_{\pmb{\theta}} \; \text{d} \pmb{\mathsf{r}}_1 \ldots \text{d} \pmb{\mathsf{r}}_N
$$

$$
\mathsf{E}(\pmb{\theta}) = \int \ldots \int \, \Psi_{\pmb{\theta}}^* \Psi_{\pmb{\theta}} \, \, \frac{\hat{\mathsf{H}} \Psi_{\pmb{\theta}}}{\Psi_{\pmb{\theta}}} \, \mathsf{\pmb{dr}}_1 \ldots \mathsf{\pmb{dr}}_N
$$

$$
\textbf{E}(\pmb{\theta}) = \int \ldots \int \, \left(\frac{\hat{H} \Psi_{\pmb{\theta}}}{\Psi_{\pmb{\theta}}} \right) \, \, |\Psi_{\pmb{\theta}}|^2 \; \text{d}\pmb{r}_1 \ldots \text{d}\pmb{r}_N
$$

$$
\texttt{E}(\boldsymbol{\theta}) = \mathbb{E}_{\Psi_{\boldsymbol{\theta}}^2}\left[\frac{\hat{\texttt{H}}\Psi_{\boldsymbol{\theta}}(\textbf{r}_1, \dots, \textbf{r}_\textsf{N})}{\Psi_{\boldsymbol{\theta}}(\textbf{r}_1, \dots, \textbf{r}_\textsf{N})}\right]
$$

Given an approximate ground-state wavefunction $\Psi_{\theta}(\mathbf{r}_1, \dots, \mathbf{r}_N)$, we can improve it by minimizing

$$
\text{E}(\boldsymbol{\theta}) = \mathbb{E}_{\Psi_{\boldsymbol{\theta}}^2}\left[\frac{\hat{H}\Psi_{\boldsymbol{\theta}}(\boldsymbol{r}_1, \dots, \boldsymbol{r}_N)}{\Psi_{\boldsymbol{\theta}}(\boldsymbol{r}_1, \dots, \boldsymbol{r}_N)}\right]
$$

QM comes with a built-in loss function. We can use it to learn wavefunctions without recourse to external data

Neural Variational Monte Carlo

Estimate the energy expectation value and its gradients with respect to the network parameters using Monte Carlo sampling.

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Neural Variational Monte Carlo

- Estimate the energy expectation value and its gradients with respect to the network parameters using Monte Carlo sampling.
- Adapt the weights and biases of the neural network to lower the energy.
- No training dataset is required. We are "learning from equations."

The Pauli Principle

Many-electron wavefunctions must be totally antisymmetric:

$$
\Psi_{\theta}(\mathbf{r}_1, \ldots, \mathbf{r}_i, \ldots, \mathbf{r}_j, \ldots, \mathbf{r}_N) \n= -\Psi_{\theta}(\mathbf{r}_1, \ldots, \mathbf{r}_j, \ldots, \mathbf{r}_i, \ldots, \mathbf{r}_N)
$$

for all electron labels i and j.

Hartree-Fock Theory

$$
\Psi \approx \begin{bmatrix}\n\varphi_1(\mathbf{r}_1) & \varphi_1(\mathbf{r}_2) & \dots & \varphi_1(\mathbf{r}_N) \\
\varphi_2(\mathbf{r}_1) & \varphi_2(\mathbf{r}_2) & \dots & \varphi_2(\mathbf{r}_N) \\
\vdots & \vdots & \ddots & \vdots \\
\varphi_N(\mathbf{r}_1) & \varphi_N(\mathbf{r}_2) & \dots & \varphi_N(\mathbf{r}_N)\n\end{bmatrix}
$$

FermiNet

Nothing requires the orbitals to be functions of the coordinates of a single electron:

$$
\Psi \approx \begin{pmatrix}\n\varphi_1(\mathbf{r}_1, \{\mathbf{r}_{/1}\}) & \varphi_1(\mathbf{r}_2, \{\mathbf{r}_{/2}\}) & \dots & \varphi_1(\mathbf{r}_N, \{\mathbf{r}_{/N}\}) \\
\varphi_2(\mathbf{r}_1, \{\mathbf{r}_{/1}\}) & \varphi_2(\mathbf{r}_2, \{\mathbf{r}_{/2}\}) & \dots & \varphi_2(\mathbf{r}_N, \{\mathbf{r}_{/N}\}) \\
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 $\overline{}$ $\overline{}$ $\overline{}$ $\overline{}$ $\overline{}$ $\overline{}$ $\overline{}$ $\overline{}$ $\overline{}$ $\overline{1}$

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 Ψ remains antisymmetric as long as $\varphi_{\mathsf{i}}(\mathbf{r}_{\mathsf{j}}; \{\mathbf{r}_{\mathsf{j}}\})$ is invariant under any change in the order of the arguments after **r**^j .

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(A drastic generalisation of a backflow wavefunction)

 $\overline{}$ $\overline{}$ $\overline{}$ $\overline{}$ $\overline{}$ $\overline{}$ $\overline{}$ $\overline{}$ $\overline{}$ $\overline{1}$

FermiNet and PsiFormer

FermiNet

D. Pfau, J.S. Spencer, A.G.D.G. Matthews, and W.M.C. Foulkes Phys. Rev. Res. **2**, 033429 (2020)

Psiformer

I. von Glehn, J.S. Spencer and D. Pfau arXiv:2211.13672 (2022)

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Recent Review

J. Hermann, J. Spencer, K. Choo, A. Mezzacapo, W.M.C. Foulkes, D. Pfau, G. Carleo, and F. Noé Nat. Rev. Chem. **7**, 692 (2023)

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Atoms

• 16 FermiNet determinants.

> (conventional VMC & DMC used 50–100 CSFs.)

• FermiNet consistently captures 99.7% of correlation energy.

VMC, DMC: P. Seth, P. López Ríos and R.J. Needs J. Chem. Phys. **134**, 084105 (2011)

Exact: S.J. Chakravorty et al. Phys. Rev. B **47**, 3649 (1993)

Molecules

 H_{10}

Motta et al., PRX 031059 (2017)

Chemical Reactions

Bicyclobutane → Butadiene

(Energies in kcal/mol relative to bicyclobutane)

Solids

- \bullet Simulation cell with sides $\mathbf{a}_1, \mathbf{a}_2, \mathbf{a}_3$.
- Any point in cell can be written

$$
\textbf{r}=r_1\textbf{a}_2+r_2\textbf{a}_2+r_3\textbf{a}_3
$$

• Replace position inputs r_i by explicitly periodic functions $sin(2\pi r_i)$ and $cos(2\pi r_i)$.

LiH crystal

See also Li, Li and Chen, Nat. Commun. **13**, 7895 (2022)

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Nat. Commun. **15**, 5214 (2024)

e^+ binding energy (mHa)

e^+ binding energy (mHa)

$\left[\mathsf{BeO},\mathsf{e}^+ \right]$

Wigner Crystals

At very low densities, a uniform electron liquid freezes into a

Wigner crystal

Wigner crystallisation has not yet been observed in 3D, but was recently imaged in 2D.

(ETH Zurich, July 2021)

Simulations

Continuum QMC methods cannot easily find quantum phase transitions; you usually have to guess them first.

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- Continuum QMC methods cannot easily find quantum phase transitions; you usually have to guess them first.
- FermiNet found the 3D Wigner crystal spontaneously.

Results

27 electrons in a body-centered cubic simulation cell.

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- Simulation cell is far too small to get the right transition density. Formation of a $3 \times 3 \times 3$ Wigner crystal is strongly favoured. 1.0

Phys. Rev. Lett. **130**, 036401 (2023)

Results

- 27 electrons in a body-centered cubic simulation cell.
- Simulation cell is far too small to get the right transition density. Formation of a $3 \times 3 \times 3$ Wigner crystal is strongly favoured.

Phys. Rev. Lett. **130**, 036401 (2023)

Comparing with conventional VMC and DMC simulations for the same simulation cell is nevertheless valid.

Relative Energies

Order Parameter

The Floating Crystal

The Wigner crystal was seen in the one-electron density. If we omit the one-electron input features we see a true floating crystal, but the energies are significantly higher.

The Unitary Fermi Gas

$$
\hat{\mathsf{H}} = -\frac{1}{2}\sum_{i=1}^{\mathsf{N}}\nabla_i^2 + \sum_{i=1}^{\mathsf{N}^\uparrow}\sum_{j=1}^{\mathsf{N}^\downarrow}\mathsf{U}(\textbf{r}_i^\uparrow - \textbf{r}_j^\downarrow)
$$

M. Randeria and E. Tayor, Annu. Rev. Condens. Matter Phys. **5**, 209 (2014)

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Attractive interaction U chosen such that s-wave scattering length a^s diverges.

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- Attractive interaction U chosen such that s-wave scattering length a^s diverges.
- (Range of U)/ $a_s \rightarrow 0$.
- \bullet Only remaining length scale is $1/k_F$.

M. Randeria and E. Tayor, Annu. Rev. Condens. Matter Phys. **5**, 209 (2014)

• All thermodynamic quantities depend on k_F only. For example

$$
\mathsf{E} = \xi \mathsf{E}_{\mathsf{FG}} = \xi \frac{3}{5} \frac{\hbar^2 \mathsf{k}_{\mathsf{F}}^2}{2\mathsf{m}}
$$

where ξ is a dimensionless constant known as the Bertsch parameter.

Energy as a function of number of particles

Ideas for Improvement?

Use a bigger network or a linear combination of more terms?

- Slightly increases maximum system size that works.
- Costly.

Phys. Rev. X **14**, 021030 (2024)

Ideas for Improvement?

Use a bigger network or a linear combination of more terms?

- Slightly increases maximum system size that works.
- Costly.
- Build a neural generalisation of the BCS (AGPs) wave function:

$$
\Psi^{2\mathsf{N}} = \mathcal{A}\left[\Phi(\textbf{r}_1^\uparrow,\textbf{r}_2^\downarrow)\Phi(\textbf{r}_3^\uparrow,\textbf{r}_4^\downarrow)\ldots\Phi(\textbf{r}_{2\mathsf{N}-1}^\uparrow,\textbf{r}_{2\mathsf{N}}^\downarrow)\right]
$$

Energies calculated using pairing wave function

Superfluid fraction

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Positives

- Neural wave functions are way better than other approximate wave functions for molecules with more than a few atoms.
- Rival and sometimes outperform the best quantum chemical methods, many of which have been in development for 50+ years.
- Seem to be good at dealing with unusual systems (positronic molecules; Wigner crystals) where conventional methods need painful modification.
- Might be able to deal with strong correlations.
- Might be able to discover new phases.

Questions and Negatives

- How general is the FermiNet wavefunction?
- Does it work for non-Fermi liquids?
- Size consistency and extensivity?
- Optimization is slow and sometimes problematic.
- Scaling.

Range of Ideas and Approaches

Typical Network Configuration

- **16 FermiNet determinants**
- 256 hidden units for the one-electron stream
- 32 hidden units for the two-electron stream
- Approximately 700,000 parameters
- Batch size 4096
- Training iterations 2×10^5
- Steps between parameter updates 10

Scaling

